## ON LAX OPERATORS

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ABSTRACT. We define a Lax operator as a monic pseudodifferential operator  $L(\partial)$  of order  $N \geq 1$ , such that the Lax equations  $\frac{\partial L(\partial)}{\partial t_k} = [(L^{\frac{k}{N}}(\partial))_+, L(\partial)]$ are consistent and non-zero for infinitely many positive integers  $k$ . Consistency of an equation means that its flow is defined by an evolutionary vector field. In the present paper we demonstrate that the traditional theory of the KP and the N-th KdV hierarchies holds for arbitrary scalar Lax operators.

### **CONTENTS**



### 1. Introduction

<span id="page-0-0"></span>In his seminal paper [\[Lax68\]](#page-37-1) Lax observed that the famous KdV equation

$$
u_t = \frac{1}{4}u_{xxx} + \frac{3}{2}uu_x
$$

is equivalent to an equation of the form  $\frac{\partial L(\partial)}{\partial t} = [B(\partial), L(\partial)],$  on the differential operator  $L(\partial) = \partial^2 + u$  for a certain differential operator  $B(\partial)$  of order 3. This has lead to the Lax equation approach in the theory of integrable systems. As is well known, writing an evolution equation in a Lax form allows one to construct its higher symmetries and integrals of motion.

The first beautiful application of this approach was developed by Gelfand and Dickey [\[GD76\]](#page-37-2) by considering a differential operator  $L(\partial)$  of the form

<span id="page-0-1"></span>
$$
L(\partial) = \partial^{N} + \sum_{j=1}^{N-1} u_j \partial^{N-j-1} , \quad N \ge 2 ,
$$
 (1.1)

where  $u_1, \ldots, u_{N-1}$  are the generators of the algebra of differential polynomials  $V_{N-1}$  in  $N-1$  differential variables, cf. [\(2.5\)](#page-4-1). They showed that the hierarchy of Lax equations

$$
\frac{\partial L(\partial)}{\partial t_k} = [B_k(\partial), L(\partial)], \text{ where } B_k(\partial) = (L^{\frac{k}{N}}(\partial))_+, k \in \mathbb{Z}_{\geq 1}, \qquad (1.2)
$$

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is a hierarchy of compatible evolution equations, i.e. equations of the form

$$
\frac{\partial u_j}{\partial t_k} = R_{j,k} , \quad j = 1, \ldots, N-1, \ k \in \mathbb{Z}_{\geq 1}, \ R_{j,k} \in \mathcal{V}_{N-1},
$$

which is called the N-th KdV hierarchy. Recall that the subscript  $+$  stands for the differential part of a pseudodifferential operator and compatibility means that the partial derivatives  $\frac{\partial}{\partial t_n}$  commute. The case  $N = 2$  of  $(1.1)$  is the KdV hierarchy, of which [\(1.1\)](#page-0-1) for  $k = 3$  is the KdV equation. (For arbitrary N and  $k = 1$  equation [\(1.1\)](#page-0-1) is the trivial evolution equation  $\frac{\partial u_j}{\partial t_1} = u'_j$ , and for  $k \in N\mathbb{Z}$  equation (1.1) is the zero equation  $\frac{\partial u}{\partial t_j} = 0.$ )

Note however that a Lax equation is not necessarily an evolution equation. For example, taking  $L(\partial) = \partial^3 + u$ , we have  $B_2(\partial) = (L^{\frac{2}{3}}(\partial))_+ = \partial^2$ , hence the corresponding Lax equation is

$$
\frac{\partial u}{\partial t_2} = [B_2(\partial), L(\partial)] = 2u'\partial + u''\,,
$$

which is a linear system

$$
u'=0\ ,\ \ \frac{\partial u}{\partial t_2}=u''\,.
$$

This is not an evolution equation and there exists no evolution equation on  $\mathcal{V}_1$  whose flow defines the flow of this system. This equation is thus called "inconsistent".

The general definition of "consistency" is as follows. Let  $u = (u_{\alpha})_{\alpha \in I}$  be the column vector of differential variables of  $\mathcal{V} = \mathcal{V}_{\ell}$  and let  $P = (P_{\alpha})_{\alpha \in I} \in \mathcal{V}^{\ell}$ . The evolution equation associated to  $P$  is defined as

<span id="page-1-0"></span>
$$
\frac{\partial u}{\partial t} = P. \tag{1.3}
$$

By the chain rule, this equation induces the evolution of an arbitrary  $f \in \mathcal{V}$ :  $\frac{\partial f}{\partial t} = X_P(f)$ , where  $X_P$  is the following derivation of  $\mathcal V$  commuting with  $\partial$ , called an evolutionary vector field:

$$
X_P = \sum_{\alpha \in I} \sum_{n \in \mathbb{Z}_{\geq 0}} (\partial^n P_{\alpha}) \frac{\partial}{\partial u_{\alpha}^{(n)}}.
$$

Let now  $M(\partial)$  be an  $m \times \ell$  matrix differential operator over  $\mathcal{V}$ , and let  $Q = (Q_i)_{i=1}^m \in$  $\mathcal{V}^m$ . The corresponding linear system of quasi-evolution equations is defined as:

<span id="page-1-1"></span>
$$
M(\partial)\frac{\partial u}{\partial t} = Q.
$$
\n(1.4)

We call such system *consistent* if Q lies in the image of  $M(\partial): \mathcal{V}^{\ell} \to \mathcal{V}^{m}$ , i.e. there exists  $P \in \mathcal{V}^{\ell}$  such that

$$
M(\partial)P = Q. \tag{1.5}
$$

In this case the flow, defined by the evolution equation  $(1.3)$ , is "consistent" with the linear system [\(1.4\)](#page-1-1).

A consistent linear system  $(1.4)$  is called *compatible* with another consistent linear system of quasi-evolution equations  $\tilde{M}(\partial)\frac{\partial u}{\partial \tilde{t}} = \tilde{Q}$ , such that  $\tilde{M}(\partial)\tilde{P} = \tilde{Q}$ , if P and  $\tilde{P}$  can be chosen in such a way that the corresponding evolutionary vector fields  $X_P$  and  $X_{\tilde{P}}$  commute:

$$
X_P X_{\tilde{P}} = X_{\tilde{P}} X_P . \tag{1.6}
$$

In the present paper we consider an arbitrary scalar monic pseudodifferential operator of order  $N \geq 1$  over the algebra of differential polynomials  $V_{\ell}$  in  $\ell$  differential variables  $(\ell \text{ may be infinite})$ 

<span id="page-2-1"></span>
$$
L(\partial) = \partial^{N} + \sum_{i=0}^{\infty} a_{i} \partial^{N-1-i} , \quad a_{i} \in \mathcal{V}_{\ell}.
$$
 (1.7)

Given a differential operator  $B(\partial) \in V_{\ell}[\partial]$ , the Lax equation associated to  $L(\partial)$  and  $B(\partial)$  is defined as

<span id="page-2-0"></span>
$$
\frac{\partial L(\partial)}{\partial t_B} = [B(\partial), L(\partial)].
$$
\n(1.8)

The Lax equation  $(1.8)$  is a linear system of quasi-evolution equations, which can be written down explicitly as follows. Note that, for a pseudodifferential operator  $L(\partial)$  as in [\(1.7\)](#page-2-1) and a differential operator  $B(\partial)$  of order k, one has (see Section [3](#page-10-0) for details):

$$
[B(\partial), L(\partial)] = \sum_{i=-k}^{\infty} q_i \partial^{N-1-i}, \quad q_i \in \mathcal{V}_{\ell}.
$$

Let

$$
D_{i,\alpha}(\partial) = \sum_{n=0}^{\infty} \frac{\partial a_i}{\partial u_{\alpha}^{(n)}} \partial^n \in \mathcal{V}_{\ell}[\partial], \quad i \ge 0, \, \alpha \in I,
$$

be the Frechet derivative of the coefficient  $a_i \in V_\ell$  of  $\partial^{N-1-i}$  in  $L(\partial)$ . Then, the Lax equation [\(1.8\)](#page-2-0) becomes the following linear system of quasi-evolution equations

<span id="page-2-2"></span>
$$
0 = q_i \text{ for } i \in \{-k, \dots, -1\},\tag{1.9}
$$

$$
\sum_{\alpha \in I} D_{i,\alpha}(\partial) \frac{\partial u_{\alpha}}{\partial t_B} = q_i \quad \text{for } i \ge 0.
$$
 (1.10)

Hence, this linear system of quasi-evolution equations is consistent if and only if  $q_i = 0$  for  $i < 0$  and  $Q = (q_i)_{i=0}^{\infty} \in V_{\ell}^{\infty}$  lies in the image of the matrix differential operator  $D^{(a)}(\partial) = (D_{i,\alpha}(\partial)) : \mathcal{V}_{\ell}^I \to \mathcal{V}_{\ell}^{\infty}$ .

Let  $B_k(\partial) = (L^{\frac{k}{N}}(\partial))_+$ . In Section [3](#page-10-0) we shall see that condition [\(1.9\)](#page-2-2) holds if and only if  $B(\partial)$  is, up to an adding an element of V, a linear combinations with constant coefficients of the differential operators  $B_k(\partial), k \in \mathbb{Z}_{\geq 1}$ .

Let  $\tilde{\mathcal{Z}}_L$  (respectively,  $\mathcal{Z}_L$ ) denote the set of all positive integers  $k$  for which the Lax equation  $\frac{\partial L(\partial)}{\partial t_k} = [B_k(\partial), L(\partial)]$  is non-zero (i.e. RHS  $\neq$  0) and consistent (respectively, is consistent). Clearly,  $\tilde{Z}_L \subset \mathcal{Z}_L$ . We say that  $L(\partial)$  is a Lax operator if the set  $\tilde{Z}_L$  is infinite. For example, the operators [\(1.1\)](#page-0-1) are Lax operators for every  $N \geq 2$ , while  $L(\partial) = \partial^3 + u$  is NOT.

If  $L(\partial)$  is a Lax operator, then we get the corresponding infinite hierarchy of consistent Lax equations

<span id="page-2-3"></span>
$$
\frac{\partial L(\partial)}{\partial t_k} = [B_k(\partial), L(\partial)], \quad k \in \mathcal{Z}_L.
$$
\n(1.11)

One of the basic results of the theory of integrable systems is that the hierarchy [\(1.11\)](#page-2-3) is compatible (see Proposition [3.2\)](#page-11-0). (Note that, if  $B(\partial) \in \mathcal{V}$ , then the Lax equation [\(1.8\)](#page-2-0) is consistent, but it may fail to be compatible with the equations of the hierarchy [\(1.11\)](#page-2-3).)

The next important development in the theory of Lax equations is the work of Sato [\[Sat81\]](#page-37-3) and his disciples [\[DJKM81,](#page-37-4) [DJKM83\]](#page-37-5) on the KP hierarchy and its analogues and reductions. The KP hierarchy is defined as the following system of Lax equations on the Sato operator  $L(\partial) = \partial + \sum_{i=1}^{\infty} u_i \partial^{-i}$ , where  $u_1, u_2, \ldots$  are the differential variables of  $\mathcal{V}_{\infty}$ :

$$
\frac{\partial L(\partial)}{\partial t_k} = [B_k(\partial), L(\partial)], \quad k \in \mathbb{Z}_{\geq 1},
$$
\n(1.12)

for which  $\mathcal{Z}_L = \mathbb{Z}_{\geq 1}$ .

Around the same time Drinfeld and Sokolov [\[DS85\]](#page-37-6) associated to each simple Lie algebra  $\mathfrak g$  and its principal nilpotent element f a hierarchy of bi-Hamiltonian PDE. In particular, for each classical g they constructed the corresponding Lax operator  $L(\partial)$ . In the case  $\mathfrak{g} = \mathfrak{sl}_N$ ,  $L(\partial)$  is the Gelfand-Dickey operator [\(1.1\)](#page-0-1) so that the corresponding hierarchy is the *n*-th KdV hierarchy. For  $g = s\mathfrak{o}_{2n+1}$  (respectively  $\mathfrak{sp}_{2n}$ ,  $L(\partial) \circ \partial$  (resp.  $L(\partial)$ ) is the "generic" skewadjoint (resp. selfadjoint) monic differential operator of order  $2n + 1$  (resp. 2n), while, for  $\mathfrak{g} = \mathfrak{so}_{2n}$ ,  $L(\partial) \circ \partial$  is the sum of the "generic" skewadjoint differential operator of order  $2n - 1$  and the pseudodifferential operator  $u\partial^{-1} \circ u$ .

The ideas of [\[GD76\]](#page-37-2) and [\[DS85\]](#page-37-6) were further developed in [\[DSKV15,](#page-37-7) [DSKV16a,](#page-37-8) [DSKV18\]](#page-37-9). In particular, in [\[DSKV16a\]](#page-37-8) (respectively [\[DSKV18\]](#page-37-9)) to each nilpotent element of the Lie algebra  $\mathfrak{g} = \mathfrak{sl}_N$  (resp.  $\mathfrak{so}_N$  and  $\mathfrak{sp}_N$ ) corresponding to a partition p of N, we constructed an  $r \times r$ -matrix monic pseudodifferential operator  $L_n(\partial)$  of order p, where p is the maximal part of p and r is its multiplicity. It follow from [\[DSKV16a\]](#page-37-8) (resp. [\[DSKV18\]](#page-37-9)) that  $L(\partial) = L_p(\partial)$  is a Lax operator with  $\mathcal{Z}_L = \mathbb{Z}_{\geq 1}$ (resp.  $\mathcal{Z}_L = 1+2\mathbb{Z}_{\geq 0}$ ), and that the Lax equations associated to it are Hamiltonian for the Poisson structure of the corresponding W-algebra  $\mathcal{W}(\mathfrak{g}, p)$ .

The simplest Lax operator  $L_p(\partial)$  corresponds to the partition  $\underline{p} = N$ , consisting of one part N, which is associated to the principal nilpotent element of  $\mathfrak{sl}_N$  (resp.  $\mathfrak{so}_N$  for odd N and  $\mathfrak{sp}_N$  for even N). In this cases  $L_p(\partial)$  coincides with [\(1.1\)](#page-0-1), the Lax operator for the N-th Gelfand-Dickey hierarchy (resp. the Lax operator associated to the  $\mathfrak{so}_N$  Drinfeld-Sokolov hierarchies composed with  $\partial$ , and the Lax operator associated to the  $\mathfrak{sp}_N$  Drinfeld-Sokolov hierarchies). For  $\mathfrak{g} = \mathfrak{so}_N$ , N even,  $p = (N-1, 1)$  and  $L_p(\partial)$  is a certain pseudodifferential operator [\[DS85\]](#page-37-6).

Another well-known example in the case of  $\mathfrak{g} = \mathfrak{sl}_N$  corresponds to the partition  $p = (p, 1^{N-p})$ , where  $N > p$ . In this case, after a reduction by non-evolving variables,  $L_p(\partial) = L_{(p)}(\partial) + \sum_{i=1}^{N-p} p_i \partial^{-1} \circ q_i$ , which is therefore a Lax operator and the corresponding hierarchy  $(1.11)$  of Lax equations is a subsystem of a hierarchy of Hamiltonian equations [\[CDSKVvdL20,](#page-37-10) Example 14.1]. This produces the wellknown  $(N-p)$ -vector p-constrained KP hierarchy [\[KSS91,](#page-37-11) [Che92,](#page-37-12) [KS92,](#page-37-13) [SS93\]](#page-37-14). The case  $p = 2$  produces the  $(N-2)$ -component Yajima-Oikawa hierarchy [\[Y076,](#page-38-0) [Ma81\]](#page-37-15).

In the present paper we establish a series of results in the generality of arbitrary scalar Lax operators, similar to the traditional results for the Sato Lax operator. Our proofs are often identical to those of the Kyoto school [\[DJKM83,](#page-37-5) [Shi86\]](#page-37-16) and of the book [\[Dic03\]](#page-37-17), sometimes they are simpler and more rigorous. We show that the hierarchy of Lax equations, associated to a Lax operator  $L(\partial)$  is compatible (Proposition [3.2\)](#page-11-0), and is equivalent to

- (i) Zakharov-Shabat equations  $[2S74]$  (Proposition [4.1\(](#page-0-2)a));
- (ii) complementary Zakharov-Shabat equations (Proposition [4.1\(](#page-0-2)b));
- (iii) Sato equations on the dressing operator (Theorem [5.1\)](#page-20-1);
- (iv) the linear problem on the wave function (Theorem  $5.3$ );
- (v) the bilinear equation on the wave function (Theorem [5.4\)](#page-22-0).

As for the Sato Lax operator, this leads to the construction of the tau-function for  $L(\partial)$ . The tau-function for all Lax operators  $L_p(\partial)$ , for all partitions p, were constructed in [\[CDSKVvdL20\]](#page-37-10) as those for the r-component KP hierarchy satisfying a simple constraint. Of course, the importance of the tau-function  $\tau$  for the Lax operator  $L(\partial)$  stems from the fact that, as for KP, all solutions of the corresponding hierarchy [\(1.11\)](#page-2-3) can be expressed via  $\tau$  (see formula [\(6.5\)](#page-25-0)).

For the Sato Lax operator formula [\(6.5\)](#page-25-0) establishes an essentially bijective correspondence between tau-functions and the solutions of the KP-hierarchy. For

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arbitrary p-reductions of the KP hierarchy the corresponding tau-functions satisfy a simple constraint, which allowed for their explicit construction in [\[CDSKVvdL20\]](#page-37-10).

Unfortunately, it is still an open problem for an arbitrary Lax operator to find the constraints on the tau-function imposed by the constraints on the coefficients of  $L(\partial)$ . This problem has been solved for the Lax operators of the BKP and CKP hierarchies [\[DJKM81,](#page-37-4) [CW13,](#page-37-18) [KZ20\]](#page-37-19), but not, for example, for the Lax operator corresponding to the KN equation discovered by Sokolov [\[Sok84\]](#page-38-2).

The above discussion can be extended to  $r \times r$  matrix pseudodifferential operators  $L(\partial)$ . This will be treated in a forthcoming publication.

Throughout the paper the base field  $\mathbb F$  is a field of characteristic zero.

# 2. Algebraic setup

<span id="page-4-5"></span><span id="page-4-0"></span>2.1. **Functions on space-time.** Throughout the paper we let  $\mathcal F$  be a given commutative, associative, unital algebra over  $\mathbb F$  endowed with commuting derivations

$$
\partial\,\Big(\frac{\partial}{\partial x}\Big),\,\frac{\partial}{\partial t_k}\,:\,\mathcal{F}\to\mathcal{F}\,,\quad k\in\mathcal{Z}\,,
$$

indexed by an index set  $\mathcal{Z} \subset \mathbb{Z}_{\geq 1}$ .

The elements of  $\mathcal F$  are called *functions on space-time* (or simply *functions*), and will be usually denoted as  $f = f(x, t) = f(x, t_k, k \in \mathcal{Z})$  (In the usual terminology, the "space variable" is x, and there are many "time variables"  $t_k, k \in \mathcal{Z}$ .)

We assume that the common kernel of all space and time derivatives is the base field  $\mathbb{F}$ :

$$
\mathbb{F} = \left\{ c \in \mathcal{F} \, \middle| \, \partial c = 0 \text{ and } \frac{\partial c}{\partial t_k} = 0 \text{ for all } k \in \mathcal{Z} \right\}. \tag{2.1}
$$

We also assume that F is endowed with a surjective algebra homomorphism  $\mathcal{F} \to \mathbb{F}$ , restricting to the identity map on  $\mathbb{F} \subset \mathcal{F}$ , which we shall call the *evaluation at*  $x = 0, t = 0$ , and we shall denote as

<span id="page-4-3"></span>
$$
f = f(x, t) \to f(0) = f|_{x=0, t=0} \in \mathbb{F}.
$$
 (2.2)

<span id="page-4-6"></span>**Definition 2.1.** Given elements  $g, f_k \in \mathcal{F}, k \in \mathcal{Z}$ , consider a system of equations on the unknown function  $\varphi \in \mathcal{F}$ :

<span id="page-4-2"></span>
$$
\partial \varphi = g \; , \; \frac{\partial \varphi}{\partial t_k} = f_k \text{ for all } k \in \mathcal{Z} \, . \tag{2.3}
$$

The system  $(2.3)$  is called *compatible* if the following conditions hold:

<span id="page-4-7"></span>
$$
\frac{\partial g}{\partial t_k} = \partial f_k \text{ for all } k \in \mathcal{Z}, \text{ and } \frac{\partial f_k}{\partial t_h} = \frac{\partial f_h}{\partial t_k} \text{ for all } h, k \in \mathcal{Z}. \tag{2.4}
$$

The algebra of functions on space-time F is said to be *integrable* if, for every  $c \in \mathbb{F}$ and every compatible system of equations [\(2.3\)](#page-4-2), there exists a unique solution  $\varphi \in \mathcal{F}$ such that  $\varphi(0) = c$ .

**Example 2.2.** The algebra  $\mathbb{F}[x, t_1, t_2, t_3, \ldots]$  of polynomials in infinitely many variables, and the algebra  $\mathbb{F}[[x, t_1, t_2, t_3, \dots]]$  of formal power series, are both integrable.

<span id="page-4-4"></span>2.2. Unknown (dependent) variables. We let  $u_1, \ldots, u_\ell$  ( $\ell$  may be infinite) be the "unknown functions" on space-time (=dependent variables), and we let  $V_{\ell}$  be the algebra of differential polynomials in the variables  $u_{\alpha}, \alpha \in I = \{1, \ldots, \ell\},\$ 

<span id="page-4-1"></span>
$$
\mathcal{V}_{\ell} = \mathbb{F}[u_{\alpha}^{(n)} \mid \alpha \in I, n \in \mathbb{Z}_{\geq 0}]. \tag{2.5}
$$

It is a differential algebra with respect to the derivation  $\partial$  defined by  $\partial u_{\alpha}^{(n)} = u_{\alpha}^{(n+1)},$  $\alpha \in I, n \in \mathbb{Z}_{\geq 0}$ . Note that the partial derivatives  $\frac{\partial}{\partial u_{\alpha}^{(n)}}, \alpha \in I, n \in \mathbb{Z}_{\geq 0}$ , are commuting derivations of  $\mathcal{V}_{\ell}$  which satisfy the following commutation relations:

$$
\left[\frac{\partial}{\partial u_{\alpha}^{(n)}}, \partial\right] = \frac{\partial}{\partial u_{\alpha}^{(n-1)}}, \ \ \alpha \in I, \ n \in \mathbb{Z}_{\geq 0}, \tag{2.6}
$$

where the RHS is considered to be 0 for  $n = 0$ .

By the universal property of the algebra of differential polynomials, for every collection of functions on space-time  $f_\alpha \in \mathcal{F},$   $\alpha \in I,$  there exists a unique differential algebra homomorphism  $\mathcal{V}_\ell \to \mathcal{F}$ , mapping  $u_\alpha^{(n)} \mapsto \partial^n f_\alpha$ ,  $\alpha \in I$ ,  $n \in \mathbb{Z}_{\geq 0}$ , which we shall call *evaluation at*  $u = f$ , and we shall denote as

<span id="page-5-3"></span>
$$
P = P(u, u', u'', \dots) \mapsto P(f) = P(f, f', f'', \dots) \in \mathcal{F}.
$$
 (2.7)

Remark 2.3. The results of the present paper can be generalized to the case when  $\mathcal{V}_{\ell}$ is replaced by an algebra of differential functions extending the algebra of differential polynomials [\(2.5\)](#page-4-1) [\[BDSK09\]](#page-37-20). On the other hand, for the KP hierarchy and all the hierarchies arising from classical affine  $W$ -algebras, the underlying differential algebra of unknown functions is an algebra of differential polynomials (in infinitely many variables for the KP hierarchy, and finitely many variables for  $W$ -algebras).

2.3. Consistent linear systems of quasi-evolution equations. Let  $\mathcal{V} = \mathcal{V}_{\ell}$  for simplicity of notation. Recall that an *evolutionary vector field* on  $\mathcal V$  is a derivation  $X: V \to V$  commuting with  $\partial$ . It is immediate to see that all evolutionary vector fields on the algebra  $V$  are of the form

<span id="page-5-1"></span>
$$
X_P = \sum_{\alpha \in I} \sum_{n=0}^{\infty} (\partial^n P_{\alpha}) \frac{\partial}{\partial u_{\alpha}^{(n)}},\tag{2.8}
$$

for some  $P = (P_\alpha)_{\alpha \in I}, P_\alpha \in \mathcal{V}$ . Hence, we have a bijective map from  $\mathcal{V}^\ell$  to the space of evolutionary vector fields, mapping  $P$  to  $X_P$  . Note that  $[X_P,X_Q]=X_{[P,Q]},$ where

<span id="page-5-2"></span>
$$
[P,Q] = X_P(Q) - X_Q(P). \tag{2.9}
$$

For  $f \in \mathcal{V}$ , we define its Frechet derivative  $D_{\alpha}^{(f)}(\partial) \in \mathcal{V}[\partial], \alpha \in I$ , as

<span id="page-5-4"></span>
$$
D_{\alpha}^{(f)}(\partial) = \sum_{n=0}^{\infty} \frac{\partial f}{\partial u_{\alpha}^{(n)}} \partial^{n}.
$$
\n(2.10)

Obviously, for  $P = (P_{\alpha})_{\alpha \in I} \in \mathcal{V}^I$  and  $f \in \mathcal{V}$ , we have:

<span id="page-5-5"></span>
$$
X_P(f) = \sum_{\alpha \in I} D_{\alpha}^{(f)}(\partial) P_{\alpha} . \tag{2.11}
$$

Let  $u = \{u_\alpha\}_{\alpha \in I}$  be the column vector of generators and let  $P = (P_\alpha)_{\alpha \in I} \in \mathcal{V}^\ell$ . By definition, an *evolution equation* on  $V$  has the form

$$
\frac{\partial u}{\partial t} = P. \tag{2.12}
$$

By the chain rule,  $\frac{\partial}{\partial t}$  extends uniquely to the evolutionary vector field  $\frac{\partial}{\partial t} = X_P$  on  $\nu$ 

We generalize the notion of evolution equation as follows: a *linear system of* quasi-evolution equations on  $V$  is a system of the form

<span id="page-5-0"></span>
$$
M(\partial)\frac{\partial u}{\partial t} = Q\,,\tag{2.13}
$$

where  $Q = (Q_j)_{j \in J}$ , with  $Q_j \in V$ ,  $j \in \{1, ..., m\}$ , and  $M(\partial)$  is an  $m \times \ell$  matrix differential operator over V. It is assumed that  $M(\partial)$  has only finitely many nonzero entries in each row if  $\ell$  is infinite, so that  $M(\partial)$  defines a linear map  $M(\partial)$ :

 $\mathcal{V}^{\ell} \to \mathcal{V}^{m}$ . Of course, an evolution equation is a special case of a linear system  $(2.13)$ , where  $M(\partial) = \mathbb{1}$  is the identity matrix.

<span id="page-6-4"></span>**Definition 2.4.** A linear system of evolution equations [\(2.13\)](#page-5-0) is *consistent* if  $Q \in$ Im  $M(\partial) \subset \mathcal{V}^m$ .

Suppose, for example, that the matrix differential operator  $M(\partial)$  has a right inverse  $A(\partial)B(\partial)^{-1}$ , where  $A(\partial)$  is an  $\ell \times m$  matrix differential operator and  $B(\partial)$  is an invertible  $m \times m$  matrix differential operator. In other words  $M(\partial)A(\partial) = B(\partial)$ is invertible as a pseudodifferential operator, i.e its inverse lies in  $\text{Mat}_{m\times m}\,\mathcal{V}((\partial^{-1})).$ If, moreover, the vector  $Q = (Q_i)_{i \in J}$  lies in the image of  $B(\partial)$ , then the system [\(2.13\)](#page-5-0) is obviously consistent.

We also define a *hierarchy of linear systems of quasi-evolution equations* as a collection of linear systems,

<span id="page-6-0"></span>
$$
M_k(\partial)\frac{\partial u}{\partial t_k} = Q_k \ , \ Q_k \in \mathcal{V}^m \ , k \in \mathcal{Z} \ , \tag{2.14}
$$

parametrized by a set  $\mathcal{Z} \subset \mathbb{Z}_{\geq 1}$ . Suppose that each equation [\(2.14\)](#page-6-0) of the hierarchy is consistent for every  $k \in \mathcal{Z}$ , i.e. there exists  $P_k \in \mathcal{V}^{\ell}$  such that

<span id="page-6-1"></span>
$$
M_k(\partial)P_k = Q_k, \ k \in \mathcal{Z} \,. \tag{2.15}
$$

We also let  $X_k = X_{P_k} : \mathcal{V} \to \mathcal{V}$  be the corresponding evolutionary vector fields as in [\(2.8\)](#page-5-1).

**Definition 2.5.** The hierarchy [\(2.14\)](#page-6-0) is *compatible* if we can choose elements  $P_k \in$  $\mathcal{V}^{\ell}, k \in \mathcal{Z}$ , so that  $(2.15)$  holds and

$$
X_k(P_h) = X_h(P_k) \text{ for all } h, k \in \mathcal{Z}. \tag{2.16}
$$

By equation [\(2.9\)](#page-5-2), this means that the evolutionary vector fields  $X_k$ ,  $k \in \mathcal{Z}$ , commute:  $[X_k, X_h] = 0$  for  $k, h \in \mathcal{Z}$ .

Definition 2.6. A *solution* of the hierarchy of linear systems of quasi-evolution equations [\(2.14\)](#page-6-0) is a collection of functions  $\varphi = {\varphi_{\alpha} \mid \alpha \in I} \subset \mathcal{F}$ , such that

$$
M_k(\varphi; \partial) \frac{\partial \varphi}{\partial t_k} = Q_k(\varphi), \quad k \in \mathcal{Z}, \tag{2.17}
$$

where  $M_k(\varphi; \partial)$  and  $Q_k(\varphi)$  are obtained by applying the evaluation map [\(2.7\)](#page-5-3).

Remark 2.7. Note that the consistency of a linear system of quasi-evolution equations is NOT a necessary condition for the existence of solutions. For example, the hierarchy of equations, corresponding to  $L(\partial) = \partial^3 + u$ ,

$$
0 = 0 \frac{\partial u}{\partial t_1} = u' , \quad \frac{\partial u}{\partial t_2} = 0 ,
$$

is not consistent, but it admits the solution  $u(x, t_1, t_2) = 1 \in \mathcal{F}$ .

2.4. Pseudodifferential operators. We will consider the algebra  $V((\partial^{-1}))$  of scalar pseudodifferential operators with coefficients in  $\mathcal V$ . Given such an operator

<span id="page-6-3"></span>
$$
A(\partial) = \sum_{i=-\infty}^{N} a_i \partial^i , \quad a_i \in \mathcal{V}, \tag{2.18}
$$

its symbol is defined as  $A(z) = \sum_{i=-\infty}^{N} a_i z^i \in \mathcal{V}((z^{-1}))$ . Recall that the product of two pseudodifferential operators  $A(\partial)$  and  $B(\partial)$  is defined via their symbols by

<span id="page-6-2"></span>
$$
(AB)(z) = (A \circ B)(z) = A(z + \partial)(B(z)).
$$
\n(2.19)

From  $(2.19)$  we see that the subspace  $V[\partial] \subset V((\partial^{-1}))$  of differential operators is a subalgebra. Here and further, a Laurent series involving negative powers of  $z + \partial$  is always considered to be expanded using geometric series expansion in the domain of large z.

Let  $A(\partial)$  be as in [\(2.18\)](#page-6-3). We denote by  $A^*(\partial) = \sum_i (-\partial)^i \circ a_i \in \mathcal{V}((\partial^{-1}))$  its i formal adjoint, by  $A(\partial)_+ = \sum_{i=0}^N a_i \partial^i \in \mathcal{V}[\partial]$  its differential part, by  $A(\partial)_- =$  $\sum_{i=-\infty}^{-1} a_i \partial^i \in \mathcal{V}[[\partial^{-1}]]\partial^{-1}$  its singular part.

The following notation will be used throughout the paper: given  $A(\partial) \in V((\partial^{-1}))$ as in  $(2.18)$  and  $b, c \in V$ , we let:

<span id="page-7-0"></span>
$$
A(z+x)\left(\big|_{x=\partial}b\right)c = \sum_{i=-\infty}^{N} a_i((z+\partial)^i b)c \in \mathcal{V}.
$$
 (2.20)

Using the notation  $(2.20)$ , we can rewrite

<span id="page-7-1"></span>
$$
A^*(z) = (|_{x=\partial} A(-z-x)) = \sum_{i=-\infty}^{N} (-z-\partial)^i a_i,
$$
 (2.21)

and the RHS of equation [\(2.19\)](#page-6-2) can be rewritten as  $A(z+x)(|_{x=\partial}B(z))$ .

Furthermore, for  $A(\partial) \in \mathcal{V}((\partial^{-1}))$  as in [\(2.18\)](#page-6-3) (respectively, its symbol  $A(z) \in$  $\mathcal{V}((z^{-1}))$  we define its residue as  $\operatorname{Res}_{\partial} A(\partial) = a_{-1}$  (respectively,  $\operatorname{Res}_{z} A(z) = a_{-1}$ ). For a series involving negative powers of  $z \pm w$  we shall use the notation  $i_z$  or  $i_w$  to denote geometric series expansion in the domain of large  $z$  or of large  $w$  respectively. For example,  $\iota_z(z-w)^{-1} = \sum_{n \in \mathbb{Z}_+} z^{-n-1} w^n$ . For  $A(z) \in \mathcal{V}((z^{-1}))$  as above, we have

<span id="page-7-6"></span>
$$
\operatorname{Res}_{z} A(z) \iota_{z}(z-w)^{-1} = A(w)_{+} , \quad \operatorname{Res}_{z} A(z) \iota_{w}(z-w)^{-1} = -A(w)_{-} . \quad (2.22)
$$

The order of a pseudodifferential operator  $A(\partial)$  as in [\(2.18\)](#page-6-3) is ord $(A) = N$  if  $a_N \neq 0$ . We also say that  $A(\partial)$  is monic if  $a_N = 1$ . Note that, if  $A(\partial)$  and  $B(\partial) \in V((\partial))$  have orders ord $(A) = N$  and ord $(B) = M$  respectively, then

<span id="page-7-5"></span>
$$
\text{ord}(A \circ B) = M + N \text{ and } \text{ord}([A, B]) \le M + N - 1,\tag{2.23}
$$

where  $[A, B]$  denotes the commutator of the pseudodifferential operators A and B. The following results will be important in Section [3.](#page-10-0)

<span id="page-7-2"></span>Lemma 2.8. [\[DSKV16,](#page-37-21) Lem.2.1a)] Given pseudodifferential operators  $A(\partial), B(\partial) \in$  $\mathcal{V}((\partial^{-1}))$ , we have

$$
\operatorname{Res}_{z} A(z)B^*(-z) = \operatorname{Res}_{z}(A \circ B)(z).
$$

*Proof.* First, observe that, for an arbitrary  $f(z) \in \mathcal{V}((z^{-1}))$ , Res<sub>z</sub>  $f(z+x)$ , expanded in the domain  $|z| > |x|$ , is independent of x. Next, observe that replacing z with  $-z-x$  in equation [\(2.21\)](#page-7-1), we have the identity  $A(z) = \left(\big|_{x=\partial} A^*(-z-x)\right)$ . It follows that

$$
\operatorname{Res}_{z} A(z) B^*(-z) = \operatorname{Res}_{z} A(z+x) \left( \Big|_{x=\partial} B^*(-z-x) \right)
$$
  
= 
$$
\operatorname{Res}_{z} A(z+\partial) B(z) = \operatorname{Res}_{z} (A \circ B)(z).
$$

<span id="page-7-4"></span>**Lemma 2.9.** Given pseudodifferential operators  $A(\partial), B(\partial) \in \mathcal{V}((\partial^{-1}))$ , the equation  $\text{Res}_{z}((z+\partial)^{n}A(z))B^{*}(-z) = 0$  holds for all  $n \geq 0$  if and only if  $(A \circ B)(\partial) =$ 0.

Proof. By Lemma [2.8,](#page-7-2) we have

<span id="page-7-3"></span>
$$
Res_z((z+\partial)^n A(z))B^*(-z) = Res_z(z+\partial)^n A(z+\partial)B(z)
$$
  
= Res\_z(z+\partial)^n (A \circ B)(z) = Res\_z z^n (A \circ B)^\*(-z). (2.24)

For the first equality of  $(2.24)$  we applied Lemma [2.8](#page-7-2) to the pair of pseudodifferential operators  $\partial^n A(\partial)$  and  $B(\partial)$ , while for the last equality of [\(2.24\)](#page-7-3) we applied Lemma [2.8](#page-7-2) to the pair  $\partial^n$  and  $A(\partial)B(\partial)$ . Obviously, the RHS of [\(2.24\)](#page-7-3) is 0 for every  $n \geq 0$ if and only if  $(A \circ B)^*(-z)$  is a polynomial in z, which is equivalent to saying that  $(A \circ B)^*(\partial) = 0$ , which in turn is equivalent to saying that  $(A \circ B)(\partial) = 0$ . □

<span id="page-8-0"></span>Lemma 2.10. [\[Adl79\]](#page-37-22),[\[DSKV16,](#page-37-21) Lem.2.1b)] For  $A(\partial), B(\partial) \in V((\partial^{-1}))$ , we have

$$
\operatorname{Res}_{\partial}[A(\partial), B(\partial)] \in \partial \mathcal{V}.
$$

*Proof.* It suffices to prove the claim for  $A(\partial) = a\partial^i$ ,  $B(\partial) = b\partial^j$ , with  $a, b \in V$  and  $i, j \in \mathbb{Z}$ . We have

$$
\text{Res}_{\partial}[a\partial^i, b\partial^j] = {i \choose i+j+1}ab^{(i+j+i)} - {j \choose i+j+1}ba^{(i+j+1)},
$$

where the binomial coefficient  $\begin{pmatrix} p \\ q \end{pmatrix}$  is intended to be 0 for negative q. Note that  $ba^{(m)} \equiv (-1)^m ab^{(m)} \mod \partial \mathcal{V}$ . The claim thus follows from the identity  $\binom{j}{i+j+1} =$  $(-1)^{i+j+1} {i \choose i+j+1}$ . □

Given a collection of functions  $f_{\alpha} \in \mathcal{F}$ ,  $\alpha \in I$ , we have the homomorphism  $\mathcal{V}((\partial^{-1})) \to \mathcal{F}((\partial^{-1}))$  extending the evaluation map  $(2.7)$ , which we shall denote as

<span id="page-8-4"></span>
$$
A = A(\partial) \mapsto A(f; \partial) \in \mathcal{F}((\partial^{-1})). \tag{2.25}
$$

Of course all the results proved for the algebra  $\mathcal{V}((\partial^{-1}))$  (including Lemmas [2.8,](#page-7-2) [2.9](#page-7-4) and [2.10](#page-8-0) above) hold for the algebra  $\mathcal{F}((\partial^{-1}))$ .

# 2.5.  $N$ -th root of a monic pseudodifferential operator of order  $N$ .

<span id="page-8-2"></span>**Lemma 2.11.** Let  $L(\partial) = \partial^N + a_0 \partial^{N-1} + a_1 \partial^{N-2} + \cdots \in \mathcal{V}((\partial^{-1}))$  be a monic pseudodifferential operator of order  $N \geq 1$ .

- (a) The inverse  $L^{-1}(\partial) \in V((\partial^{-1}))$  exists (and is unique) and it is a monic pseu $dodifferential operator of order -N.$
- $(b)$  There exists a unique monic pseudodifferential operator of order 1, denoted  $L^{\frac{1}{N}}(\partial),$  which is an N-th root of  $L(\partial),$  i.e. such that

<span id="page-8-1"></span>
$$
\left(L^{\frac{1}{N}}(\partial)\right)^N = L(\partial).
$$
\n(2.26)

*Proof.* (See [\[Dic03\]](#page-37-17)) Part (a) is clear, since  $L^{-1}(\partial)$  can be obtained by geometric series expansion. For part (b), let

$$
L^{\frac{1}{N}}(\partial) = \partial + b_0 + b_1 \partial^{-1} + b_2 \partial^{-2} + \dots
$$

The equation  $(L^{\frac{1}{N}}(\partial))^{N} = L(\partial)$  translate, by looking at the coefficients of  $\partial^{-i}$ , to a system of equations of the form

$$
Nb_i + P_i = a_i , i \geq 0,
$$

where  $P_0 = 0$  and, for  $i \geq 1$ ,  $P_i$  is a polynomial with constant coefficients in the variables  $b_j$ , with  $0 \leq j < i$ , and their derivatives. (In fact,  $P_i$  is homogeneous with respect to the grading defined by  $\deg b_i^{(n)} = 1 + i + n$ .) Clearly, such system can be solved recursively and it admits a unique solution. □

<span id="page-8-3"></span>**Lemma 2.12.** Let  $D_1$  and  $D_2$  be derivations of  $\mathcal{V}((\partial^{-1}))$  and let  $L(\partial)$  be a monic pseudodifferential operator of order N. If  $D_1(L(\partial))=D_2(L(\partial))$ , then  $D_1(L^{\frac{k}{N}}(\partial))=$  $D_2(L^{\frac{k}{N}}(\partial))$  for every  $k \in \mathbb{Z}$ .

*Proof.* By the Leibniz rule, it suffices to prove the claim for  $k = 1$ . We have

$$
0 = D_1(L(\partial)) - D_2(L(\partial)) = D_1((L^{\frac{1}{N}}(\partial))^N) - D_2((L^{\frac{1}{N}}(\partial))^N)
$$
  
= 
$$
\sum_{h=0}^{N-1} L^{\frac{h}{N}}(\partial) \left( D_1(L^{\frac{1}{N}}(\partial)) - D_2(L^{\frac{1}{N}}(\partial)) \right) L^{\frac{N-h-1}{N}}(\partial).
$$

Since  $L^{\frac{1}{N}}(\partial)$  is monic, the leading coefficient of the RHS is N times the leading coefficient of  $D_1(L^{\frac{1}{N}}(\partial)) - D_2(L^{\frac{1}{N}}(\partial))$ , which must therefore be zero.

<span id="page-9-1"></span>2.6. Action of pseudodifferential operators on oscillating functions. The algebra of differential operators  $V[\partial]$  acts naturally on V, and likewise the algebra of differential operators over  $\mathcal F$  acts naturally on  $\mathcal F$ . On the other hand, pseudodifferential operators do not act in any way on  $V$  or  $\mathcal{F}$ .

To overcome this problem suppose that  $\mathcal F$  is an algebra of functions on spacetime, with space variable x and time variables  $t_k, k \in \mathcal{Z}$ . In this case, we define the space of oscillating functions as

$$
\mathcal{F}((z^{-1}))e^{z\cdot t} = \left\{S(z)e^{z\cdot t} \, \middle| \, S(z) \in \mathcal{F}((z^{-1})) \right\},\,
$$

where  $e^{z \cdot t}$  is just a formal symbol, defined by the following rules

<span id="page-9-0"></span>
$$
\frac{\partial}{\partial x}e^{z\cdot t} = ze^{z\cdot t}, \quad \frac{\partial}{\partial t_k}e^{z\cdot t} = z^k e^{z\cdot t}, \quad k \in \mathcal{Z}, \quad e^{z\cdot t}|_{x=0, t=0} = 1. \tag{2.27}
$$

Namely, we should think of  $z \cdot t$  as  $zx + \sum_{k \in \mathcal{Z}} z^k t_k$ .

We have a natural representation of the algebra  $\mathcal{F}((\partial^{-1}))$  on the space of oscillating functions, given by

$$
P(\partial)(S(z)e^{z \cdot t}) = (P \circ S)(z)e^{z \cdot t}, \text{ for every } P(\partial), S(\partial) \in \mathcal{F}((\partial^{-1}))
$$

On the space of oscillating functions we have a well-dened (commuting) action of all time derivatives  $\frac{\partial}{\partial t_k}$ ,  $k \in \mathcal{Z}$ , induced by their action on  $\mathcal F$  and by [\(2.27\)](#page-9-0):

<span id="page-9-3"></span>
$$
\frac{\partial}{\partial t_k}(S(z)e^{z\cdot t}) = \left(\frac{\partial S(z)}{\partial t_k} + z^k S(z)\right) e^{z\cdot t}.
$$
\n(2.28)

Recall that on the algebra of functions  $\mathcal F$  we have the evaluation map  $\mathcal F \to \mathbb F$ defined in  $(2.2)$ . It induces an *evaluation map* on the space of oscillating functions  $\mathcal{F}((z^{-1}))e^{z \cdot t} \to \mathbb{F}((z^{-1}))$ , defined by

$$
w(z) = S(z)e^{z \cdot t} \mapsto w(z)|_{x=0, t=0} := S(z)|_{x=0, t=0}.
$$
 (2.29)

<span id="page-9-2"></span>**Lemma 2.13.** Let  $w(z) = S(z)e^{z \cdot t}$  be an oscillating function associated to a monic pseudodifferential operator  $S(\partial) = \partial^N + s_0 \partial^{N-1} + s_1 \partial^{N-2} + \cdots \in \mathcal{F}((\partial^{-1}))$ . If  $P(\partial) \in \mathcal{F}((\partial^{-1}))$  is a pseudodifferential operator such that  $P(\partial)w(z) = 0$ , then  $P(\partial) = 0.$ 

*Proof.* It is a consequence of the obvious fact that a monic pseudodifferential operator is not a zero divisor in  $\mathcal{F}((\partial^{-1}$ )).  $\Box$ 

We also define the space of *anti-oscillating functions* as  $\mathcal{F}((z^{-1}))e^{-z\cdot t}$ . We define the action of  $\mathcal{F}((\partial^{-1}))$  on the space of anti-oscillating functions  $\mathcal{F}((z^{-1}))e^{-z \cdot t}$  by

$$
P(\partial)(S(-z)e^{-z\cdot t}) = (P \circ S)(-z)e^{-z\cdot t}.
$$

Furthermore, we define the action of the partial derivatives  $\frac{\partial}{\partial t_k}$ ,  $k \in \mathcal{Z}$ , on the space of anti-oscillating functions  $\mathcal{F}((z^{-1}))e^{-z \cdot t}$  by

$$
\frac{\partial}{\partial t_k}(S(z)e^{-z\cdot t}) = \left(\frac{\partial S(z)}{\partial t_k} - z^k S(z)\right)e^{-z\cdot t}.
$$
\n(2.30)

Given an oscillating function  $w(z) = S(z)e^{z \cdot t}$ , associated to an invertible pseudodifferential operator  $S(\partial) \in \mathcal{F}((\partial^{-1}))$ , we define the corresponding *adjoint* antioscillating function as

<span id="page-10-6"></span>
$$
w^*(z) := (S^*)^{-1}(-z)e^{-z \cdot t}.
$$
\n(2.31)

Note that two oscillating functions, or two anti-oscillating functions, cannot be multiplied. On the other hand, we can multiply an oscillating function  $w(z)$  $P(z)e^{zt}$  and an anti-oscillating function  $\omega(z) = Q(-z)e^{-zt}$ , the result being a Laurent series:

<span id="page-10-1"></span>
$$
w(z)\omega(z) = P(z)Q(-z) \in \mathcal{F}((z^{-1})).
$$
\n(2.32)

It then also makes sense to take its residue  $\text{Res}_z(w(z)\omega(z)) \in \mathcal{F}$ , which is, by definition, the coefficient of  $z^{-1}$  of  $(2.32)$ .

## 3. Lax operators and hierarchies of Lax equations

<span id="page-10-0"></span>3.1. **Fractional powers.** As in Section [2,](#page-4-0) we let  $V$  be the algebra [\(2.5\)](#page-4-1) of differential polynomials in  $\ell$  variables  $u_{\alpha}, \alpha \in I$ . Moreover, throughout this section we fix a monic pseudodifferential operator  $L(\partial) \in \mathcal{V}((\partial^{-1}))$  of order  $N \geq 1$ :

<span id="page-10-2"></span>
$$
L(\partial) = \partial^N + a_0 \partial^{N-1} + a_1 \partial^{N-2} + a_2 \partial^{N-3} + \dots, \quad a_i \in \mathcal{V}.
$$
 (3.1)

Recalling the definition [\(2.26\)](#page-8-1) of the N-th root  $L^{\frac{1}{N}}(\partial)$ , we let, for  $k \in \mathbb{Z}$ ,

<span id="page-10-5"></span>
$$
B_k(\partial) = (L^{\frac{k}{N}}(\partial))_+ \in \mathcal{V}[\partial], \quad B_k^{(-)}(\partial) = (L^{\frac{k}{N}}(\partial))_- \in \mathcal{V}[[\partial^{-1}]]\partial^{-1}, \quad (3.2)
$$

so that

<span id="page-10-4"></span>
$$
L^{\frac{k}{N}}(\partial) = B_k(\partial) + B_k^{(-)}(\partial). \tag{3.3}
$$

Recall from [\(2.10\)](#page-5-4) the definition of the Frechet derivative  $D_{\alpha}^{(f)}(\partial) \in \mathcal{V}[\partial], \alpha \in I$ , associated to  $f \in V$ . It can be viewed as a differential operator

$$
D^{(f)}(\partial):\mathcal{V}^I\to\mathcal{V}\,,
$$

mapping  $P = (P_{\alpha})_{\alpha \in I} \mapsto \sum_{\alpha \in I} D_{\alpha}^{(f)}(\partial) P_{\alpha}$ . We generalize this map by replacing  $f \in V$  with an arbitrary pseudodifferential operator  $L(\partial) \in V((\partial^{-1}))$  as in  $(3.1)$ . As a result we obtain a map

$$
D^{(L)}: \mathcal{V}^I \to \mathcal{V}((\partial^{-1}))\,,
$$

defined by taking the Frechet derivative of the coefficients  $a_i \in V$  of  $L(\partial)$ . In other words, recalling [\(2.11\)](#page-5-5)

$$
D^{(L)}(P) = \sum_{i=0}^{\infty} X_P(a_i) \partial^{N-1-i}.
$$

We denote by  $D^{(L)}(\mathcal{V}^I)$  the image of this map:

$$
D^{(L)}(\mathcal{V}^I) = \{ D^{(L)}(P) | P \in \mathcal{V}^I \} \subset \mathcal{V}((\partial^{-1})).
$$

In the present section we review some results about scalar Lax operators, see e.g. [\[GD76,](#page-37-2) [DS85\]](#page-37-6). Let  $L(\partial)$  be as in  $(3.1)$ , let  $B(\partial) \in \mathcal{V}[\partial]$ , and consider the associated Lax equation

<span id="page-10-3"></span>
$$
\frac{\partial L(\partial)}{\partial t_B} = [B(\partial), L(\partial)].
$$
\n(3.4)

Recall that, if  $\frac{\partial u}{\partial t_B} = P \in \mathcal{V}^I$ , then, by the chain rule,  $L(\partial)$  evolves according to

$$
\frac{\partial L(\partial)}{\partial t_B} = \sum_{i=0}^{\infty} X_P(a_i) \partial^{N-1-i} = D^{(L)}(P) \in \mathcal{V}((\partial^{-1})).
$$

Hence, equation  $(3.4)$  can be rewritten as

<span id="page-11-3"></span>
$$
D^{(L)}\left(\frac{\partial u}{\partial t_B}\right) = [B(\partial), L(\partial)].
$$
\n(3.5)

We then obtain that equation  $(3.4)$  is consistent (cf. Definition [2.4\)](#page-6-4) if and only if

<span id="page-11-1"></span>
$$
[B(\partial), L(\partial)] = D^{(L)}(P)(=X_P(L(\partial))) \in \mathcal{V}((\partial^{-1})), \tag{3.6}
$$

for some  $P \in \mathcal{V}^I$ . With a slight abuse of notation, if  $(3.6)$  holds, we shall denote by  $\frac{\partial}{\partial t_B}$  the (rather an) associated evolutionary vector field  $X_P$ . The element  $P =$  $(P_{\alpha})_{\alpha \in I} \in \mathcal{V}^{I}$  is defined up to adding an element in the kernel of  $D^{(L)}$ .

We want to describe more explicitly the space of operators  $B(\partial)$  for which [\(3.6\)](#page-11-1) holds. Clearly, a necessary condition for  $(3.6)$  to hold is that ord $([B(\partial), L(\partial)])$  $\text{ord}(X_P(L)) \leq N-1.$ 

<span id="page-11-5"></span>**Lemma 3.1.** For  $B(\partial) \in V[\partial]$  we have  $\text{ord}([B(\partial), L(\partial)]) \leq N - 1$  if and only if  $B(\partial) \in \mathcal{V} + \text{Span}_{\mathbb{F}}\{B_k(\partial)\}_{k\geq 1} \subset \mathcal{V}[\partial].$ 

In fact,

<span id="page-11-6"></span>
$$
\operatorname{ord}([B_k(\partial), L(\partial)]) \le N - 2 \quad \text{for every } k \ge 1. \tag{3.7}
$$

*Proof.* If  $f \in V$ , by  $(2.23)$ ,  $[f, L(\partial)]$  has order at most  $N-1$ . Obviously,  $[L^{\frac{k}{N}}(\partial), L(\partial)] =$ 0 (since  $L(\partial) = (L^{\frac{1}{N}}(\partial))^N$ ). Hence, by [\(3.3\)](#page-10-4), we have

$$
[B_k(\partial), L(\partial)]=- [B_k^{(-)}(\partial), L(\partial)],
$$

which has order at most  $N-2$ . This proves the "if" part.

For the opposite implication, let  $B(\partial) \in V[\partial]$  be such that ord $([B(\partial), L(\partial)]) \leq$  $N-1$ . We shall prove, by induction on  $k = \text{ord}(B(\partial)) \geq 0$ , that  $B(\partial) \in \mathcal{V}$  +  $\text{Span}_{\mathbb{F}}\{B_h(\partial)\}_{h>1}$ . For  $k=0$ , the claim is obvious, so we can assume  $k\geq 1$ . Let  $a \in V$  be the leading coefficient of  $B(\partial)$ :

$$
B(\partial) = a\partial^k + \text{ lower order terms}.
$$

We have

$$
[B(\partial), L(\partial)] = -Na'\partial^{N+k-1} + \text{ lower order terms}.
$$

Hence,  $a' = 0$ , i.e.  $a \in \mathbb{F}$ . Since, by Lemma [2.11\(](#page-8-2)b),  $L^{\frac{1}{N}}(\partial)$  (hence  $B_k(\partial)$ ) is monic,

$$
ord(B(\partial) - aB_k(\partial)) \leq k - 1,
$$

and, by the "if" part,

$$
ord([B(\partial)-aB_k(\partial),L(\partial)])\leq N-1.
$$

The claim follows by the assumption of induction.  $\Box$ 

Recall from Section [2.2](#page-4-4) the notion of compatibility. We can study when two consistent Lax equations [\(3.4\)](#page-10-3) are compatible. This is discussed in the following proposition.

<span id="page-11-0"></span>**Proposition 3.2.** Let  $B(\partial), C(\partial) \in V[\partial]$  be two differential operators satisfying the consistency condition [\(3.6\)](#page-11-1). Then, assuming that  $B(\partial), C(\partial) \in \text{Span}_{\mathbb{F}}\{B_k(\partial)\}_{k>1}$ , the corresponding Lax equations

<span id="page-11-2"></span>
$$
\frac{\partial L(\partial)}{\partial t_B} = [B(\partial), L(\partial)] \quad and \quad \frac{\partial L(\partial)}{\partial t_C} = [C(\partial), L(\partial)] \tag{3.8}
$$

are compatible on the subalgebra of V generated by the coefficients of  $L(\partial)$ . In other words, the mixed second derivatives of  $L(\partial)$ , in virtue of  $(3.8)$ , coincide:

<span id="page-11-4"></span>
$$
\left(\frac{\partial^2 L(\partial)}{\partial t_B \partial t_C} = \right) \frac{\partial}{\partial t_B} \left[ C(\partial), L(\partial) \right] = \frac{\partial}{\partial t_C} \left[ B(\partial), L(\partial) \right].
$$
 (3.9)

*Proof.* (See [\[Dic03\]](#page-37-17).) Recall that the Lax equations  $(3.8)$  are equivalent to linear systems of quasi-evolution equations as in  $(3.5)$ . Moreover, consistency of  $(3.8)$ means that there exist  $P_B, P_C \in \mathcal{V}^I$  such that  $[B(\partial), L(\partial)] = X_{P_B}(L(\partial))$  and  $[C(\partial), L(\partial)] = X_{P_C}(L(\partial))$ . As before, by  $\frac{\partial}{\partial t_B}$  and  $\frac{\partial}{\partial t_C}$  we denote the evolutionary vector fields  $X_{P_B}$  and  $X_{P_C}$  respectively. The compatibility condition [\(3.9\)](#page-11-4) is then equivalent to

$$
[X_{P_B}, X_{P_C}](L(\partial)) = 0.
$$

Since  $\frac{\partial}{\partial t_B}$  and  $[B(\partial), \cdot]$  are both derivations of  $\mathcal{V}((\partial^{-1}))$  and  $\frac{\partial L(\partial)}{\partial t_B} = [B(\partial), L(\partial)],$ by Lemma  $2.12$  we have

<span id="page-12-0"></span>
$$
\frac{\partial L^{\frac{k}{N}}(\partial)}{\partial t_B} = \left[ B(\partial), L^{\frac{k}{N}}(\partial) \right],\tag{3.10}
$$

for every  $k \geq 0$ . By assumption,  $B(\partial) = \sum_{k \in \mathbb{Z}_{\geq 0}} \beta_k B_k(\partial), C(\partial) = \sum_{k \in \mathbb{Z}_{\geq 0}} \gamma_k B_k(\partial),$ with  $\beta_k, \gamma_k \in \mathbb{F}$ . Let us denote  $B^{(-)}(\partial) = \sum_k \beta_k B_k^{(-)}$  $k_k^{(-)}(\partial)$  and similarly for  $C^{(-)}(\partial)$ . Hence, by the Leibniz rule,

<span id="page-12-1"></span>
$$
[X_{P_B}, X_{P_C}](L(\partial)) = \frac{\partial}{\partial t_B} [C(\partial), L(\partial)] - \frac{\partial}{\partial t_C} [B(\partial), L(\partial)]
$$
  
\n
$$
= \sum_{k \in \mathbb{Z}_{\geq 0}} \left( \gamma_k \Big[ \left( \frac{\partial L^{\frac{k}{N}}(\partial)}{\partial t_B} \right)_+, L(\partial) \Big] - \beta_k \Big[ \left( \frac{\partial L^{\frac{k}{N}}(\partial)}{\partial t_C} \right)_+, L(\partial) \Big] \right)
$$
  
\n
$$
+ \Big[ C(\partial), \frac{\partial L(\partial)}{\partial t_B} \Big] - \Big[ B(\partial), \frac{\partial L}{\partial t_C} \Big]
$$
  
\n
$$
= \sum_{k \in \mathbb{Z}_{\geq 0}} \left( \gamma_k \Big[ [B(\partial), L^{\frac{k}{N}}(\partial)]_+, L(\partial) \Big] + \beta_k \Big[ [L^{\frac{k}{N}}(\partial), C(\partial)]_+, L(\partial) \Big] \right)
$$
  
\n
$$
+ \Big[ C(\partial), \Big[ B(\partial), L(\partial) \Big] \Big] - \Big[ B(\partial), \Big[ C(\partial), L(\partial) \Big] \Big]
$$
  
\n
$$
= \Big[ [B(\partial), C(\partial) + C^{(-)}(\partial)]_+, L(\partial) \Big] + \Big[ [B(\partial) + B^{(-)}(\partial), C(\partial)]_+, L(\partial) \Big]
$$
  
\n
$$
- \Big[ [B(\partial), C(\partial)], L(\partial) \Big].
$$
 (9)

In the second equality we used the obvious fact that applying  $\frac{\partial}{\partial t}$  to an element of  $\mathcal{V}((\partial^{-1}))$  commutes with taking its positive part, and in the third equality we used the Lax equations  $(3.8)$  and  $(3.10)$ . Note that

$$
\begin{aligned} &\left[B(\partial)+B^{(-)}(\partial),C(\partial)\right]_{+}-\left[B(\partial),C(\partial)\right] \\ &=\left[B(\partial)+B^{(-)}(\partial),C(\partial)\right]_{+}-\left[B(\partial),C(\partial)\right]_{+}\\ &=\left[B^{(-)}(\partial),C(\partial)\right]_{+}=\left[B^{(-)}(\partial),C(\partial)+C^{(-)}(\partial)\right]_{+}. \end{aligned}
$$

Hence, the RHS of [\(3.11\)](#page-12-1) becomes

$$
\begin{aligned}\n\left[ \left[ B(\partial) + B^{(-)}(\partial), C(\partial) + C^{(-)}(\partial) \right]_+, L(\partial) \right] \\
&= \sum_{h,k \in \mathbb{Z}_{\geq 0}} \beta_h \gamma_k \left[ \left[ L^{\frac{h}{N}}(\partial), L^{\frac{k}{N}}(\partial) \right]_+, L(\partial) \right] = 0.\n\end{aligned}
$$

The Lax equation [\(3.4\)](#page-10-3) is useful since it admits a large family of integrals of motion. For  $h \in V$  we denote  $\int h$  its image in the quotient space  $V/\partial V$ . Recall that an *integral of motion* for an evolution equation is an element  $\int h \in V/\partial V$  such that  $\frac{d}{dt} \int h = 0$ . The following fact is well known (see e.g. [\[Adl79,](#page-37-22) [Dic03\]](#page-37-17)).

Proposition 3.3. Assume that the Lax equation  $(3.4)$  is consistent, for a pseudodifferential operator  $L(\partial)$ , monic of order  $N \geq 1$ , and a differential operator  $B(\partial)$ . Then,  $\int \text{Res}_{\partial} L^{\frac{n}{N}}(\partial)$  is an integral of motion of [\(3.4\)](#page-10-3) for every  $n \geq 0$ .

*Proof.* It immediately follows from  $(3.10)$  and Lemma [2.10.](#page-8-0) □

3.2. Lax operators. In view of Lemma [3.1](#page-11-5) and Proposition [3.2](#page-11-0) we introduce the following:

**Definition 3.4.** A (scalar) Lax operator is a pseudodifferential operator  $L(\partial) \in$  $\mathcal{V}((\partial^{-1}))$  such that the Lax equation

<span id="page-13-0"></span>
$$
\frac{\partial L(\partial)}{\partial t_k} = [B_k(\partial), L(\partial)],\tag{3.12}
$$

is consistent and non-zero (i.e. the RHS is non-zero) for infinitely may values of  $k \geq 1$ . We denote by  $\mathcal{Z}_L \subset \mathbb{Z}_{\geq 1}$  the (infinite) set of k's such that equation [\(3.12\)](#page-13-0) is consistent.

As an immediate consequence of Lemma [3.1](#page-11-5) and Proposition [3.2](#page-11-0) we have the following result.

Corollary 3.5. For a scalar Lax operator  $L(\partial) \in \mathcal{V}((\partial^{-1}))$  all equations of the hierarchy [\(3.12\)](#page-13-0), with  $k \in \mathcal{Z}_L$ , are consistent and are compatible on the differential subalgebra of V generated by the coefficients of  $L(\partial)$ .

3.3. Examples of Lax operators. Let  $L(\partial)$  be a monic pseudodifferential oper-ator as in [\(3.1\)](#page-10-2). By [\(3.7\)](#page-11-6), the Lax equation [\(3.12\)](#page-13-0) implies  $\frac{\partial a_0}{\partial t_k} = 0$ . Hence, without loss of generality, we may assume that  $a_0 = 0$ . In this case  $B_1(\partial) = \partial$ , hence [\(3.12\)](#page-13-0) for  $k = 1$  is the trivial equation  $\frac{\partial a_j}{\partial t_1} = \partial a_j$ .

**Example 3.6.** Let  $V$  be the algebra of differential polynomials in infinitely many differential variables  $u_i, i \in \mathbb{Z}_{\geq 1}$ . Let  $L(\partial)$  be the Sato operator, i.e. the pseu-dodifferential operator [\(3.1\)](#page-10-2), with  $a_0 = 0$  and  $a_i = u_i$  for all  $i \ge 1$ . It follows from [\(3.7\)](#page-11-6) that the Lax equation [\(3.12\)](#page-13-0) is consistent for all  $k > 1$ . So, in this case  $\mathcal{Z}_L = \mathbb{Z}_{\geq 1}$  and all Lax equations are non-zero, hence  $L(\partial)$  is a Lax operator. The corresponding hierarchy of (compatible) Lax equations is called the KP hierarchy. Note that this hierarchy is isomorphic to that with  $N = 1$  [\[DSKV15\]](#page-37-7). As will be shown in Example [4.2,](#page-19-0) from the Lax equation one derives, in this case, the classical KP equation.

<span id="page-13-1"></span>**Example 3.7.** Let V be the algebra of differential polynomials in  $u_1, \ldots, u_{N-1}$ , where  $N \geq 2$ , and let  $L(\partial) = \partial^N + u_1 \partial^{N-2} + \cdots + u_{N-1}$ , which is called the Gelfand-Dickey operator. Again thanks to [\(3.7\)](#page-11-6), all Lax equations [\(3.12\)](#page-13-0) are consistent and they are nonzero if k is not divisible by N. Hence  $\mathcal{Z}_L = \mathbb{Z}_{\geq 1}$  and  $L(\partial)$  is a Lax operator. The corresponding hierarchy of evolution equation is called the N-th Gelfand-Dickey or N-th KdV hierarchy. For  $N = 2$ , V is an algebra of differential polynomials in a single variable  $u = u_1$ , the Lax operator is  $L(\partial) = \partial^2 + u$ , and the first non-trivial Lax equation is  $\frac{\partial L(\partial)}{\partial t_3} = [B_3(\partial), L(\partial)]$ , which is the KdV equation  $\frac{\partial u}{\partial t_3} = \frac{1}{4}u''' + \frac{3}{2}uu'$ . For  $N = 3$ ,  $\mathcal{V}$  is an algebra of differential polynomials in two variables  $u = u_1, v = u_2$  and  $L(\partial) = \partial^3 + u\partial + v$ . One can easily get  $B_2(\partial) = \partial^2 + \frac{2}{3}u$ . Hence, the Lax equation [\(3.12\)](#page-13-0) for  $k = 2$  becomes the Boussinesq equation

$$
\frac{\partial u}{\partial t_2} = 2v' - u'' \ , \ \frac{\partial v}{\partial t_2} = v'' - \frac{2}{3}u''' - \frac{2}{3}uu' \ .
$$

However, one can show that  $L(\partial) = \partial^3 + u$  is not a Lax operator.

<span id="page-13-2"></span>**Example 3.8.** Let  $V$  be the algebra of differential polynomials in the variables  $u_1, \ldots, u_{n-1}, p_1, q_1, \ldots, p_s, q_s$ , and consider the pseudodifferential operator

$$
L(\partial) = \partial^{n} + u_1 \partial^{n-2} + \dots + u_{n-1} + \sum_{i=1}^{s} p_i \partial^{-1} \circ q_i.
$$

This is a Lax operator with  $\mathcal{Z}_L = \mathbb{Z}_{\geq 1}$  [\[Che92,](#page-37-12) [KSS91,](#page-37-11) [KS92,](#page-37-13) [SS93\]](#page-37-14). The corresponding Lax hierarchy  $(3.12)$  is called the s-vector *n*-constrained KP hierarchy. It is obtained by a certain reduction of the KP hierarchy, and it follows from [\[DSKV16a,](#page-37-8) [CDSKVvdL20\]](#page-37-10) that it is compatible. The simplest special case  $L(\partial) = \partial + p\partial^{-1} \circ q$  gives the famous non-linear Schroedinger hierarchy, while  $L(\partial) = \partial^2 + u + p\partial^{-1} \circ q$  gives the Yajima-Oikawa hierarchy [\[Y076\]](#page-38-0).

<span id="page-14-0"></span>Example 3.9. A generalization of both Examples [3.7](#page-13-1) and [3.8](#page-13-2) is considered in [\[DSKV16a\]](#page-37-8). Namely, for each partition p of N we construct an  $r \times r$  matrix Lax operator  $L(\partial)$ , where r is the multiplicity of the largest part in p, for which  $\mathcal{Z}_L$  =  $\mathbb{Z}_{\geq 1}$ . The special case of the partition  $N = N$  corresponds to the N-th Gelfand Dickey hierarchy of Example [3.7,](#page-13-1) while the partition  $N = n + 1 + \cdots + 1$  (with  $s =$  $N-n$  ones) corresponds to the s-vector n-constrained KP hierarchy in Example [3.8.](#page-13-2) This is obtained by making use of the classical  $\mathcal{W}$ -algebra associated to the nilpotent element of  $\mathfrak{sl}_N$  with Jordan form corresponding to the partition p, generalizing the construction of [\[DS85\]](#page-37-6).

<span id="page-14-4"></span>Remark 3.10. The  $r \times r$  matrix Lax operator  $L(\partial)$  in Example [3.9](#page-14-0) is obtained as the quasideterminant (cf. [\[DSKV16a\]](#page-37-8))

$$
L(\partial) = (JA(\partial)^{-1}I)^{-1}
$$

where  $f^{p_1-1} = JI$  is the "canonical decomposition" of the  $p_1 - 1$ -th power of the nilpotent element f of  $\mathfrak{sl}_N$  with Jordan form corresponding to the partition p and  $A(\partial)$  is a certain first order matrix differential operator. For  $\mathfrak{g} = \mathfrak{so}_N$  or  $\mathfrak{sp}_N$  we have that  $A^*(\partial) = -A(\partial)$  (cf. [\[DSKV18\]](#page-37-9)). Hence, using Lemma 5.5 in [DSKV18] (with  $T = f^{p_1 - 1}$ ) we get that the Lax operator for the integrable hierarchy corresponding to the classical W-algebra associated to  $\mathfrak g$  and the partition  $p$  satisfies the further condition

$$
L^*(\partial) = (-1)^{p_1} L(\partial),
$$

where the adjoint is computed with respect to a bilinear form on an r-dimensional vector space which has parity  $\epsilon(-1)^{p_1-1}$ , with  $\epsilon=1$  for  $\mathfrak{so}_N$  and  $\epsilon=-1$  for  $\mathfrak{sp}_N$ . In this case we have  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0} + 1.$ 

<span id="page-14-1"></span>Remark 3.11. Let  $L(\partial)$  be a monic pseudodifferential operator of order N which is  $(-1)^N$ -adjoint, i.e. selfadjoint for even N and skewadjoint for odd N. Then by uniqueness of the N-th root of a monic operator,  $L^{\frac{1}{N}}(\partial)$  is necessarily a skewadjoint operator. Therefore,  $L^{\frac{k}{N}}(\partial)$  is  $(-1)^k$ -adjoint, hence  $B_k(\partial)$  is  $(-1)^k$ -adjoint as well. Therefore  $[B_k(\partial), L(\partial)]$  is  $(-1)^{N+k+1}$ -adjoint. And therefore the corresponding Lax equation  $(3.12)$  can be consistent only for odd k.

<span id="page-14-3"></span>**Example 3.12.** Let V be the algebra of differential polynomials in the variables  $v_j$ with  $j \ge 1$  odd, and let  $L(\partial)$  be the "generic" monic skewadjoint pseudodifferential operator of order 1, namely

<span id="page-14-2"></span>
$$
L(\partial) = \partial + \frac{1}{2} \sum_{j \ge 1, \text{ odd}} (\partial^{-j} \circ v_j + v_j \partial^{-j}). \tag{3.13}
$$

Then, by [\(3.7\)](#page-11-6),  $[B_k(\partial), L(\partial)]$  has negative order, and by Remark [3.11](#page-14-1) it is skewadjoint for every odd k. Hence, since  $L(\partial)$  is generic, the corresponding Lax equation [\(3.12\)](#page-13-0) is consistent and non-zero for every odd k, so that  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0} + 1$  and  $(3.13)$  is a Lax operator. The corresponding hierarchy of Lax equation is called the CKP hierarchy [\[DJKM81\]](#page-37-4). Note that as in the KP case we could have taken instead  $L(\partial)$  be a generic monic  $(-1)^N$ -adjoint pseudodifferential operator of order N, again obtaining the CKP hierarchy.

<span id="page-15-3"></span>Remark 3.13. Let  $L(\partial)$  be a monic pseudodifferential operator of order N satisfying

<span id="page-15-1"></span>
$$
L^*(\partial) = (-1)^N \partial \circ L(\partial) \partial^{-1}.
$$
\n(3.14)

We call such an operator  $(-1)^N$ -sesquiadjoint. Observe that any such operator can be written as

<span id="page-15-2"></span>
$$
L(\partial) = \partial^{N} + \frac{1}{2} \sum_{j \in \mathbb{Z}_{\geq 1}} (a_j \partial^{N-2j} + \partial^{N-2j-1} \circ a_j \partial), \qquad (3.15)
$$

with  $a_i \in V$  for every j. Observe also that, if N is odd, then

<span id="page-15-0"></span>
$$
\operatorname{Res}_{\partial} L(\partial)\partial^{-1} = 0. \tag{3.16}
$$

Indeed, since N is odd, neither  $a_j \partial^{N-2j}$  nor  $\partial^{N-2j-1} \circ a_j \partial$  can contribute to the residue [\(3.16\)](#page-15-0) for any  $j \in \mathbb{Z}_{\geq 0}$ . It is immediate to check, by the uniqueness of the N-th root of a monic operator, that  $L^{\frac{1}{N}}(\partial)$  is automatically  $(-1)$ -sesquiadjoint and, therefore,  $L^{\frac{k}{N}}(\partial)$  is  $(-1)^k$ -sesquiadjoint. As a consequence, for odd k,

$$
B_k(\partial)^* = ((L^{\frac{k}{N}}(\partial))_+)^* = ((L^{\frac{k}{N}}(\partial))^*)_+ = (-1)^k (\partial \circ L^{\frac{k}{N}}(\partial) \partial^{-1})_+ = -\partial \circ B_k(\partial) \partial^{-1}
$$

since, by [\(3.16\)](#page-15-0),  $L^{\frac{k}{N}}(\partial)$  does not have constant term. Hence,  $B_k(\partial)$  is  $(-1)^k$ sesquiadjoint as well. It follows that  $[B_k(\partial), L(\partial)]$  is  $(-1)^N$ -sesquiadjoint only for odd k, so the corresponding Lax equation  $(3.12)$  can be consistent only for odd k.

**Example 3.14.** Let V be the same as in Example [3.12](#page-14-3) and let  $L(\partial)$  be the "generic" monic pseudodifferential operator of order 1 which is  $(-1)$ -sesquiadjoint, as defined in  $(3.14)$ . Namely, by  $(3.15)$ ,

<span id="page-15-4"></span>
$$
L(\partial) = \partial + \frac{1}{2} \sum_{j \ge 1, \text{ odd}} (v_j \partial^{-j} + \partial^{-j-1} \circ v_j \partial). \tag{3.17}
$$

Then by [\(3.7\)](#page-11-6)  $[B_k(\partial), L(\partial)]$  has negative order and by Remark [3.13](#page-15-3) it is (-1)sesquiadjoint for every odd k. Hence, since  $L(\partial)$  is generic, the corresponding Lax equation [\(3.12\)](#page-13-0) is consistent and non-zero for every odd k, so that  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0} + 1$ and [\(3.17\)](#page-15-4) is a Lax operator. The corresponding hierarchy of Lax equations is called the BKP hierarchy [\[DJKM81\]](#page-37-4). As in Example [3.12](#page-14-3) we could have taken  $L(\partial)$  of order N, obtaining again the BKP hierarchy.

**Example 3.15.** Let  $L(\partial)$  be the "generic" monic  $(-1)^N$ -adjoint differential operator of order N, namely

<span id="page-15-5"></span>
$$
L(\partial) = \partial^{N} + \frac{1}{2} \sum_{j=1}^{\left[\frac{N}{2}\right]} (\partial^{N-2j} \circ v_j + v_j \partial^{N-2j}), \qquad (3.18)
$$

where  $\{v_j\}_{j=1}^{\lceil \frac{N}{2} \rceil}$  are the differential variables generating the algebra of differential polynomials V. Then, by [\(3.7\)](#page-11-6),  $[B_k(\partial), L(\partial)]$  is a differential operator of order bounded by  $N-2$ , and by Remark [3.11](#page-14-1) it is  $(-1)^N$ -adjoint for every odd k. Hence, since  $L(\partial)$  is generic, the corresponding Lax equation [\(3.12\)](#page-13-0) is consistent for every odd k and it is non-zero provided that k is not divisible by N, so that  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0}+1$ and [\(3.18\)](#page-15-5) is a Lax operator. The corresponding hierarchy of Lax equations is called the Drinfeld-Sokolov hierarchy associated to  $(A_N^{(2)})$  $\binom{2}{N-1}$ ,  $c_0$ ) for odd N and to  $\mathfrak{sp}_N$  for even N, [\[DS85\]](#page-37-6). The simplest case  $N = 2$  gives the KdV hierarchy, while  $N = 3$ gives the Kaup-Kupershmidt operator  $L(\partial) = \partial^3 + u\partial + \frac{1}{2}u'$ , and the Lax equation  $(3.12)$  with  $k = 5$  is the Kaup-Kupershmidt equation [\[Kaup80\]](#page-37-23).

**Example 3.16.** Let  $L(\partial)$  be the "generic" monic  $(-1)^N$ -sesquiadjoint differential operator of order N, namely

<span id="page-16-0"></span>
$$
L(\partial) = \partial^{N} + \frac{1}{2} \sum_{j=1}^{\left[\frac{N-1}{2}\right]} (v_j \partial^{N-2j} + \partial^{N-2j-1} \circ v_j \partial), \qquad (3.19)
$$

where  $\{v_j\}_{j=1}^{\lceil \frac{N-1}{2} \rceil}$  are the differential variables generating the algebra of differential polynomials V. Then, by  $(3.7), [B_k(\partial), L(\partial)]$  $(3.7), [B_k(\partial), L(\partial)]$  is a differential operator of order bounded by  $N-2$ , and by Remark [3.11](#page-14-1) it is  $(-1)^N$ -sesquiadjoint for every odd k. Hence, since L is generic, the corresponding Lax equation  $(3.12)$  is consistent for every odd  $k$  and it is non-zero provided that  $k$  is not divisible by  $N$ , so that  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0} + 1$  and [\(3.19\)](#page-16-0) is a Lax operator. The corresponding hierarchy of Lax equations is called the Drinfeld-Sokolov hierarchy associated to  $(A_N^{(2)})$  $\frac{c_1(z)}{N-1}, c_{\frac{N-1}{2}}$  for odd N and associated to  $(D_{\frac{N}{2}}^{(2)}, c_0)$  for even N. The simplest case  $N = 3$  gives  $L(\partial) = \partial^3 + u\partial$  and the Lax equation [\(3.12\)](#page-13-0) for  $k = 5$  is the Sawada-Kotera equation [\[SK74\]](#page-37-24).

<span id="page-16-1"></span>**Example 3.17.** For odd N and any  $s > 1$ , let

$$
L(\partial) = \partial^{N} + \frac{1}{2} \sum_{j=1}^{\frac{N-1}{2}} (\partial^{N-2j} \circ v_j + v_j \partial^{N-2j}) + \sum_{i=1}^{s} p_i \partial^{-1} \circ p_i,
$$

where  $v_1, \ldots, v_{N-1}, p_1, \ldots, p_s$  are the differential variables generating the algebra of differential polynomials V. For  $s = 1$  it is a Lax operator with  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0} + 1$ , since the corresponding hierarchy of Lax equations is the Drinfeld-Sokolov hierarchy associated to  $(A_N^{(2)}, c_{N-1})$ . The simplest case  $N = 2$  gives the operator  $L(\partial) =$  $\partial + u \partial^{-1} \circ u$  and the corresponding hierarchy is the modified KdV hierarchy. For arbitrary  $s \geq 1$  the pseudodifferential operator  $L(\partial)$  can be characterized as the "generic" skewadjoint pseudodifferential operator of the form given in Example  $3.8$ . Hence, by Remark [3.10,](#page-14-4) this is the Lax operator for the integrable hierarchy (after a suitable Dirac reduction) constructed in  $[DSKV18]$  for the classical affine W-algebra associated to  $\mathfrak{so}_{N+s}$  and a nilpotent element f with Jordan form corresponding to the partition  $p = (N, 1<sup>s</sup>)$ .

**Example 3.18.** For even N and even  $s > 1$ , let

$$
L(\partial) = \partial^N + \frac{1}{2} \sum_{j=1}^{\frac{N}{2}} (\partial^{N-2j} \circ v_j + v_j \partial^{N-2j}) + \sum_{i=1}^{\frac{s}{2}} (p_i \partial^{-1} \circ q_i - q_i \partial^{-1} \circ p_i),
$$

where  $v_1, \ldots, v_{\frac{N}{2}}, p_1, q_1, \ldots, p_{\frac{s}{2}}, q_{\frac{s}{2}}$  are the differential variables generating the algebra of differential polynomials  $\mathcal V$ . As in Example [3.17,](#page-16-1) by Remark [3.10,](#page-14-4) it follows that it is a Lax operator (with  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0} + 1$ ) for the integrable hierarchy (after a suitable Dirac reduction) constructed in  $[DSKV18]$  for the classical affine W-algebra associated to  $sp_{N+s}$  and a nilpotent element f whose Jordan form corresponds to the partition  $p = (N, 1<sup>s</sup>)$ . In fact,  $L(\partial)$  can be characterized as the "generic" selfadjoint pseudodifferential operator of the form given in Example  $3.8$ .

Example 3.19. Consider the "generic" selfadjoint differential operator of order 4:

<span id="page-16-3"></span>
$$
L(\partial) = \partial^4 + \partial \circ a\partial + b,\tag{3.20}
$$

over the algebra  $\mathcal V$  of differential polynomials in a and b. Its 4-th root has the form

<span id="page-16-2"></span>
$$
L^{\frac{1}{4}}(\partial) = \partial + \sum_{j=1}^{\infty} v_j \partial^{-j}, \qquad (3.21)
$$

where the coefficients  $v_j \in V$  satisfy a certain explicit recurrence relation, obtained from the identity  $(L^{\frac{1}{4}}(\partial))^4 = L(\partial)$ . In particular,

<span id="page-17-0"></span>
$$
v_1 = \frac{a}{4}, \ v_2 = -\frac{a'}{8}.
$$
 (3.22)

We want to write explicitly the Lax equation [\(3.12\)](#page-13-0) associated to  $L(\partial)$  for  $k = 3$ . By  $(3.21)$  and  $(3.22)$  we get

$$
B_3(\partial) = (L^{\frac{3}{4}}(\partial))_+ = \partial^3 + 3v_1\partial + 3(v_2 + v_1') = \partial^3 + \frac{3}{4}a\partial + \frac{3}{8}a'.
$$

Hence, by direct computation,

$$
[B_3(\partial), L(\partial)] = \partial \circ P\partial + Q,
$$

where

$$
P = 3b' - \frac{5}{4}a''' - \frac{3}{4}aa', \quad Q = b''' - \frac{3}{8}a^{(5)} - \frac{3}{8}(aa''' + a'a'') + \frac{3}{4}ab'.
$$

As a consequence, the Lax equation [\(3.12\)](#page-13-0) for  $k = 3$  and the operator [\(3.20\)](#page-16-3) is the following system of evolution equations on  $a$  and  $b$ :

<span id="page-17-1"></span>
$$
\frac{\partial a}{\partial t} = 3b' - \frac{5}{4}a''' - \frac{3}{4}aa' , \quad \frac{\partial b}{\partial t} = b''' - \frac{3}{8}a^{(5)} - \frac{3}{8}(aa''' + a'a'') + \frac{3}{4}ab' .
$$
 (3.23)

Next, we notice that the system  $(3.23)$  remains consistent once we impose the constraint

<span id="page-17-2"></span>
$$
b = \frac{1}{2}a'' + \frac{1}{4}a^2.
$$
 (3.24)

Indeed, as one can easily check,

$$
\frac{\partial b}{\partial t} = Q = \frac{1}{2}P'' + \frac{1}{2}aP = \frac{\partial}{\partial t} \left(\frac{1}{2}a'' + \frac{1}{4}a^2\right),\,
$$

provided that  $(3.24)$  holds. As a consequence, the system  $(3.23)$  can be greatly simplified with the change of variables  $(a, b) \rightarrow (a, u)$ , where

<span id="page-17-5"></span>
$$
b = \frac{1}{2}a'' + \frac{1}{4}a^2 - \frac{c}{3}u,
$$
\n(3.25)

and c is an arbitrary non-zero constant. Indeed, equations  $(3.23)$  can be rewritten in terms of these new variables as

<span id="page-17-3"></span>
$$
\frac{\partial u}{\partial t} = -\frac{1}{2}u''' - \frac{3}{4}au', \quad \frac{\partial a}{\partial t} = \frac{1}{4}a''' + \frac{3}{4}aa' - cu'.
$$
 (3.26)

Again we can ask under which constraint relating a and u the system  $(3.26)$  remains consistent. A solution to this question is given by the following constraint (see [\[Sok84\]](#page-38-2)):

<span id="page-17-4"></span>
$$
a = -\frac{u'''}{u'} + \frac{1}{2} \frac{(u'')^2}{(u')^2} - \frac{1}{3} \frac{h(u)}{(u')^2},
$$
\n(3.27)

where h is an arbitrary polynomial of degree 3 in  $u$  and the leading coefficient  $c$ . In this case the second equation of  $(3.26)$  remains consistent with the first, which reduces to the famous Krichever-Novikov equation [\[KN80\]](#page-37-25)

<span id="page-17-6"></span>
$$
4\frac{\partial u}{\partial t} = u''' - \frac{3}{2}\frac{(u'')^2}{u'} + \frac{h(u)}{u'}.
$$
 (3.28)

Here consistency means that applying  $D_a(\partial)$ , for a as in [\(3.27\)](#page-17-4), to the RHS of the first equation in  $(3.26)$ , we obtain the RHS of the second equation.

Consider the differential operator  $L(\partial)$  obtained from  $(3.20)$  by substituting a and b as  $(3.25)$  and  $(3.27)$ , obtaining an operator  $L(\partial)$  of order 4 discovered by Sokolov [\[Sok84\]](#page-38-2), who showed that [\(3.28\)](#page-17-6) is equivalent to the Lax equation  $\frac{\partial L(\partial)}{\partial t_3}$  =  $[B_3(\partial), L(\partial)]$ . One may expect that all equations  $\frac{\partial L(\partial)}{\partial t_k} = [B_k(\partial), L(\partial)]$  for odd k are consistent, but this is still an open problem.

<span id="page-18-0"></span>Let  $L(\partial) \in \mathcal{V}((\partial^{-1}))$  be as in [\(3.1\)](#page-10-2). In the terminology of Section [2.2](#page-4-4) we can talk about the solutions of the hierarchy [\(3.12\)](#page-13-0), as a collection of functions on space-time  $\varphi_{\alpha} \in \mathcal{F}$ ,  $\alpha \in I$ , such that

<span id="page-18-1"></span>
$$
\frac{\partial L(\varphi;\partial)}{\partial t_k} = [B_k(\varphi;\partial), L(\varphi;\partial)] \quad k \in \mathcal{Z}_L \,, \tag{4.1}
$$

where  $L(\varphi;\partial)$  and  $B_k(\varphi;\partial)$  are obtained applying the evaluation map  $\mathcal{V}((\partial^{-1})) \to$  $\mathcal{F}((\partial^{-1}))$  defined in [\(2.25\)](#page-8-4) to  $L(\partial)$  and  $B_k(\partial)$  respectively. In [\(4.1\)](#page-18-1) we assume that the functions on space-time  $\varphi \in \mathcal{F}$  depend on time variables  $t_k, k \in \mathcal{Z}_L$ .

**Proposition 4.1.** (a) Let  $L(\partial)$  be a Lax operator, let  $h, k \in \mathcal{Z}_L$ , and let  $\frac{\partial}{\partial t_k}, \frac{\partial}{\partial t_k}$ :  $V \rightarrow V$  denote (as usual) the associated evolutionary vector fields. Then the following Zakharov-Shabat equations  $(=$  zero-curvature equations) hold

<span id="page-18-2"></span>
$$
\frac{\partial B_k(\partial)}{\partial t_h} - \frac{\partial B_h(\partial)}{\partial t_k} + [B_k(\partial), B_h(\partial)] = 0, \ h, k \in \mathcal{Z}_L,
$$
\n(4.2)

as well as the following complementary Zakharov-Shabat equations

<span id="page-18-3"></span>
$$
\frac{\partial B_k^{(-)}(\partial)}{\partial t_h} - \frac{\partial B_h^{(-)}(\partial)}{\partial t_k} - [B_k^{(-)}(\partial), B_h^{(-)}(\partial)] = 0, \ h, k \in \mathcal{Z}_L,
$$
 (4.3)

where  $B_k^{(-)}$  $\binom{n}{k}$  [\(3.2\)](#page-10-5).

- $(b)$  Conversely, suppose that there exist infinitely many non-zero evolutionary vector fields  $\frac{\partial}{\partial t_k} : \mathcal{V} \to \mathcal{V}$ ,  $k \in \tilde{Z}_L$ , such that the Zakharov-Shabat equation [\(4.2\)](#page-18-2) holds for all  $h, k \in \tilde{Z}_L$ . Then, the Lax equation [\(3.12\)](#page-13-0) holds as well, for every  $k \in \mathcal{Z}_L$ . In particular, L is a Lax operator and  $\mathcal{Z}_L \subset \mathcal{Z}_L$ .
- (c) Suppose  $L(\partial)$  is a Lax operator. A collection of functions  $\varphi_{\alpha} \in \mathcal{F}$ ,  $\alpha \in I$ , is a solution of the hierarchy of Lax equations  $(3.12)$  if and only if it is a solution of the Zakharov-Shabat equations [\(4.2\)](#page-18-2) for every  $h, k \in \mathcal{Z}_L$ . Furthermore, if  $\varphi_{\alpha} \in \mathcal{F}, \ \alpha \in I, \text{ is a solution of either (3.12) or (4.2), then it is a solution of}$  $\varphi_{\alpha} \in \mathcal{F}, \ \alpha \in I, \text{ is a solution of either (3.12) or (4.2), then it is a solution of}$  $\varphi_{\alpha} \in \mathcal{F}, \ \alpha \in I, \text{ is a solution of either (3.12) or (4.2), then it is a solution of}$  $\varphi_{\alpha} \in \mathcal{F}, \ \alpha \in I, \text{ is a solution of either (3.12) or (4.2), then it is a solution of}$  $\varphi_{\alpha} \in \mathcal{F}, \ \alpha \in I, \text{ is a solution of either (3.12) or (4.2), then it is a solution of}$ the complementary Zakharov-Shabat equations [\(4.3\)](#page-18-3).

*Proof.* (See [\[Shi86\]](#page-37-16)) Let  $N \ge 1$  be the order of  $L(\partial)$ . If  $h, k \in \mathcal{Z}_L$ , then by [\(3.10\)](#page-12-0) we have

$$
\frac{\partial B_k(\partial)}{\partial t_h} + \frac{\partial B_k^{(-)}(\partial)}{\partial t_h} = \frac{\partial L^{\frac{k}{N}}(\partial)}{\partial t_h} = \left[ B_h(\partial), L^{\frac{k}{N}}(\partial) \right] = \left[ B_h(\partial), B_k(\partial) + B_k^{(-)}(\partial) \right].
$$
\n(4.4)

As a consequence, we get

<span id="page-18-4"></span>
$$
\begin{split}\n&\left(\frac{\partial B_k(\partial)}{\partial t_h} - \frac{\partial B_h(\partial)}{\partial t_k} + \left[B_k(\partial), B_h(\partial)\right]\right) \\
&+ \left(\frac{\partial B_k^{(-)}(\partial)}{\partial t_h} - \frac{\partial B_h^{(-)}(\partial)}{\partial t_k} - \left[B_k^{(-)}(\partial), B_h^{(-)}(\partial)\right]\right) \\
&= \left[B_h(\partial), B_k(\partial) + B_k^{(-)}(\partial)\right] - \left[B_k(\partial), B_h(\partial) + B_h(\partial)^{(-)}\right] \\
&+ \left[B_k(\partial), B_h(\partial)\right] - \left[B_k^{(-)}(\partial), B_h^{(-)}(\partial)\right] \\
&= -\left[B_k(\partial) + B_k^{(-)}(\partial), B_h(\partial) + B_h^{(-)}(\partial)\right] = -\left[L^{\frac{h}{N}}(\partial), L^{\frac{k}{N}}(\partial)\right] = 0.\n\end{split} \tag{4.5}
$$

Note that the two summands in the LHS of  $(4.5)$  are respectively in  $\mathcal{V}[\partial]$  and  $\mathcal{V}[[\partial^{-1}]]\partial^{-1}$ . Hence, they must both vanish, proving part (a).

Next, let us prove part (b). For  $k, h \in \tilde{Z}_L$ , we have, by [\(4.2\)](#page-18-2),

<span id="page-19-1"></span>
$$
\frac{\partial L^{\frac{k}{N}}(\partial)}{\partial t_{h}} - [B_{h}(\partial), L^{\frac{k}{N}}(\partial)] = \frac{\partial B_{k}(\partial)}{\partial t_{h}} + \frac{\partial B_{k}^{(-)}(\partial)}{\partial t_{h}} - [B_{h}(\partial), B_{k}(\partial) + B_{k}^{(-)}(\partial)]
$$
  
\n
$$
= \frac{\partial B_{k}(\partial)}{\partial t_{h}} - [B_{h}(\partial), B_{k}(\partial)] + \frac{\partial B_{k}^{(-)}(\partial)}{\partial t_{h}} - [B_{h}(\partial), B_{k}^{(-)}(\partial)]
$$
  
\n
$$
= \frac{\partial B_{h}(\partial)}{\partial t_{k}} + \frac{\partial B_{k}^{(-)}(\partial)}{\partial t_{h}} - [B_{h}(\partial), B_{k}^{(-)}(\partial)].
$$
\n(4.6)

Since  $B_k^{(-)}$  $\binom{n}{k}(\partial)$  has negative order, the RHS of  $(4.6)$  is a pseudodifferential operator of order less than or equal to  $h-1$  (independently of k). Assume, by contradiction, that

$$
R(\partial) = \frac{\partial L^{\frac{1}{N}}(\partial)}{\partial t_h} - [B_h(\partial), L^{\frac{1}{N}}(\partial)]
$$

is a pseudodifferential operator of order  $d \in \mathbb{Z}$  and leading coefficient  $v \in \mathcal{V}\backslash\{0\}$ . By the Leibniz rule, we have

<span id="page-19-2"></span>
$$
\frac{\partial L^{\frac{k}{N}}(\partial)}{\partial t_h} - \left[ B_h(\partial), L^{\frac{k}{N}}(\partial) \right] = \sum_{i=0}^{k-1} L^{\frac{i}{N}}(\partial) R(\partial) L^{\frac{k-1-i}{N}}(\partial). \tag{4.7}
$$

The RHS of [\(4.7\)](#page-19-2) has order equal to  $d+k-1$ , the leading coefficient being equal to  $kv \neq 0$ . Since the LHS of [\(4.6\)](#page-19-1) and [\(4.7\)](#page-19-2) coincide, we deduce that  $d + (k-1)N \leq$  $h-1$  for every  $k \in \mathcal{Z}_L$ , which is impossible since  $\mathcal{Z}_L$  is infinite. Hence, R must be 0. It follows, setting  $k = N$  in [\(4.7\)](#page-19-2), that the Lax equation [\(3.12\)](#page-13-0) holds for every  $k \in \tilde{Z}_L$ , so that  $\tilde{Z}_L \subset \mathcal{Z}_L$ , as claimed.

The proof of part (c) is the same as for parts (a) and (b), by replacing everywhere  $L(\partial), B_k(\partial)$  and  $B_h(\partial)$  with their evaluations  $L(\varphi; \partial), B_k(\varphi; \partial)$  and  $B_h(\varphi; \partial)$  in  $\mathcal{F}((\partial^{-1}% (\theta),\theta))=\mathcal{F}_{T_{\mathcal{A},\mathcal{C}}^{(n)}}(\theta),$ )). □

<span id="page-19-0"></span>**Example 4.2.** Let  $L(\partial) = \partial + \sum_{i=1}^{\infty} u_i \partial^{-i}$  be the Sato Lax operator. By a straightforward computation we have  $B_1(\partial) = \partial$ ,  $B_2(\partial) = \partial^2 + 2u_1$  and  $B_3(\partial) =$  $\partial^3 + 3u_1\partial + 3u_2 + 3u'_1$ . The Zakharov-Shabat equation [\(4.2\)](#page-18-2) for  $k = 2$  and  $h = 3$ gives

$$
\left(6u_2' + 3u_1'' - 3\frac{\partial u_1}{\partial t_2}\right)\partial + 2\frac{\partial u_1}{\partial t_3} - 3\frac{\partial u_2}{\partial t_2} - 3\frac{\partial u_1'}{\partial t_2} - 6u_1u_1' + 3u_2'' + u_1'''.
$$
 (4.8)

This is equivalent to

$$
\frac{\partial u_1}{\partial t_2} = 2u_2' + u_1'', \qquad 2\frac{\partial u_1}{\partial t_3} - 3\frac{\partial u_2}{\partial t_2} = 6u_1u_1' + 3u_2'' + 2u_1''' = 0.
$$

Applying  $\frac{\partial}{\partial t_2}$  to the first equation and  $\partial$  to the second we get

$$
\frac{\partial^2 u_1}{\partial t_2^2} = 2\frac{\partial u_2'}{\partial t_2} + \frac{\partial u_1''}{\partial t_2} = 2\frac{\partial u_2'}{\partial t_2} + 2u_2''' + u_1^{(4)}
$$

and

$$
2\frac{\partial u_1'}{\partial t_3} = 3\frac{\partial u_2'}{\partial t_2} + 6(u_1\partial u_1)' + 3u_2''' + 2u_1^{(4)},
$$

from which follows that

$$
3\frac{\partial^2 u_1}{\partial t_2^2} - 4\partial\frac{\partial u_1}{\partial t_3} = -\partial^2 \left(6u_1^2 + \partial^2 u_1\right) \, .
$$

Renaming  $t_2 = y$ ,  $t_3 = t$ , and  $u = 2u_1$ , we get

$$
3\frac{\partial^2 u}{\partial y^2} = \left(4\frac{\partial u}{\partial t} - u''' - 6uu'\right)'.
$$

which is the Kadomtsev-Petviashvili (KP) equation.

## <span id="page-20-0"></span>5. Lax and Sato equations, linear problem and bilinear equation

Let  $L(\partial) \in \mathcal{V}((\partial^{-1}))$  be as in [\(3.1\)](#page-10-2). Let, as in Section [2.1,](#page-4-5) F be an algebra of functions in space-time, with time derivatives  $\frac{\partial}{\partial t_k}$  indexed by  $k \in \mathcal{Z}_L$ .

# <span id="page-20-7"></span>5.1. Lax equations and Sato equations.

<span id="page-20-1"></span>**Theorem 5.1.** Let  $\varphi = (\varphi_\alpha)_{\alpha \in I}$  be a collection of functions  $\varphi_\alpha \in \mathcal{F}$ .

(a) If there exists a monic pseudodifferential operator of order  $0$  over the algebra of functions on space-time  $\mathcal{F},$ 

<span id="page-20-2"></span>
$$
S(\partial) = 1 + \sum_{i=1}^{\infty} s_i \partial^{-i} \in \mathcal{F}((\partial^{-1})),
$$
\n(5.1)

satisfying the following dressing equation

<span id="page-20-3"></span>
$$
L(\varphi; \partial) = S(\partial) \partial^{N} S^{-1}(\partial), \qquad (5.2)
$$

and Sato equations

<span id="page-20-4"></span>
$$
\frac{\partial S(\partial)}{\partial t_k} = -B_k^{(-)}(\varphi; \partial)S(\partial), \quad \text{for every } k \in \mathcal{Z}_L \,, \tag{5.3}
$$

then  $\varphi$  is a solution of the hierarchy of Lax equations [\(3.12\)](#page-13-0).

(b) Conversely, assume that the algebra of functions  $\mathcal F$  is integrable (cf. Definition [2.1\)](#page-4-6) and assume that  $\varphi$  is a solution of the hierarchy of Lax equations [\(3.12\)](#page-13-0) such that  $a_0(\varphi) = 0$   $(a_0(\varphi) \in \mathcal{F}$  being the evaluation at  $u = \varphi$ ). Then, for every monic constant coefficients pseudodifferential operator  $S_0(\partial) \in \mathbb{F}((\partial^{-1}))$ of order  $0$ , there exists a unique order  $0$  monic pseudodifferential operator  $S(\partial) \in \mathcal{F}((\partial^{-1}))$  as in [\(5.1\)](#page-20-2) satisfying the dressing equation [\(5.2\)](#page-20-3) and the Sato equations [\(5.3\)](#page-20-4), and such that  $S(\partial)|_{x=t=0} = S_0(\partial)$ .

*Proof.* (See [\[Shi86\]](#page-37-16)) Applying  $\frac{\partial}{\partial t_k}$ :  $\mathcal{F} \to \mathcal{F}$  to both sides of the dressing equation  $(5.2)$ , and using the Sato equations  $(5.3)$ , we get

$$
\frac{\partial L(\varphi;\partial)}{\partial t_k} = \frac{\partial S(\partial)}{\partial t_k} \partial^N S^{-1}(\partial) - S(\partial) \partial^N S^{-1}(\partial) \frac{\partial S(\partial)}{\partial t_k} S^{-1}(\partial) \n= -B_k^{(-)}(\varphi;\partial)S(\partial) \partial^N S^{-1}(\partial) + S(\partial) \partial^N S^{-1}(\partial) B_k^{(-)}(\varphi;\partial) \n= -[B_k^{(-)}(\varphi;\partial), L(\varphi;\partial)] = [B_k(\varphi;\partial), L(\varphi;\partial)].
$$

This proves (a).

In order to prove part (b), assume that  $\varphi_{\alpha} \in \mathcal{F}$ ,  $\alpha \in I$ , is a solution of the Lax equations  $(3.12)$ . The dressing equation  $(5.2)$  can be rewritten as

<span id="page-20-5"></span>
$$
(1 + s_1 \partial^{-1} + s_2 \partial^{-2} + ...) \partial^N
$$
  
=  $(\partial^N + a_0(\varphi) \partial^{N-1} + a_1(\varphi) \partial^{N-2} + ...)(1 + s_1 \partial^{-1} + s_2 \partial^{-2} + ...).$  (5.4)

By looking at the coefficient of  $\partial^{N-1}$  in both sides of equation [\(5.4\)](#page-20-5), we get  $a_0(\varphi)$  = 0, while, by looking at the coefficient of  $\partial^{N-1-i}$ , for  $i \geq 1$ , in both sides of [\(5.4\)](#page-20-5), we get

<span id="page-20-6"></span>
$$
s'_{i} = -\frac{1}{N} a_{i}(\varphi) - \frac{1}{N} \sum_{n=0}^{i-2} {N \choose n+2} s_{i-n-1}^{(n)} - \frac{1}{N} \sum_{j=1}^{i-1} \sum_{n=0}^{i-j-1} {N-1-j \choose n} a_{j}(\varphi) s_{i-j-n}^{(n)}.
$$
\n(5.5)

Next, the Sato equations  $(5.3)$  can be rewritten as

<span id="page-21-1"></span>
$$
\frac{\partial s_1}{\partial t_k} \partial^{-1} + \frac{\partial s_2}{\partial t_k} \partial^{-2} + \dots
$$
  
=  $(c_{1,k}(\varphi)\partial^{-1} + c_{2,k}(\varphi)\partial^{-2} + \dots)(1 + s_1\partial^{-1} + s_2\partial^{-2} + \dots),$  (5.6)

where we are letting  $-B_k^{(-)}$  $\mathcal{L}_{k}^{(-)}(\partial) = \sum_{j=1}^{\infty} c_{j,k} \partial^{-j} \in \mathcal{V}[[\partial^{-1}]]\partial^{-1}$ , and we denote, as usual, by  $c_{j,k}(\varphi) \in \mathcal{F}$  the evaluation of  $c_{j,k} \in \mathcal{V}$  at  $u = \varphi$ . By looking at the coefficient of  $\partial^{-i}$ , for  $i \geq 1$ , in both sides of equation [\(5.6\)](#page-21-1), we get

<span id="page-21-2"></span>
$$
\frac{\partial s_i}{\partial t_k} = c_{i,k}(\varphi) + \sum_{j=1}^{i-1} \sum_{n=0}^{i-j-1} (-1)^n {j+n-1 \choose n} c_{j,k}(\varphi) s_{i-j-n}^{(n)}, \quad k \in \mathcal{Z}_L. \tag{5.7}
$$

For every  $i \geq 1$ , equations [\(5.5\)](#page-20-6) and [\(5.7\)](#page-21-2) form a system of equations in the only unknown function  $s_i \in \mathcal{F}$ , if we assume that we have solved, recursively, the previous equations on  $s_j$  with  $j < i$ . Indeed, the RHS of both [\(5.5\)](#page-20-6) and [\(5.7\)](#page-21-2) only involves the functions  $a_j(\varphi) \in \mathcal{F}$ , which are given, and functions  $s_j^{(n)}$  with  $j < i$ , which are assumed to be known by the recursive construction. By the integrability assumption on F (cf. Definition [2.1\)](#page-4-6), this system has a unique solution  $s_i \in \mathcal{F}$  satisfying the initial condition  $s_i(0) = s_{i,0}$ , the coefficient of  $\partial^{-i}$  in  $S_0(\partial)$ , provided that the compatibility conditions  $(2.4)$  hold. Recalling that  $(5.5)$  is equivalent to the dressing equation  $S(\partial)\partial^{N}S^{-1}(\partial) = L(\varphi;\partial)$ , the first of the compatibility conditions [\(2.4\)](#page-4-7) for the system  $(5.5)-(5.7)$  $(5.5)-(5.7)$  $(5.5)-(5.7)$  amounts to

$$
\frac{\partial L(\varphi;\partial)}{\partial t_k} = -[B_k^{(-)}(\varphi;\partial), L(\varphi;\partial)],
$$

which holds since, by assumption,  $\varphi$  is a solution of the Lax equations [\(3.12\)](#page-13-0). On the other hand, [\(5.7\)](#page-21-2) is equivalent to the Sato equation  $\frac{\partial S(\partial)}{\partial t_k} = -B_k^{(-)}$  $S_k^{(-)}(\varphi;\partial)S(\partial),$ hence, the second of the compatibility conditions  $(2.4)$  for the system  $(5.5)-(5.7)$  $(5.5)-(5.7)$  $(5.5)-(5.7)$ amounts to

$$
-\frac{\partial B_k^{(-)}(\varphi;\partial)}{\partial t_h}S(\partial)+B_k^{(-)}(\varphi;\partial)B_h^{(-)}(\varphi;\partial)S(\partial)\\ =-\frac{\partial B_h^{(-)}(\varphi;\partial)}{\partial t_k}S(\partial)+B_h^{(-)}(\varphi;\partial)B_k^{(-)}(\varphi;\partial)S(\partial)\,,
$$

which is the complementary Zakharov-Shabat equation  $(4.3)$ , and therefore holds by Proposition  $4.1(c)$  $4.1(c)$ .

<span id="page-21-4"></span>Remark 5.2. Note that, by the uniqueness of the  $N$ -th root of a monic pseu-dodifferential operator (cf. Lemma [2.11\(](#page-8-2)b)), the dressing equation  $L(\varphi; \partial)$  =  $S(\partial)\partial^{N}S^{-1}(\partial)$  on the pseudodifferential operator  $S(\partial) \in \mathcal{F}((\partial^{-1}))$ , monic of order 0, is equivalent to the equation

<span id="page-21-5"></span>
$$
L^{\frac{1}{N}}(\varphi;\partial) = S(\partial)\partial S^{-1}(\partial).
$$
\n(5.8)

5.2. Sato equation and the linear problem for the wave function. Consider the space of oscillating functions  $\mathcal{F}((z^{-1}))e^{z \cdot t}$  defined in Section [2.6.](#page-9-1)

<span id="page-21-0"></span>**Theorem 5.3.** Let  $\varphi = (\varphi_{\alpha})_{\alpha \in I}$ ,  $\varphi_{\alpha} \in \mathcal{F}$ . Let  $S(\partial) \in \mathcal{F}((\partial^{-1}))$  be a monic pseudodifferential operator of order 0, and let  $w(z) = S(z)e^{z \cdot t} \in \mathcal{F}((z^{-1}))e^{z \cdot t}$  be the corresponding oscillating function.

(a) The dressing equation [\(5.2\)](#page-20-3) on  $S(\partial)$  is equivalent to the following eigenvalue problem *on the* wave function  $w(z)$ :

<span id="page-21-3"></span>
$$
L(\varphi; \partial)w(z) = z^N w(z). \tag{5.9}
$$

(b) Assuming that the dressing equation [\(5.2\)](#page-20-3) holds, the Sato equation [\(5.3\)](#page-20-4) for  $S(\partial)$  is equivalent to the following linear problem for the wave function  $w(z)$ :

<span id="page-22-2"></span>
$$
\frac{\partial w(z)}{\partial t_k} = B_k(\varphi; \partial) w(z), \quad k \in \mathcal{Z}_L.
$$
 (5.10)

*Proof.* Equation  $(5.9)$  can be rewritten, using the first equation in  $(2.27)$ , as

$$
(L(\varphi; \partial)S(\partial) - S(\partial)\partial^{N})e^{z \cdot t} = 0.
$$
\n(5.11)

Claim (a) follows by Lemma [2.13.](#page-9-2) Next, let us prove claim (b). Note that, by Remark [5.2,](#page-21-4) the dressing equation [\(5.2\)](#page-20-3) is also equivalent to

$$
L^{\frac{1}{N}}(\varphi;\partial)w(z) = zw(z).
$$

Hence we have

<span id="page-22-1"></span>
$$
z^k w(z) = L^{\frac{k}{N}}(\varphi; \partial)w(z) = \left(B_k(\varphi; \partial) + B_k^{(-)}(\varphi; \partial)\right)w(z).
$$
 (5.12)

By  $(2.28)$  and  $(5.12)$ , we get

$$
\frac{\partial w(z)}{\partial t_k} = \left(\frac{\partial S(\partial)}{\partial t_k} + z^k S(z)\right) e^{z \cdot t} = \left(\frac{\partial S(\partial)}{\partial t_k} S^{-1}(\partial) + z^k\right) w(z) \n= \left(\frac{\partial S(\partial)}{\partial t_k} S^{-1}(\partial) + B_k(\varphi; \partial) + B_k^{(-)}(\varphi; \partial)\right) w(z).
$$

Hence, if Sato equation [\(5.3\)](#page-20-4) holds, we automatically get the linear problem [\(5.10\)](#page-22-2). Conversely, if the linear problem [\(5.10\)](#page-22-2) holds, Sato equation [\(5.3\)](#page-20-4) holds due to Lemma [2.13.](#page-9-2)  $\Box$ 

### <span id="page-22-4"></span>5.3. Lax equations and the bilinear equation on the wave function.

<span id="page-22-0"></span>**Theorem 5.4.** Let  $S(\partial) \in \mathcal{F}((\partial^{-1}))$  be an invertible pseudodifferential operator, let  $w(z) = S(z)e^{z \cdot t} \in \mathcal{F}((z^{-1}))e^{z \cdot t}$  be the corresponding oscillating function, and let  $w^*(z) = (S^*)^{-1}(-z)e^{-z \cdot t}$  be the corresponding adjoint anti-oscillating function  $(cf. (2.31))$  $(cf. (2.31))$  $(cf. (2.31))$ . Then, the following conditions are equivalent:

(i) For every  $k \in \mathcal{Z}_L$ , we have

$$
\left(\frac{\partial S(\partial)}{\partial t_k}S^{-1}(\partial)+S(\partial)\partial^k S^{-1}(\partial)\right)_-=0.
$$

(ii) For every  $k \in \mathcal{Z}_L$ , there exists a differential operator  $\tilde{B}_k(\partial) \in \mathcal{F}[\partial]$  such that

$$
\frac{\partial w(z)}{\partial t_k} = \tilde{B}_k(\partial)w(z) .
$$

(iii) For every  $s \geq 0$ ,  $k_1, \ldots, k_s \in \mathcal{Z}_L$  and  $n_0, n_1, \ldots, n_s \geq 0$ , we have

$$
\operatorname{Res}_z\Big(\frac{\partial^{n_0+n_1+\cdots+n_s}w(z)}{\partial x^{n_0}\partial t_{k_1}^{n_1}\cdots\partial t_{k_s}^{n_s}}w^\star(z)\Big)=0\,.
$$

Remark 5.5. Condition (iii) can be suggestively rewritten as the following bilinear equation on the wave function  $w(z) = w(x, t; z)$ :

<span id="page-22-3"></span>
$$
\text{Res}_{z} (w(x, t; z) w^*(x', t'; z)) = 0.
$$
\n(5.13)

Indeed, the equation in (iii) formally coincides with the coefficient of

$$
\frac{(x-x')^{n_0}(t_{k_1}-t'_{k_1})^{n_1}\dots(t_{k_s}-t'_{k_s})^{n_s}}{n_0!n_1!\dots n_s!}
$$

in the Taylor series expansion of  $(5.13)$  around  $x = x'$ ,  $t = t'$ , see [\[Dic03\]](#page-37-17).

*Proof of Theorem [5.4.](#page-22-0)* By the definition [\(2.28\)](#page-9-3) of the action of  $\frac{\partial}{\partial t_k}$  on oscillating functions, we have

$$
\frac{\partial w(z)}{\partial t_k} = \left(\frac{\partial S(\partial)}{\partial t_k} S^{-1}(\partial) + S(\partial)\partial^k S^{-1}(\partial)\right) w(z).
$$

The equivalence of conditions (i) and (ii) is an obvious consequence of this identity and of Lemma [2.13.](#page-9-2) Indeed, when this holds, we have

$$
\tilde{B}_k(\partial) = \left(\frac{\partial S(\partial)}{\partial t_k} S^{-1}(\partial) + S(\partial)\partial^k S^{-1}(\partial)\right)_+.
$$

Next, let us prove that condition (ii) implies (iii). We claim, by induction on  $n_0 + n_1 + \cdots + n_s$ , that

<span id="page-23-0"></span>
$$
\frac{\partial^{n_0+n_1+\cdots+n_s}w(z)}{\partial x^{n_0}\partial t_{k_1}^{n_1}\cdots\partial t_{k_s}^{n_s}} = P(\partial)w(z)
$$
\n(5.14)

for some differential operator  $P(\partial) \in \mathcal{F}[\partial]$  (depending on  $n_0, n_1, \ldots, n_s, k_1, \ldots, k_s$ ). For  $n_0 = 1, n_1 = \cdots = n_s = 0$ , [\(5.14\)](#page-23-0) holds with  $P(\partial) = \partial$ . For  $n_0 = 0$  and  $n_1 + \cdots + n_s = 1$ , [\(5.14\)](#page-23-0) follows by condition (ii). Assuming, by induction, that  $(5.14)$  holds, we have

$$
\frac{\partial}{\partial x}\frac{\partial^{n_0+n_1+\cdots+n_s}w(z)}{\partial x^{n_0}\partial t_{k_1}^{n_1}\cdots\partial t_{k_s}^{n_s}}=\partial P(\partial)w(z)\,,
$$

and, for  $h \in \mathcal{Z}_L$ ,

$$
\frac{\partial}{\partial t_h} \frac{\partial^{n_0+n_1+\cdots+n_k}w(z)}{\partial x^{n_0}\partial t_{k_1}^{n_1}\cdots\partial t_{k_s}^{n_s}}=\Big(\frac{\partial P(\partial)}{\partial t_h}+P(\partial)\tilde{B}_h(\partial)\Big)w(z)\,,
$$

proving  $(5.14)$ . Then we have

$$
\operatorname{Res}_{z} \left( \frac{\partial^{n_0+n_1+\cdots+n_s} w(z)}{\partial x^{n_0} \partial t_{k_1}^{n_1} \cdots \partial t_{k_s}^{n_s}} w^{\star}(z) \right) = \operatorname{Res}_{z} \left( P(\partial) w(z) \right) w^{\star}(z)
$$

$$
= \operatorname{Res}_{z} (P \circ S)(z) (S^{\star})^{-1} (-z) = \operatorname{Res}_{z} P(z) = 0.
$$

In the third equality we used Lemma [2.8.](#page-7-2)

Finally, we prove that condition (iii) implies (i). For every  $n \geq 0$  and  $k \in \mathcal{Z}_L$ , we have, by condition (iii) and the definition [\(2.28\)](#page-9-3) of the action of  $\frac{\partial}{\partial t_k}$  on oscillating functions,

$$
0 = \text{Res}_{z} \left( \partial^{n} \frac{\partial w(z)}{\partial t_{k}} \right) w^{\star}(z) = \text{Res}_{z} \left( (z + \partial)^{n} \left( \frac{\partial S(z)}{\partial t_{k}} + S(z) z^{k} \right) \right) (S^{\ast})^{-1} (-z).
$$

We can now use Lemma  $2.9$  to conclude that condition (i) holds.

$$
\qquad \qquad \Box
$$

5.4. Summarizing statement. We can summarize the results of Sections [5.1](#page-20-7)–[5.3](#page-22-4) in the following Corollary:

<span id="page-23-1"></span>Corollary 5.6. Let  $\varphi = (\varphi_\alpha)_{\alpha \in I} \in \mathcal{F}^\ell$ . Let  $S(\partial) \in \mathcal{F}((\partial^{-1}))$  be a monic pseudodifferential operator of order 0 and let  $w(z) = S(z)e^{z \cdot t} \in \mathcal{F}((z^{-1}))e^{z \cdot t}$  be the corresponding oscillating function. Then, the dressing equation [\(5.2\)](#page-20-3) on  $S(\partial)$  is equivalent to the eigenvalue problem [\(5.9\)](#page-21-3) on the wave function  $w(z)$ . Furthermore, if either  $(5.2)$  or  $(5.9)$  holds, then the following conditions are equivalent:

- (a) the Sato equation  $(5.3)$  holds;
- (b) the linear problem  $(5.10)$  holds;
- $(c)$  the bilinear equation  $(5.13)$  holds.

If, moreover, any of the three equivalent conditions (a)–(c) hold, then  $\varphi$  is a solution to the Lax equation  $(3.12)$  (i.e.  $(4.1)$  holds).

*Proof.* The equivalence of  $(5.2)$  and  $(5.9)$  is given by Theorem  $5.3(a)$  $5.3(a)$ . Moreover, assuming  $(5.2)$ , the equivalence of (a) and (b) is given by Theorem  $5.3(b)$  $5.3(b)$ , and we need to prove their equivalence to (c). Theorem [5.4](#page-22-0) says, in particular, that (a) implies (c). Let us prove the converse implication. Since  $S(\partial)$  is monic of order 0, then  $\frac{\partial S(\partial)}{\partial t_k}S^{-1}(\partial)$  is a pseudodifferential operator of negative order. Hence, by Theorem  $5.4(i)$  $5.4(i)$  and  $(5.8)$ , we have

$$
\frac{\partial S(\partial)}{\partial t_k}S^{-1}(\partial)=-\big(S(\partial)\partial^kS^{-1}(\partial)\big)_-=-\big(L^{\frac{k}{N}}(\varphi,\partial)\big)_-=-B_k^{(-)}(\varphi;\partial)\,,
$$

proving the Sato equation [\(5.3\)](#page-20-4). Finally, the last assertion of the Theorem is given by Theorem [5.1\(](#page-20-1)a).  $\Box$ 

### 6. Wave functions and tau-functions

<span id="page-24-0"></span>6.1. Tau-functions of KP type. We review in this section the construction of tau-functions when  $\mathcal{Z}_L = \mathbb{Z}_{\geq 1} (= \mathcal{Z})$ , see [\[DJKM83,](#page-37-5) [Dic03\]](#page-37-17).

Throughout this section we assume that the algebra of functions  $\mathcal F$  is integrable (cf. Definition  $2.1$ ). Consider the operator

$$
e^{-z^{-1}\cdot\tilde{\partial}_t}: \mathcal{F} \to \mathcal{F}((z^{-1}))
$$
, where  $z^{-1}\cdot\tilde{\partial}_t := \sum_{k=1}^{\infty} \frac{z^{-k}}{k} \frac{\partial}{\partial t_k}$ .

By Taylor expansion, it "shifts" the time variables  $t_k \mapsto t_k - \frac{1}{kz^k}$ . Fix a monic pseudodifferential operator of order 0,  $S(\partial) \in 1 + \mathcal{F}[[\partial^{-1}]]\partial^{-1}$ . Let  $w(z) =$  $S(z)e^{z\cdot t} \in \mathcal{F}((z^{-1}))e^{z\cdot t}$  be the corresponding oscillating function, and let  $w^*(z)$  $(S^*)^{-1}(-z)e^{-z\cdot t}$  be the corresponding adjoint anti-oscillating function.

<span id="page-24-2"></span>**Theorem 6.1.** If the bilinear identity  $(5.13)$  holds, then, there exists a function  $\tau \in \mathcal{F}$  such that

<span id="page-24-1"></span>
$$
w(z) = \frac{\Gamma(t; z)\tau}{\tau} \quad \text{and} \quad w^\star(z) = \frac{\Gamma^\star(t; z)\tau}{\tau}, \tag{6.1}
$$

where  $\Gamma(t;z) = e^{z \cdot t} e^{-z^{-1} \cdot \tilde{\partial}_t}$ :  $\mathcal{F} \to \mathcal{F}((z^{-1})) e^{z \cdot t}$  and  $\Gamma^*(t;z) = e^{-z \cdot t} e^{z^{-1} \cdot \tilde{\partial}_t}$ :  $\mathcal{F} \to$  $\mathcal{F}((z^{-1}))e^{-z\cdot t}$  are called vertex operators. Moreover, the function  $\tau \in \mathcal{F}$  solving  $(6.1)$  is unique up to multiplication by a function of x (constant in t).

As an immediate corollary of Theorem [6.1](#page-24-2) and Corollary [5.6,](#page-23-1) we get the following

<span id="page-24-4"></span>**Corollary 6.2.** Let  $L(\partial) \in \mathcal{V}((\partial^{-1}))$  be as in [\(3.1\)](#page-10-2) and let  $\varphi = (\varphi_{\alpha})_{\alpha \in I}, \varphi_{\alpha} \in \mathcal{F}$ , be a solution of the hierarchy of Lax equations [\(4.1\)](#page-18-1) such that  $a_0(\varphi) = 0$ . Then, there exists  $\tau \in \mathcal{F}$  such that

<span id="page-24-3"></span>
$$
L(\varphi; \partial) = S(\partial) \partial^{N} S(\partial)^{-1} \quad \text{with} \quad S(z) = \frac{e^{-z^{-1} \cdot \partial_t \tau}}{\tau} \,. \tag{6.2}
$$

*Proof.* By Theorem [5.1\(](#page-20-1)b), there exists  $S(\partial) \in 1 + \mathcal{F}[[\partial^{-1}]]\partial^{-1}$  satisfying the dress-ing equation [\(5.2\)](#page-20-3), and by Theorem [6.1](#page-24-2) there exists  $\tau \in \mathcal{F}$  such that [\(6.2\)](#page-24-3) holds.  $\Box$ 

Corollary [6.2](#page-24-4) is saying that the Lax equations, which are systems of equations in (infinitely) many unknown functions  $a_i$ , can be "reduced" to a problem involving one single unknown function  $\tau$ . In fact, we can translate the bilinear equation [\(5.13\)](#page-22-3) (which, by Corollary  $5.6$ , is essentially equivalent to the Lax equations  $(3.12)$ ), to a system of partial differential equations on the unknown function  $\tau$ . The bilinear identity [\(5.13\)](#page-22-3) is translated to the following bilinear identity for  $\tau$ :

$$
\operatorname{Res}_{z} \left( \Gamma(\boldsymbol{t}';z) \tau \right) \left( \Gamma^{\star}(\boldsymbol{t}'';z) \tau \right) = \operatorname{Res}_{z} \left( e^{-z^{-1} \cdot \tilde{\partial}_{\boldsymbol{t}'}} \tau \right) \left( e^{z \cdot \tilde{\partial}_{\boldsymbol{t}''}} \tau \right) e^{z \cdot (\boldsymbol{t}' - \boldsymbol{t}'')} = 0 \, .
$$

By computing the above equation for  $t' = t - y$  and  $t'' = t + y$ , we get

<span id="page-25-1"></span>
$$
\sum_{k=0}^{\infty} p_k(-2\mathbf{y}) p_{k+1}(\tilde{\partial}_{\mathbf{y}}) (e^{-\mathbf{y}\cdot\partial_t}\tau) (e^{\mathbf{y}\cdot\partial_t}\tau) = 0, \qquad (6.3)
$$

where the Schur polynomials  $p_k(y) \in \mathbb{F}[y_1, \ldots, y_k]$  are defined by

<span id="page-25-2"></span>
$$
e^{\sum_{k=1}^{\infty} y_k z^k} = \sum_{k=0}^{\infty} p_k(\mathbf{y}) z^k.
$$
 (6.4)

Equation  $(6.3)$  indeed provides an infinite system of partial differential equations on  $\tau$ , by looking at the coefficient of each monomial  $y_1^{n_1} y_2^{n_2} \dots y_s^{n_s}$ .

We can express the solution  $L(\varphi; \partial)$  in terms of the function  $\tau$ . Recalling that  $w(z) = S(z)e^{z \cdot t}$ , using equations [\(6.1\)](#page-24-1) and [\(6.4\)](#page-25-2) we get  $S(z) = \sum_{k=0}^{\infty} \frac{p_k(-\tilde{\partial}_t)\tau}{\tau}z^{-k}$ . Similarly, by  $w^*(z) = (S^*)^{-1}(-z)e^{-z \cdot t}$  we get  $(S^*)^{-1}(-z) = \sum_{k=0}^{\infty} \frac{p_k(\tilde{\partial}_t)\tau}{\tau}z^{-k}$ , from which follows that  $S^{-1}(z) = \sum_{k=0}^{\infty} (z + \partial)^{-k} \frac{p_k(\tilde{\partial}_t)\tau}{\tau}$ . Hence, if  $\tau$  is a tau-function of the solution  $\varphi$  of the Lax equation [\(3.12\)](#page-13-0) such that  $a_0(\varphi) = 0$ , we obtain

<span id="page-25-0"></span>
$$
L(\varphi; \partial) = \sum_{h,k=0}^{\infty} \frac{p_h(-\tilde{\partial}_t)\tau}{\tau} \partial^{N-h-k} \circ \frac{p_k(\tilde{\partial}_t)\tau}{\tau}
$$
  
= 
$$
\sum_{k=0}^{\infty} \left( \sum_{h=0}^{k} \sum_{l=0}^{k-h} {N-k+l \choose l} \frac{p_{k-h-l}(-\tilde{\partial}_t)\tau}{\tau} \partial^l \left( \frac{p_h(\tilde{\partial}_t)\tau}{\tau} \right) \right) \partial^{N-k}.
$$
 (6.5)

By equating powers of  $\partial^{N-2}$  in both sides of the above formula we get for example

$$
a_1(\varphi) = N\partial^2 \log \tau.
$$

(We used the fact that  $\partial \log \tau = -s_1 = \frac{\partial}{\partial t_1} \log \tau$ , cf. equation [\(6.11\)](#page-26-0) below.)

The remainder of the present section is devoted to the proof of Theorem [6.1.](#page-24-2) Equation [\(6.1\)](#page-24-1) defining  $\tau$ , which we want to solve, is obviously equivalent to the the equation relating  $S(z)$  and  $\tau$  in [\(6.2\)](#page-24-3). Applying the logarithm to both sides of the latter, we get

<span id="page-25-3"></span>
$$
\log(S(z)) = \left(e^{-z^{-1}\cdot\tilde{\partial}_t} - 1\right) \log \tau. \tag{6.6}
$$

For this, we used the (obvious) fact that the operator  $e^{-z^{-1}\cdot\tilde{\partial}_t}$ , describing the "shift" of variables  $t_k \mapsto t_k - \frac{1}{k^2}$ , commutes with taking the logarithm. Note that the logarithm of  $S(z)$  is well defined by power series expansion in  $z^{-1}$ , since  $S(z) \in 1 + \mathcal{F}[[z^{-1}]]z^{-1}$ . The outline of the proof of Theorem [6.1](#page-24-2) is as follows. From equation  $(6.6)$  we shall derive an (equivalent) system of partial differential equations of type [\(2.3\)](#page-4-2) for  $\phi = \log \tau$ , depending on a function f of x. We will then prove that such system is compatible. Then, by the integrability assumption on F, it will uniquely define  $\phi$  up to the choice of the additive function f, i.e. it will uniquely define  $\tau$  up to a factor depending only on x, as claimed by Theorem [6.1.](#page-24-2)

First, we derive from [\(6.6\)](#page-25-3) an equation for  $\partial \log \tau$ . For this, we need the following result.

<span id="page-25-5"></span>Lemma 6.3. The following identities hold:

a)  $S(z)e^{-z^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-z)=1.$ b)  $\partial \log S(z) = (1 - e^{-z^{-1} \cdot \tilde{\partial}_t}) s_1$ .

*Proof.* We set  $t'_{k} = t_{k} - \frac{1}{kw^{k}}$ ,  $k \geq 1$ , in the bilinear identity [\(5.13\)](#page-22-3), to get

<span id="page-25-4"></span>
$$
\operatorname{Res}_{z} \left( S(x, t; z) e^{z \cdot t} e^{-z \cdot t'} (S^*)^{-1} (x', t'; -z) \right) = 0. \tag{6.7}
$$

A straightforward computation leads to

<span id="page-26-4"></span>
$$
e^{z \cdot t} e^{-z \cdot t'} = \iota_w \left(1 - \frac{z}{w}\right)^{-1} e^{z(x - x')}.
$$
\n(6.8)

Hence, equation [\(6.7\)](#page-25-4) gives

Res<sub>z</sub> 
$$
(\iota_w(z-w)^{-1}e^{z(x-x')}S(x,\mathbf{t};z)e^{-w^{-1}\cdot\tilde{\partial}_{\mathbf{t}}}(S^*)^{-1}(x',\mathbf{t};-z))=0,
$$

which, by the second equation in  $(2.22)$ , gives

<span id="page-26-1"></span>
$$
\left(e^{z(x-x')}S(x,t;z)e^{-w^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(x',t;-z)\right)\Big|_{z=w}=0\,,\tag{6.9}
$$

where the − means that we need to take the negative powers of z, before setting  $z = w$ . Clearly, we can rewrite [\(6.9\)](#page-26-1) as

$$
e^{w(x-x')}S(x,t;w)e^{-w^{-1}\cdot\tilde{\partial}_{t}}(S^{*})^{-1}(x',t;-w)
$$
  
= 
$$
\left(e^{z(x-x')}S(x,t;z)e^{-w^{-1}\cdot\tilde{\partial}_{t}}(S^{*})^{-1}(x',t;-z)\right)_{+}\Big|_{z=w}.
$$
 (6.10)

<span id="page-26-2"></span>Setting  $x' = x$  and recalling that  $S(z) \in 1 + \mathcal{F}[[z^{-1}]]z^{-1}$ , we immediately get claim (a). Instead, applying  $\frac{\partial}{\partial x}$  to both sides of [\(6.10\)](#page-26-2) and then setting  $x' = x$ , the LHS is

$$
((w+\partial)S(w))e^{-w^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-w)=w+\frac{\partial S(w)}{S(w)}=w+\partial\log S(w),
$$

where we used part (a). On the other hand, the RHS of [\(6.10\)](#page-26-2) gives, since  $S(z)$  =  $1 + s_1 z^{-1} + \dots$  and  $(S^*)^{-1}(-z) = 1 - s_1 z^{-1} + \dots$ ,

$$
\left( ((z+\partial)S(z))e^{-w^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-z)\right)_+|_{z=w} = w + s_1 - e^{-w^{-1}\cdot\tilde{\partial}_t}s_1.
$$
  
Claim (b) follows.

Comparing claim (b) in Lemma  $6.3$  and equation  $(6.6)$ , we get

 $(e^{-z^{-1}\cdot\tilde{\partial}_t}-1)(\partial \log \tau+s_1)=0$ .

On the other hand, the kernel of  $e^{-z^{-1}\cdot\tilde{\partial}_t}-1$  contains all functions of x, which are independent of t. We thus conclude that  $\tau$  must satisfy

<span id="page-26-0"></span>
$$
\partial \log \tau = -s_1 + f \,, \tag{6.11}
$$

where f is such that  $\frac{\partial f}{\partial t_k} = 0$  for every  $k \geq 1$ , and we may set it equal to 0 without loss of generality.

Next, we want to derive from [\(6.6\)](#page-25-3) an equation for  $\frac{\partial}{\partial t_k} \log \tau, k \geq 1$ . We introduce the following differential operator

<span id="page-26-5"></span>
$$
N(z) = \frac{\partial}{\partial z} - \sum_{k=1}^{\infty} z^{-k-1} \frac{\partial}{\partial t_k}.
$$
 (6.12)

<span id="page-26-3"></span>**Lemma 6.4.**  $N(z) (e^{-z^{-1} \cdot \tilde{\partial}_t} f) = 0$  for every  $f \in \mathcal{F}$ .

Proof. We have

$$
\frac{\partial}{\partial z}e^{-z^{-1}\cdot\tilde{\partial}_t} = \frac{\partial}{\partial z}\exp\{-\sum_k \frac{1}{kz^k}\frac{\partial}{\partial t_k}\} = \left(\sum_{k=1}^{\infty}z^{-k-1}\frac{\partial}{\partial t_k}\right)\circ e^{-z^{-1}\cdot\tilde{\partial}_t}.
$$

We then apply the operator  $N(z)$  to equation [\(6.6\)](#page-25-3) and use Lemma [6.4](#page-26-3) to get

<span id="page-27-0"></span>
$$
N(z)\log(S(z)) = \sum_{k=1}^{\infty} z^{-k-1} \frac{\partial}{\partial t_k} \log \tau.
$$
 (6.13)

Therefore, equation  $(6.1)$  is equivalent to  $(6.13)$  which, by looking at the various powers of  $z$ , is equivalent to the following system of equations:

<span id="page-27-1"></span>
$$
\frac{\partial}{\partial t_k} \log \tau = \text{Res}_z \, z^k N(z) \log(S(z)) =: g_k \,. \tag{6.14}
$$

In order to prove Theorem [6.1,](#page-24-2) we only need to show that the system of equations consisting of [\(6.11\)](#page-26-0) and [\(6.14\)](#page-27-1) admits a solution  $\phi := \log \tau \in \mathcal{F}$ , unique up to an additive constant. Equivalently, since by assumption  $\mathcal F$  is integrable, we only need to prove that that system is compatible, i.e.

<span id="page-27-2"></span>
$$
\partial g_k = -\frac{\partial s_1}{\partial t_k} \quad \text{for all } k \ge 1 \,, \tag{6.15}
$$

and

<span id="page-27-3"></span>
$$
\frac{\partial g_k}{\partial t_h} = \frac{\partial g_h}{\partial t_k} , \quad \text{for every } h, k \ge 1.
$$
 (6.16)

Applying  $N(z)$  to both sides of the equation in Lemma [6.3\(](#page-25-5)b) and taking the coefficient of  $z^{-k-1}$ , we immediately get [\(6.15\)](#page-27-2) thanks to Lemma [6.4.](#page-26-3) Next, we prove equation [\(6.16\)](#page-27-3). For this, we shall need the following result.

# <span id="page-27-5"></span>Lemma 6.5. The following identities hold:

(a)  $(e^{-w^{-1}\cdot\tilde{\partial}_t} - 1) \log S(z) = (e^{-z^{-1}\cdot\tilde{\partial}_t} - 1) \log S(w).$ (b)  $N(z)N(w) \log S(z) = N(z)N(w) \log S(w)$ .

<span id="page-27-4"></span>*Proof.* Setting 
$$
x' = x
$$
 and  $t'_k = t_k - \frac{1}{kw_1^k} - \frac{1}{kw_2^k}$ , we have, as in (6.8),  
\n
$$
e^{z \cdot t} e^{-z \cdot t'} = \frac{w_1 w_2}{w_2 - w_1} \left( t_{w_1} (w_1 - z)^{-1} - t_{w_2} (w_2 - z)^{-1} \right).
$$
\n(6.17)

Applying the bilinear identity  $(5.13)$  and using equation  $(6.17)$  we get

$$
0 = \text{Res}_{z} \left( \left( \iota_{w_{1}} (w_{1} - z)^{-1} - \iota_{w_{2}} (w_{2} - z)^{-1} \right) S(z) e^{-w_{1}^{-1} \cdot \tilde{\partial}_{t}} e^{-w_{2}^{-1} \cdot \tilde{\partial}_{t}} (S^{*})^{-1} (-z) \right)
$$
  
\n
$$
= \left( S(z) e^{-w_{1}^{-1} \cdot \tilde{\partial}_{t}} e^{-w_{2}^{-1} \cdot \tilde{\partial}_{t}} (S^{*})^{-1} (-z) \right)_{-} \Big|_{z=w_{2}}
$$
  
\n
$$
- \left( S(z) e^{-w_{1}^{-1} \cdot \tilde{\partial}_{t}} e^{-w_{2}^{-1} \cdot \tilde{\partial}_{t}} (S^{*})^{-1} (-z) \right)_{-} \Big|_{z=w_{1}}
$$
  
\n
$$
= S(w_{2}) e^{-w_{1}^{-1} \cdot \tilde{\partial}_{t}} \frac{1}{S(w_{2})} - S(w_{1}) e^{-w_{2}^{-1} \cdot \tilde{\partial}_{t}} \frac{1}{S(w_{1})} .
$$

In the second equality we used the second equation in  $(2.22)$ . The  $-$  sign here means that we need to take negative powers in z before substituting  $z = w_2$  in the first parenthesis and  $z = w_1$  in the second. In the third equality we used Lemma  $6.3(a)$  $6.3(a)$ . We rewrite the above identity as

$$
\frac{e^{-w^{-1}\cdot\tilde{\partial}_t}S(z)}{S(z)} = \frac{e^{-z^{-1}\cdot\tilde{\partial}_t}S(w)}{S(w)}
$$

and taking logarithm of both sides, we get claim (a). Moreover, applying  $N(z)N(w)$ to both sides of (a) and using Lemma [6.4,](#page-26-3) we get claim (b).  $\Box$ 

We can finally prove equation [\(6.16\)](#page-27-3). We have, by the definition [\(6.14\)](#page-27-1) of  $g_k$  and Lemma  $6.5(b)$  $6.5(b)$ ,

$$
\sum_{h,k=1}^{\infty} \left( \frac{\partial g_h}{\partial t_k} - \frac{\partial g_k}{\partial t_h} \right) z^{-h-1} w^{-k-1}
$$
\n
$$
= \sum_{k=1}^{\infty} w^{-k-1} \frac{\partial}{\partial t_k} \left( \sum_{h=1}^{\infty} g_h z^{-h-1} \right) - \sum_{h=1}^{\infty} z^{-h-1} \frac{\partial}{\partial t_h} \left( \sum_{k=1}^{\infty} g_k w^{-k-1} \right)
$$
\n
$$
= -N(w)N(z) \log S(z) + N(z)N(w) \log S(w) = 0.
$$

This proves [\(6.16\)](#page-27-3) and concludes the proof of Theorem [6.1.](#page-24-2)

6.2. Tau-functions of BKP and CKP type. We next want to construct taufunctions when  $\mathcal{Z}_L = 2\mathbb{Z}_{\geq 0}+1$ . The generic situation of this is provided by the BKP and CKP hierarchies, studied in [\[DJKM81,](#page-37-4) [DJKM83,](#page-37-5) [DMH09,](#page-37-26) [CW13,](#page-37-18) [KZ20\]](#page-37-19) and other papers. Throughout this section we assume that the algebra of functions  $\mathcal{F}$ , in the space variable x and odd time variables  $t_k, k \in \mathcal{Z}_L$ , is integrable (cf. Definition [2.1\)](#page-4-6). Recall the oscillating/antioscillating functions  $e^{\pm z \cdot t} = e^{\pm (zx + \sum_{k \in \mathcal{Z}_L} z^k t_k)}$ , and introduce the operator

$$
e^{-2z^{-1}\cdot\tilde{\partial}_t}
$$
:  $\mathcal{F} \to \mathcal{F}((z^{-1}))$ , where  $2z^{-1}\cdot\tilde{\partial}_t = 2\sum_{k\in\mathcal{Z}_L} \frac{z^{-k}}{k} \frac{\partial}{\partial t_k}$ .

By Taylor expansion, it "shifts" the time variables  $t_k \mapsto t_k - \frac{2}{kz^k}$ ,  $k \in \mathcal{Z}_L$ . Fix a monic pseudodifferential operator  $S(\partial) = 1 + \sum_{i=1}^{\infty} s_i \partial^{-i} \in \mathbb{1} + \mathcal{F}[[\partial^{-1}]]\partial^{-1}$  of order 0. Let  $w(z) = S(z)e^{z \cdot t} \in \mathcal{F}((z^{-1}))e^{z \cdot t}$  be the corresponding oscillating function, and let  $w^*(z) = (S^*)^{-1}(-z)e^{-z \cdot t}$  be the corresponding adjoint anti-oscillating function.

The analogue of Theorem [6.1](#page-24-2) in the case when  $\mathcal{Z}_L$  consists of odd integers is the following.

<span id="page-28-2"></span>Theorem 6.6. Suppose that the bilinear identity [\(5.13\)](#page-22-3) holds, and that  $S(\partial)$  satisfies the following condition

<span id="page-28-0"></span>
$$
S^*(\partial) = \partial^n S(\partial)^{-1} \partial^{-n}.
$$
\n(6.18)

for  $n = 0$ , or 1. Then, there exists a function  $\tau \in \mathcal{F}$  such that

<span id="page-28-4"></span>
$$
w(z) = \left(1 + z^{-1} \mathcal{E}(z)\right)^{\frac{1-n}{2}} \frac{\Gamma(z)\tau}{\tau},\tag{6.19}
$$

where

$$
\Gamma(z) = \Gamma(x, t; z) = e^{z \cdot t} e^{-2z^{-1} \cdot \tilde{\partial}_t} : \mathcal{F} \to \mathcal{F}((z^{-1})) e^{z \cdot t},
$$

and

<span id="page-28-5"></span>
$$
\mathcal{E}(z) = (e^{-2z^{-1}\cdot\tilde{\partial}_t} - 1)\partial \log \tau \in \mathcal{F}[[z^{-1}]]z^{-1}.
$$
 (6.20)

Moreover, the function  $\tau \in \mathcal{F}$  solving [\(6.1\)](#page-24-1) is unique up to multiplication by a function of  $x$  (constant in  $t$ ).

<span id="page-28-3"></span>**Lemma 6.7.** Let  $L(\partial) \in \mathcal{V}((\partial^{-1}))$  be as in  $(3.1)$ , satisfying the constraint

<span id="page-28-1"></span>
$$
L^*(\partial) = (-1)^N \partial^n L(\partial) \partial^{-n} \quad \text{for } n = 0 \text{ or } 1. \tag{6.21}
$$

Assume that the dressing equation [\(5.2\)](#page-20-3) has solution. Then, there exists  $S(\partial) \in$  $1 + \mathcal{F}[[\partial^{-1}]]\partial^{-1}$  solving the dressing equation and such that condition [\(6.18\)](#page-28-0) holds.

*Proof.* Let  $\tilde{S}(\partial) \in 1 + \mathcal{F}[[\partial^{-1}]] \partial^{-1}$  be a solution to the dressing equation [\(5.2\)](#page-20-3). The constraint [\(6.21\)](#page-28-1) implies that

$$
L(\partial) = \tilde{S}(\partial)\partial^N \tilde{S}^{-1}(\partial) = (\partial^{-n} (\tilde{S}^*)^{-1}(\partial)\partial^n)\partial^N (\partial^{-n} (\tilde{S}^*)^{-1}(\partial)\partial^n)^{-1}.
$$

Hence, by Theorem [5.1,](#page-20-1) we have

<span id="page-29-0"></span>
$$
\partial^{-n} (\tilde{S}^*)^{-1} (\partial) \partial^n = \tilde{S}(\partial) K(\partial) , \qquad (6.22)
$$

where  $K(\partial) \in 1 + \mathbb{F}[[\partial^{-1}]]\partial^{-1}$ . Taking the adjoint of both sides of equation [\(6.22\)](#page-29-0) we get  $K^*(\partial) = \partial^n K(\partial) \partial^{-n} = K(\partial)$ . Hence, the symbol of  $K(\partial)$  satisfies  $K(z) =$ K(-z). Let  $S(\partial) = \tilde{S}(\partial)C(\partial)$ , with  $C(\partial) \in 1 + \mathbb{F}[[\partial^{-1}]]\partial^{-1}$  be another solution to the dressing equation  $(5.2)$ . To conclude the proof we need to show that there exists  $C(\partial)$  such that the constraint [\(6.18\)](#page-28-0) holds. Replacing  $S(\partial) = \tilde{S}(\partial)C(\partial)$  in  $(6.18)$ , we get that it is equivalent to the condition

$$
C^*(\partial)\tilde{S}^*(\partial) = \partial^n C^{-1}(\partial)\tilde{S}^{-1}(\partial)\partial^{-n} = C^{-1}(\partial)\partial^n\tilde{S}^{-1}(\partial)\partial^{-n} = C^{-1}(\partial)K^*(\partial)\tilde{S}^*(\partial).
$$

In the last equality we used the adjoint of equation  $(6.22)$ . The claim follows from the above equation and the fact that  $C(z)C(-z) = K(z)$  has a solution in  $1 + \mathbb{F}[[z^{-1}]]z^{-1}$  since  $K(z) = K(-z)$ .

As an immediate corollary of Theorem [6.6,](#page-28-2) Lemma [6.7](#page-28-3) and Corollary [5.6,](#page-23-1) we get the following

Corollary 6.8. Let  $L(\partial) \in \mathcal{V}((\partial^{-1}))$  be as in [\(3.1\)](#page-10-2), satisfying the constraint [\(6.21\)](#page-28-1). Let  $\varphi = (\varphi_{\alpha})_{\alpha \in I}$ ,  $\varphi_{\alpha} \in \mathcal{F}$ , be a solution of the hierarchy of Lax equations [\(4.1\)](#page-18-1) such that  $a_0(\varphi) = 0$ . Then, there exists  $\tau \in \mathcal{F}$  such that

<span id="page-29-1"></span>
$$
L(\varphi; \partial) = S(\partial) \partial^N S(\partial)^{-1} \quad with \quad S(z) = \left(1 + z^{-1} \mathcal{E}(z)\right)^{\frac{1-n}{2}} \frac{e^{-2z^{-1} \cdot \partial_t \tau}}{\tau} \,. \tag{6.23}
$$

*Proof.* By Theorem [5.1\(](#page-20-1)b), there exists  $S(\partial) \in 1 + \mathcal{F}[[\partial^{-1}]]\partial^{-1}$  satisfying the dress-ing equation [\(5.2\)](#page-20-3). Since  $L(\partial)$  satisfies the constraint [\(6.21\)](#page-28-1), the corresponding  $S(\partial)$ can be chosen so that  $(6.18)$  holds by Lemma [6.7.](#page-28-3) Hence by Theorem [6.6](#page-28-2) there exists  $\tau \in \mathcal{F}$  such that [\(6.23\)](#page-29-1) holds. □

The remainder of this section will be devoted to the proof of Theorem [6.6.](#page-28-2) First note that equation [\(6.19\)](#page-28-4) defining  $\tau$ , which we want to solve, is obviously equivalent to the the equation relating  $S(z)$  and  $\tau$  in [\(6.23\)](#page-29-1). Applying the logarithm to both sides of the latter, we get, as we did in  $(6.6)$ ,

<span id="page-29-2"></span>
$$
\log(S(z)) = \left(e^{-2z^{-1}\cdot\tilde{\partial}_t} - 1\right)\log\tau + \frac{1-n}{2}\log\left(1 + z^{-1}\mathcal{E}(z)\right). \tag{6.24}
$$

Follow the line of reasoning used for the proof of Theorem [6.1,](#page-24-2) the proof of Theorem [6.6](#page-28-2) will be obtained as follows. From equation  $(6.24)$  we shall derive an (equivalent) system of partial differential equations of type [\(2.3\)](#page-4-2) for  $\phi = \log \tau$ , depending on a function  $f$  of x. We will then prove that such a system is compatible. The claim of Theorem [6.6](#page-28-2) will then be a consequence of the integrability assumption on  $\mathcal{F}$ .

In analogy with  $(6.20)$ , for  $i > 1$  we denote

<span id="page-29-4"></span>
$$
\mathcal{E}_i(z) = (e^{-2z^{-1}\cdot \tilde{\partial}_t} - 1)s_i.
$$
\n(6.25)

In order to derive from  $(6.24)$  an equation for  $\partial \log \tau$ , we need the following analogue of Lemma [6.3.](#page-25-5)

**Lemma 6.9.** *a)* The following identity holds:

<span id="page-29-5"></span>
$$
S(z)e^{-2z^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-z) = 1 - \frac{1}{2}z^{-1}\mathcal{E}_1(z).
$$
 (6.26)

b) The following identity holds:

<span id="page-29-3"></span>
$$
\partial \log S(z) = -\frac{1}{2} \mathcal{E}_1(z) - \frac{1}{2} z^{-1} \left( 1 - \frac{1}{2} z^{-1} \mathcal{E}_1(z) \right)^{-1} \left( \mathcal{E}_2(z) + \partial \mathcal{E}_1(z) - s_1 \mathcal{E}_1(z) - \frac{1}{2} \mathcal{E}_1(z)^2 \right). \tag{6.27}
$$

*Proof.* In analogy with the proof of Lemma [6.3](#page-25-5) we set  $t'_{k} = t_{k} - \frac{2}{kw^{k}}$ ,  $k \ge 1$ , in the bilinear identity  $(5.13)$  or, equivalently,  $(6.7)$ . A straightforward computation leads to

<span id="page-30-6"></span>
$$
e^{z \cdot t} e^{-z \cdot t'} = \left(1 + \frac{z}{w}\right) \iota_w \left(1 - \frac{z}{w}\right)^{-1} e^{z(x - x')} \,. \tag{6.28}
$$

Hence, equation [\(6.7\)](#page-25-4) gives

$$
\operatorname{Res}_{z} \left( (w+z)\iota_{w}(w-z)^{-1} e^{z(x-x')} S(x,\boldsymbol{t};z) e^{-2w^{-1}\cdot\tilde{\partial}_{\boldsymbol{t}}}(S^{*})^{-1}(x',\boldsymbol{t};-z) \right) = 0,
$$

which, by the second equation in  $(2.22)$ , gives

<span id="page-30-0"></span>
$$
\left( (w+z)e^{z(x-x')}S(x,t;z)e^{-2w^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(x',t;-z) \right)_{-}\Big|_{z=w} = 0, \quad (6.29)
$$

where, as in  $(6.9)$ , the − means that we need to take the negative powers of z, before setting  $z = w$ . We rewrite  $(6.29)$  as

<span id="page-30-1"></span>
$$
2we^{w(x-x')}S(x,t;w)e^{-2w^{-1}\cdot\tilde{\partial}_{t}}(S^{*})^{-1}(x',t;-w)
$$
  
= 
$$
\left((w+z)e^{z(x-x')}S(x,t;z)e^{-2w^{-1}\cdot\tilde{\partial}_{t}}(S^{*})^{-1}(x',t;-z)\right)_{+}\Big|_{z=w}.
$$
 (6.30)

A straightforward computation shows that, if  $S(z) = 1 + s_1 z^{-1} + s_2 z^{-2} + \dots$ , then

<span id="page-30-2"></span>
$$
(S^*)^{-1}(-z) = 1 - s_1 z^{-1} + (s_1^2 - s_2 - s_1') z^{-2} + \dots
$$
 (6.31)

Setting  $x' = x$  in [\(6.30\)](#page-30-1), the RHS becomes

$$
\left( (w+z)S(z)e^{-2w^{-1}\cdot\tilde{\partial}_{t}}(S^{*})^{-1}(-z) \right)_{+}\Big|_{z=w} = 2w + s_{1} - e^{-w^{-1}\cdot\tilde{\partial}_{t}}s_{1} = 2w - \mathcal{E}_{1}(w).
$$

Claim (a) follows. Next, we apply  $\frac{\partial}{\partial x}$  to both sides of [\(6.30\)](#page-30-1) and then we set  $x' = x$ . The LHS gives

$$
2w((w+\partial)S(w))e^{-2w^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-w),
$$

which, by part (a), is

<span id="page-30-3"></span>
$$
2w^{2} - w\mathcal{E}_{1}(w) + 2w\left(1 - \frac{1}{2}w^{-1}\mathcal{E}_{1}(w)\right)\partial \log S(w).
$$
 (6.32)

On the other hand, the RHS of  $(6.30)$  gives, using  $(6.31)$ ,

$$
\begin{aligned} & \left( (w+z) \big( (z+\partial)S(z) \big) e^{-2w^{-1} \cdot \tilde{\partial}_t} (S^*)^{-1}(-z) \right)_+ \Big|_{z=w} \\ & = 2w^2 - 2w \mathcal{E}_1(w) - \mathcal{E}_2(w) - \partial \mathcal{E}_1(w) + e^{-2w^{-1} \cdot \tilde{\partial}_t} s_1^2 - s_1 e^{-2w^{-1} \cdot \tilde{\partial}_t} s_1 \,, \end{aligned}
$$

which can be rewritten as

<span id="page-30-4"></span>
$$
2w^{2} - 2w\mathcal{E}_{1}(w) - \mathcal{E}_{2}(w) - \partial \mathcal{E}_{1}(w) + \mathcal{E}_{1}(w)^{2} + s_{1}\mathcal{E}_{1}(w), \qquad (6.33)
$$

since, obviously,  $e^{-2w^{-1}\cdot\tilde{\partial}_t} s_1^2 = (e^{-2w^{-1}\cdot\tilde{\partial}_t} s_1)^2 = (\mathcal{E}_1(w) + s_1)^2$ . Equating [\(6.32\)](#page-30-3) and [\(6.33\)](#page-30-4), a straightforward computation leads to  $(6.27)$ .  $\Box$ 

So far we did not use the assumption that  $S(\partial)$  satisfies the condition [\(6.18\)](#page-28-0). If we do, the equation  $(6.27)$  greatly simplifies.

**Lemma 6.10.** Assume that  $S(\partial)$  satisfies [\(6.18\)](#page-28-0) for some  $n \in \mathbb{Z}$ . Then, the following equations holds

(a)  $s_2 = \frac{1}{2}s_1^2 - \frac{n+1}{2}s_1'$ . (b)  $\mathcal{E}_2(z) = \frac{1}{2}\mathcal{E}_1(z)^2 + s_1 \mathcal{E}_1(z) - \frac{n+1}{2}\partial \mathcal{E}_1(z)$ . (c) Equation [\(6.27\)](#page-29-3) reduces to

<span id="page-30-5"></span>
$$
\partial \log S(z) = -\frac{1}{2} \mathcal{E}_1(z) + \frac{1-n}{2} \partial \log \left(1 - \frac{1}{2} z^{-1} \mathcal{E}_1(z)\right). \tag{6.34}
$$

*Proof.* Expanding the identity  $S^*(\partial)\partial^n S(\partial) = \partial^n$  we get

$$
\partial^{n} + (2s_2 + (n+1)s'_1 - s_1^2)\partial^{n-2} + \cdots = \partial^{n}.
$$

Claim (a) follows from this identity. Claim (b) is an immediate consequence of (a) and the definition [\(6.25\)](#page-29-4) of  $\mathcal{E}_i(z)$ . Claim (c) follows from (b).

Comparing equations  $(6.24)$  and  $(6.34)$ , we get, recalling  $(6.20)$ 

$$
\mathcal{E}(z) + \frac{1-n}{2} \partial \log (1 + z^{-1} \mathcal{E}(z)) = -\frac{1}{2} \mathcal{E}_1(z) + \frac{1-n}{2} \partial \log (1 - \frac{1}{2} z^{-1} \mathcal{E}_1(z)),
$$

which implies

<span id="page-31-6"></span>
$$
\mathcal{E}(z) = -\frac{1}{2}\mathcal{E}_1(z). \tag{6.35}
$$

Since the kernel of  $e^{-z^{-1}\cdot\tilde{\partial}_t} - 1$  contains all functions of x, which are independent of t, we then conclude, as in [\(6.11\)](#page-26-0), that  $\tau$  must satisfy

<span id="page-31-1"></span>
$$
\partial \log \tau = -\frac{1}{2}s_1 + f,\tag{6.36}
$$

where f is a function constant in all  $t_k$ 's, which, without loss of generality, we may set equal to 0.

Next, we want to derive from [\(6.24\)](#page-29-2) an equation for  $\frac{\partial}{\partial t_k} \log \tau$ ,  $k \ge 1$ . In analogy with  $(6.12)$ , we introduce the differential operator

<span id="page-31-5"></span>
$$
N(z) = \frac{\partial}{\partial z} - 2 \sum_{k \in \mathcal{Z}_L}^{\infty} z^{-k-1} \frac{\partial}{\partial t_k},
$$
\n(6.37)

and we observe that the analogue of Lemma [6.4](#page-26-3) still holds:

<span id="page-31-0"></span>
$$
N(z)\left(e^{-2z^{-1}\cdot\tilde{\partial}_t}f\right) = 0 \text{ for every } f \in \mathcal{F}.
$$
 (6.38)

We then apply the operator  $N(z)$  to equation [\(6.24\)](#page-29-2) and use equation [\(6.38\)](#page-31-0) to get

$$
N(z)\Big(\log S(z) - \frac{1-n}{2}\log\left(1+z^{-1}\mathcal{E}(z)\right)\Big) = 2\sum_{k\in\mathcal{Z}_L}^{\infty} z^{-k-1}\frac{\partial}{\partial t_k}\log\tau.
$$

By looking at the various powers of  $z$ , this is equivalent to the following system of equations:

<span id="page-31-2"></span>
$$
\frac{\partial}{\partial t_k} \log \tau = \frac{1}{2} \operatorname{Res}_z z^k N(z) \Big( \log S(z) - \frac{1-n}{2} \log \left( 1 + z^{-1} \mathcal{E}(z) \right) \Big) =: g_k, \ k \in \mathcal{Z}_L. \tag{6.39}
$$

In order to prove Theorem [6.6,](#page-28-2) by the integrability assumption on  $\mathcal{F}$ , it remains to show that the system of equations consisting of  $(6.36)$  and  $(6.39)$  is compatible. i.e.

<span id="page-31-3"></span>
$$
\partial g_k = -\frac{1}{2} \frac{\partial s_1}{\partial t_k} \quad \text{for every } k \in \mathcal{Z}_L = 2\mathbb{Z}_{\geq 0} + 1 \,, \tag{6.40}
$$

and

<span id="page-31-4"></span>
$$
\frac{\partial g_k}{\partial t_h} = \frac{\partial g_h}{\partial t_k} , \quad \text{for every } h, k \in \mathcal{Z}_L .
$$
 (6.41)

Applying  $N(z)$  to both sides of equation [\(6.34\)](#page-30-5) and taking the coefficient of  $z^{-k-1}$ , we immediately get [\(6.40\)](#page-31-3), thanks to equation [\(6.38\)](#page-31-0). We are left to prove equation  $(6.41)$ . We can write  $(6.41)$  as a formal power series by multiplying both

sides by  $2z^{-h-1}w^{-k-1}$  and summing over  $h, k \in \mathcal{Z}_L$ . As a result, recalling the definition  $(6.37)$  of  $N(z)$ , equation  $(6.41)$  becomes

$$
0 = 2 \sum_{h,k \in \mathcal{Z}_L} z^{-h-1} w^{-k-1} \left( \frac{\partial g_k}{\partial t_h} - \frac{\partial g_h}{\partial t_k} \right)
$$
  
= 
$$
2 \sum_{h \in \mathcal{Z}_L} z^{-h-1} \frac{\partial}{\partial t_h} \sum_{k \in \mathcal{Z}_L} w^{-k-1} g_k - 2 \sum_{k \in \mathcal{Z}_L} w^{-k-1} \frac{\partial}{\partial t_k} \sum_{h \in \mathcal{Z}_L} z^{-h-1} g_h
$$

Notice that  $2\sum_{h\in\mathcal{Z}_L}z^{-h-1}\frac{\partial}{\partial t_h}$  can be replaced by  $-N(z)$ , since it acts on a functions which is constant in z. Moreover, recalling the definition  $(6.39)$  of  $g_k$ , we have that  $\sum_{k\in\mathcal{Z}_L} w^{-k-1}g_k$  coincides with the even part (i.e. even powers of w) of  $N(w)\Big(\log S(w)-\frac{1-n}{2}\log\left(1+w^{-1}\mathcal{E}(w)\right)\Big).$  Hence, recalling  $(6.35)$ , equation  $(6.41)$ is translated in the following equation

<span id="page-32-0"></span>
$$
N(z)\Big(N(w)\log S(w) - \frac{1-n}{2}N(w)\log\left(1 - \frac{1}{2}w^{-1}\mathcal{E}_1(w)\right)\Big)_{\text{even}}
$$
  
=  $N(w)\Big(N(z)\log S(z) - \frac{1-n}{2}N(z)\log\left(1 - \frac{1}{2}z^{-1}\mathcal{E}_1(z)\right)\Big)_{\text{even}}$ , (6.42)

where the index "even" means that we take only the terms with even powers of  $w$ (resp.  $z$ ) in the LHS (resp. RHS). In order to complete the proof of Theorem [6.6,](#page-28-2) we are left to prove equation  $(6.42)$ . We shall prove in the following two subsections equation [\(6.42\)](#page-32-0) separately for  $n = 0$  and  $n = 1$ .

6.3. **Proof of equation** [\(6.42\)](#page-32-0) for  $n = 0$ . Setting  $x' = x$  and  $t'_k = t_k - \frac{2}{kw_1^k} - \frac{2}{kw_2^k}$ , we have, as in  $(6.28)$ ,

$$
e^{z \cdot t} e^{-z \cdot t'} = (1 + \frac{z}{w_1})(1 + \frac{z}{w_2}) \iota_{w_1} (1 - \frac{z}{w_1})^{-1} \iota_{w_2} (1 - \frac{z}{w_2})^{-1}
$$
  
=  $1 + 2 \frac{w_1 + w_2}{w_1 - w_2} \left( -w_1 \iota_{w_1} (w_1 - z)^{-1} + w_2 \iota_{w_2} (w_2 - z)^{-1} \right).$  (6.43)

<span id="page-32-2"></span><span id="page-32-1"></span>By the bilinear identity  $(5.13)$  and equation  $(6.43)$  we get

$$
\operatorname{Res}_{z} \left( S(z)e^{-2w_{1}^{-1}\cdot\tilde{\partial}_{t}}e^{-2w_{2}^{-1}\cdot\tilde{\partial}_{t}}(S^{*})^{-1}(-z) \right)
$$
\n
$$
= 2\frac{w_{1}+w_{2}}{w_{1}-w_{2}} \operatorname{Res}_{z} \left( \left( w_{1}\iota_{w_{1}}(w_{1}-z)^{-1} - w_{2}\iota_{w_{2}}(w_{2}-z)^{-1} \right) \times \right. \left. \times S(z)e^{-2w_{1}^{-1}\cdot\tilde{\partial}_{t}}e^{-2w_{2}^{-1}\cdot\tilde{\partial}_{t}}(S^{*})^{-1}(-z) \right).
$$
\n(6.44)

Recalling that  $S(z) = 1 + s_1 z^{-1} + \dots$  and  $(S^*)^{-1}(-z) = 1 - s_1 z^{-1} + \dots$ , the LHS of [\(6.44\)](#page-32-2) is

<span id="page-32-3"></span>
$$
- \left( e^{-2w_1^{-1} \cdot \tilde{\partial}_t} e^{-2w_2^{-1} \cdot \tilde{\partial}_t} - 1 \right) s_1 \,. \tag{6.45}
$$

<span id="page-32-4"></span>Using the second equation in  $(2.22)$  we can rewrite the RHS of  $(6.45)$  as

$$
2\frac{w_1 + w_2}{w_1 - w_2} \left( w_1 \left( S(z) e^{-2w_1^{-1} \cdot \tilde{\partial}_t} e^{-2w_2^{-1} \cdot \tilde{\partial}_t} (S^*)^{-1}(-z) \right)_{-} \Big|_{z=w_1} - w_2 \left( S(z) e^{-2w_1^{-1} \cdot \tilde{\partial}_t} e^{-2w_2^{-1} \cdot \tilde{\partial}_t} (S^*)^{-1}(-z) \right)_{-} \Big|_{z=w_2} \right) \tag{6.46}
$$

<span id="page-32-5"></span>The  $-$  sign means that we need to take negative powers in z before substituting  $z = w_1$  in the first parenthesis and  $z = w_2$  in the second. We next observe that

$$
\begin{aligned} &\left(S(z)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-z)\right)_-\bigg|_{z=w_1} \\ &=S(w_1)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-w_1)-1\,,\end{aligned} \tag{6.47}
$$

<span id="page-33-0"></span>and similarly

$$
\left(S(z)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-z)\right)_-\Big|_{z=w_2}
$$
\n
$$
=S(w_2)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-w_2)-1.
$$
\n(6.48)

Combining equations [\(6.44\)](#page-32-2), [\(6.45\)](#page-32-3), [\(6.46\)](#page-32-4), [\(6.47\)](#page-32-5) and [\(6.48\)](#page-33-0), we get

$$
- (e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t} - 1)s_1
$$
  
=  $2\frac{w_1 + w_2}{w_1 - w_2} (w_1S(w_1)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-w_1)$   
 $-w_2S(w_2)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}(S^*)^{-1}(-w_2)) - 2(w_1 + w_2).$ 

Using equation [\(6.26\)](#page-29-5), the above equation becomes, after some manipulations,

<span id="page-33-1"></span>
$$
w_1 S(w_1) \frac{1 - \frac{1}{2}w_1^{-1}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_1)}{e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(w_1)} - w_2 S(w_2) \frac{1 - \frac{1}{2}w_2^{-1}e^{-2w_1^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_2)}{e^{-2w_1^{-1}\cdot\tilde{\partial}_t}S(w_2)} \quad (6.49)
$$
  
= 
$$
(w_1 - w_2) \left(1 - \frac{1}{2}(w_1 + w_2)^{-1} \left(e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t} - 1\right)s_1\right).
$$

In analogy with [\(6.25\)](#page-29-4), we let

$$
\mathcal{E}_1(z, w) = \left(e^{-2z^{-1} \cdot \tilde{\partial}_t} e^{-2w^{-1} \cdot \tilde{\partial}_t} - 1\right) s_1,
$$

and we observe that

<span id="page-33-4"></span>
$$
e^{-2w_2^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_1) = \mathcal{E}_1(w_1, w_2) - \mathcal{E}_1(w_2).
$$
 (6.50)

<span id="page-33-2"></span>Hence, equation [\(6.49\)](#page-33-1) becomes

$$
\begin{split} \left(w_1 - \frac{1}{2}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_1)\right) & \frac{S(w_1)}{e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(w_1)}\\ &- \left(w_2 - \frac{1}{2}e^{-2w_1^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_2)\right) \frac{S(w_2)}{e^{-2w_1^{-1}\cdot\tilde{\partial}_t}S(w_2)}\\ & = (w_1 - w_2)\left(1 - \frac{1}{2}(w_1 + w_2)^{-1}\mathcal{E}_1(w_1, w_2)\right). \end{split} \tag{6.51}
$$

<span id="page-33-3"></span>Next, we apply  $e^{2w_2^{-1}\cdot\tilde{\partial}_t}$  to both sides of  $(6.51)$  to get

$$
(w_1 - \frac{1}{2} \mathcal{E}_1(w_1)) \frac{e^{2w_2^{-1} \cdot \tilde{\partial}_t} S(w_1)}{S(w_1)}
$$
  
–  $(e^{2w_2^{-1} \cdot \tilde{\partial}_t} S(w_2)) e^{-2w_1^{-1} \cdot \tilde{\partial}_t} e^{2w_2^{-1} \cdot \tilde{\partial}_t} \frac{w_2 - \frac{1}{2} \mathcal{E}_1(w_2)}{S(w_2)}$   
=  $(w_1 - w_2) \Big( 1 - \frac{1}{2} (w_1 + w_2)^{-1} (\mathcal{E}_1(w_1) - \mathcal{E}_1(-w_2)) \Big),$  (6.52)

where in the RHS we used the identity

$$
e^{2w_2^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_1,w_2)=e^{-2w_1^{-1}\cdot\tilde{\partial}_t}s_1-e^{2w_2^{-1}\cdot\tilde{\partial}_t}s_1=\mathcal{E}_1(w_1)-\mathcal{E}_1(-w_2).
$$

Recall that, by assumption,  $S(\partial)$  satisfies the condition  $(S^*)^{-1}(\partial) = S(\partial)$  (cf.  $(6.18)$ ). Hence, by this assumption and equation  $(6.26)$  we have

$$
e^{2w_2^{-1}\cdot \tilde{\partial}_t}S(w_2)=e^{2w_2^{-1}\cdot \tilde{\partial}_t}(S^*)^{-1}(w_2)=\frac{1+\frac{1}{2}w_2^{-1}\mathcal{E}_1(-w_2)}{S(-w_2)}.
$$

Hence,

<span id="page-34-0"></span>
$$
(e^{2w_2^{-1}\cdot\tilde{\partial}_t}S(w_2))e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{2w_2^{-1}\cdot\tilde{\partial}_t}\frac{w_2-\frac{1}{2}\mathcal{E}_1(w_2)}{S(w_2)}
$$
  
= 
$$
\frac{1+\frac{1}{2}w_2^{-1}\mathcal{E}_1(-w_2)}{S(-w_2)}e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{2w_2^{-1}\cdot\tilde{\partial}_t}(w_2e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(-w_2))
$$
 (6.53)  
= 
$$
\frac{w_2+\frac{1}{2}\mathcal{E}_1(-w_2)}{S(-w_2)}e^{-2w_1^{-1}\cdot\tilde{\partial}_t}S(-w_2).
$$

Substituting [\(6.53\)](#page-34-0) in the LHS of [\(6.52\)](#page-33-3) and replacing  $w_2$  by  $-w_2$  we arrive at the identity

<span id="page-34-1"></span>
$$
(w_1 - \frac{1}{2} \mathcal{E}_1(w_1)) \frac{e^{-2w_2^{-1} \cdot \tilde{\partial}_t} S(w_1)}{S(w_1)} + (w_2 - \frac{1}{2} \mathcal{E}_1(w_2)) \frac{e^{-2w_1^{-1} \cdot \tilde{\partial}_t} S(w_2)}{S(w_2)}
$$
  
= 
$$
(w_1 + w_2) \Big( 1 - \frac{1}{2} (w_1 - w_2)^{-1} (\mathcal{E}_1(w_1) - \mathcal{E}_1(w_2)) \Big).
$$
 (6.54)

Next, we multiply both sides of the identities  $(6.51)$  and  $(6.54)$  to get

<span id="page-34-2"></span>
$$
(w_1 - \frac{1}{2}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_1))(w_1 - \frac{1}{2}\mathcal{E}_1(w_1))
$$
  
+ 
$$
(w_1 - \frac{1}{2}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_1))(w_2 - \frac{1}{2}\mathcal{E}_1(w_2))\frac{S(w_1)}{e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(w_1)} \cdot \frac{e^{-2w_1^{-1}\cdot\tilde{\partial}_t}S(w_2)}{S(w_2)}
$$
  
- 
$$
(w_2 - \frac{1}{2}e^{-2w_1^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_2))(w_1 - \frac{1}{2}\mathcal{E}_1(w_1))\frac{e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(w_1)}{S(w_1)} \cdot \frac{S(w_2)}{e^{-2w_1^{-1}\cdot\tilde{\partial}_t}S(w_2)}
$$
  
- 
$$
(w_2 - \frac{1}{2}e^{-2w_1^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_2))(w_2 - \frac{1}{2}\mathcal{E}_1(w_2))
$$
  
= 
$$
(w_1 + w_2 - \frac{1}{2}\mathcal{E}_1(w_1, w_2))\left(w_1 - w_2 - \frac{1}{2}(\mathcal{E}_1(w_1) - \mathcal{E}_1(w_2))\right).
$$
  
(6.55)

Using [\(6.50\)](#page-33-4), it is straightforward to check that

$$
(w_1 - \frac{1}{2}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_1))(w_1 - \frac{1}{2}\mathcal{E}_1(w_1))
$$
  
–  $(w_2 - \frac{1}{2}e^{-2w_1^{-1}\cdot\tilde{\partial}_t}\mathcal{E}_1(w_2))(w_2 - \frac{1}{2}\mathcal{E}_1(w_2))$   
=  $(w_1 + w_2 - \frac{1}{2}\mathcal{E}_1(w_1, w_2)) (w_1 - w_2 - \frac{1}{2}(\mathcal{E}_1(w_1) - \mathcal{E}_1(w_2))).$ 

Then, equation [\(6.55\)](#page-34-2) becomes

<span id="page-34-3"></span>
$$
\left(\frac{S(w_1)}{e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(w_1)}\cdot\frac{e^{-2w_1^{-1}\cdot\tilde{\partial}_t}S(w_2)}{S(w_2)}\right)^2
$$
\n
$$
=\frac{1-\frac{1}{2}w_1^{-1}\mathcal{E}_1(w_1)}{e^{-2w_2^{-1}\cdot\tilde{\partial}_t}\left(1-\frac{1}{2}w_1^{-1}\mathcal{E}_1(w_1)\right)}\cdot\frac{e^{-2w_1^{-1}\cdot\tilde{\partial}_t}\left(1-\frac{1}{2}w_2^{-1}\mathcal{E}_1(w_2)\right)}{1-\frac{1}{2}w_2^{-1}\mathcal{E}_1(w_2)}\cdot\frac{(6.56)}
$$

<span id="page-34-4"></span>Taking logarithm of both sides of [\(6.56\)](#page-34-3) and dividing by 2 we arrive at

$$
\left(e^{-2w_1^{-1}\cdot\tilde{\partial}_t}-1\right)\left(\log S(w_2)-\frac{1}{2}\log\left(1-\frac{1}{2}w_2^{-1}\mathcal{E}_1(w_2)\right)\right)
$$
\n
$$
=\left(e^{-2w_2^{-1}\cdot\tilde{\partial}_t}-1\right)\left(\log S(w_1)-\frac{1}{2}\log\left(1-\frac{1}{2}w_1^{-1}\mathcal{E}_1(w_1)\right)\right).
$$
\n(6.57)

<span id="page-35-0"></span>Finally, we apply  $N(w_1)N(w_2)$  to both sides of  $(6.57)$  and use  $(6.38)$  to get

$$
N(w_1)N(w_2)\left(\log S(w_2) - \frac{1}{2}\log\left(1 - \frac{1}{2}w_2^{-1}\mathcal{E}_1(w_2)\right)\right)
$$
  
=  $N(w_1)N(w_2)\left(\log S(w_1) - \frac{1}{2}\log\left(1 - \frac{1}{2}w_1^{-1}\mathcal{E}_1(w_1)\right)\right)$ . (6.58)

Equation [\(6.58\)](#page-35-0) implies [\(6.42\)](#page-32-0) for  $n = 0$ .

6.4. Proof of equation [\(6.42\)](#page-32-0) for  $n = 1$ . In [\[DJKM83\]](#page-37-5) the following bilinear identity for the BKP hierarchy is used

<span id="page-35-1"></span>
$$
\text{Res}_{z} (w(x, t; , z)w(x', t', -z)z^{-1}) = 1.
$$
\n(6.59)

**Lemma 6.11.** Let us assume that the bilinear identity for the BKP  $(6.59)$  holds. Then  $S(\partial)$  satisfies the condition [\(6.18\)](#page-28-0) with  $n = 1$  and the bilinear identity [\(5.13\)](#page-22-3) holds. Conversely, assume that  $S(\partial)$  satisfies the condition [\(6.18\)](#page-28-0) with  $n = 1$  and  $w(z) = S(z)e^{z \cdot t}$  satisfies the the bilinear identity [\(5.13\)](#page-22-3). Then  $w(z)$  satisfies the bilinear identity for the BKP  $(6.59)$ .

*Proof.* Let us apply  $\partial_x^m$ ,  $m \ge 0$ , to [\(6.59\)](#page-35-1) and set  $x' = x$  and  $t'_k = t_k$ ,  $k \in \mathcal{Z}_L$ . Then, we have

<span id="page-35-2"></span>
$$
\text{Res}_{z} ((z + \partial)^{m} S(z))S(-z)z^{-1}) = \delta_{m,0}.
$$
 (6.60)

Similarly to the proof of Lemma [2.9,](#page-7-4) with  $A(\partial) = S(\partial)$  and  $B^*(\partial) = -S(\partial)\partial^{-1}$ , the identities [\(6.60\)](#page-35-2) for  $m > 0$  imply  $S(\partial) \partial^{-1} \circ S^*(\partial) \in \mathcal{F}[\partial] \partial^{-1}$ . Since  $S(\partial) \in$  $1 + \mathcal{F}[[\partial^{-1}]]\partial^{-1}$  we then have  $S(\partial)\partial^{-1} \circ S^*(\partial) = a\partial^{-1}$ , for some  $a \in \mathcal{F}$ . From [\(6.60\)](#page-35-2) with  $m = 0$  we get that  $a = 1$  thus showing that  $S(\partial)$  satisfies the condition  $(6.18)$  with  $n = 1$ . In particular, we have

<span id="page-35-3"></span>
$$
(S^*)^{-1}(-z) = -(-z+\partial)S(-z)z^{-1}.
$$
\n(6.61)

Next, let us apply  $\partial_{x'}$  to both sides of [\(6.59\)](#page-35-1). We get

$$
0 = \text{Res}_{z} (w(x, t; z) \partial_{x'} w(x', t', -z) z^{-1})
$$
  
= Res\_{z} (w(x, t; z)(-z + \partial\_{x'}) S(-z) z^{-1} e^{-zx'-z \cdot t'})  
= - \text{Res}\_{z} (w(x, t; z) w^\*(x', t', z)).

In the last identity we used  $(6.61)$ . This shows that  $w(z)$  satisfies  $(5.13)$ .

Conversely, let us assume that  $S(\partial)$  satisfies [\(6.61\)](#page-35-3) (that is condition [\(6.18\)](#page-28-0) with  $n = 1$ ) and  $w(z) = S(z)e^{z \cdot t}$  satisfies the bilinear identity [\(5.13\)](#page-22-3), and let

<span id="page-35-5"></span>
$$
f = \text{Res}_{z}(w(x, t, z)w(x', t', -z)z^{-1}).
$$
\n(6.62)

By  $(6.61)$ , we have

$$
w^*(z) = (S^*)^{-1}(-z)e^{-z\cdot t} = -(-z+\partial)S(-z)z^{-1}e^{-z\cdot t} = -\partial w(-z)z^{-1}.
$$

Hence,

<span id="page-35-4"></span>
$$
\partial_{x'}f = -\operatorname{Res}_{z}(w(x, t, z)\partial_{x'}w(x', t', -z)z^{-1}) = -\operatorname{Res}_{z}(w(x, t, z)w^{\star}(x', t', z)) = 0.
$$
\n(6.63)

Moreover, by Corollary [\(5.6\)](#page-23-1) and [\(5.10\)](#page-22-2) we have  $(k \in \mathcal{Z}_L)$ 

$$
\frac{\partial f}{\partial t'_k} = \text{Res}_z \left( w(x, t, z) \frac{\partial w(x', t', -z)}{\partial t'_k} z^{-1} \right)
$$
  
= Res\_z (w(x, t, z)B\_k(\partial\_{x'})w(x', t', -z)z^{-1}) = B\_k(\partial\_{x'})f = 0.

In the last identity we used the fact that  $B_k$  has no constant term (cf. Remark [3.13](#page-15-3)) and [\(3.16\)](#page-15-0)) and [\(6.63\)](#page-35-4). This shows that f does not depend on  $x'$  and  $t'_{k}$ , thus we can set  $x' = x$  and  $t'_{k} = t_{k}$  in its definition [\(6.62\)](#page-35-5) to get

$$
f = f|_{x'=x,t'_{k}=t_{k}} = \text{Res}_{z}(S(z)S(-z)z^{-1}) = 1.
$$

This concludes the proof.  $\Box$ 

In analogy with the proof of Lemma [6.9](#page-0-2) setting  $x' = x$  and  $t'_{k} = t_{k} - \frac{2}{kw^{k}}$  in [\(6.59\)](#page-35-1) we get, using [\(6.28\)](#page-30-6)

Res<sub>z</sub> 
$$
((1 + \frac{w}{z})\iota_w(w - z)^{-1}S(z)e^{-2w^{-1}\cdot\tilde{\partial}_t}S(-z)) = 1,
$$

which, by the second equation in  $(2.22)$ , gives

<span id="page-36-0"></span>
$$
\left( (1 + \frac{w}{z}) S(z) e^{-2w^{-1} \cdot \tilde{\partial}_t} S(-z) \right)_{-}\Big|_{z=w} = 1.
$$
 (6.64)

<span id="page-36-1"></span>Note that the LHS of  $(6.64)$  can be rewritten as

$$
2S(w)e^{-2w^{-1}\cdot\tilde{\partial}_t}S(-w) - \left((1+\frac{w}{z})S(z)e^{-2w^{-1}\cdot\tilde{\partial}_t}S(-z)\right)_+\bigg|_{z=w}
$$
\n
$$
= 2S(w)e^{-2w^{-1}\cdot\tilde{\partial}_t}S(-w) - 1.
$$
\n(6.65)

Combining equations [\(6.64\)](#page-36-0) and [\(6.65\)](#page-36-1) we get

<span id="page-36-3"></span>
$$
S(w)e^{-2w^{-1}\cdot\tilde{\partial}_t}S(-w) = 1.
$$
\n(6.66)

Next, setting  $x' = x$  and  $t'_{k} = t_{k} - \frac{2}{kw_{1}^{k}} - \frac{2}{kw_{2}^{k}}$  in [\(6.59\)](#page-35-1), we have, using [\(6.43\)](#page-32-1),

$$
\operatorname{Res}_{z} \left( S(z)e^{-2w_{1}^{-1}\cdot\tilde{\partial}_{t}}e^{-2w_{2}^{-1}\cdot\tilde{\partial}_{t}}S(-z)z^{-1} \right)
$$
\n
$$
= 1 + 2\frac{w_{1} + w_{2}}{w_{1} - w_{2}} \operatorname{Res}_{z} \left( \left( w_{1}\iota_{w_{1}}(w_{1} - z)^{-1} - w_{2}\iota_{w_{2}}(w_{2} - z)^{-1} \right) \times \\ \times S(z)e^{-2w_{1}^{-1}\cdot\tilde{\partial}_{t}}e^{-2w_{2}^{-1}\cdot\tilde{\partial}_{t}}S(-z)z^{-1} \right).
$$

Since  $\text{Res}_{z}\left(S(z)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(-z)z^{-1}\right)=1$ , the above equation becomes

$$
\operatorname{Res}_{z} \left( \left( w_{1} \iota_{w_{1}} (w_{1} - z)^{-1} - w_{2} \iota_{w_{2}} (w_{2} - z)^{-1} \right) \times S(z) e^{-2w_{1}^{-1} \cdot \tilde{\partial}_{t}} e^{-2w_{2}^{-1} \cdot \tilde{\partial}_{t}} S(-z) z^{-1} \right) = 0.
$$
\n(6.67)

<span id="page-36-2"></span>Using the second equation in  $(2.22)$  we can rewrite  $(6.67)$  as

$$
w_1 \left( S(z)e^{-2w_1^{-1}\cdot\tilde{\partial}_t} e^{-2w_2^{-1}\cdot\tilde{\partial}_t} S(-z)z^{-1} \right)_- \Big|_{z=w_1}
$$
  
= 
$$
w_2 \left( S(z)e^{-2w_1^{-1}\cdot\tilde{\partial}_t} e^{-2w_2^{-1}\cdot\tilde{\partial}_t} S(-z)z^{-1} \right)_- \Big|_{z=w_2}.
$$
 (6.68)

The  $-$  sign means that we need to take negative powers in z before substituting  $z = w_1$  in the first parenthesis and  $z = w_2$  in the second. We next observe that both terms in parenthesis belong to  $\mathcal{F}[[z^{-1}]]z^{-1}$ , hence we can remove the – sign and get

$$
S(w_1)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(-w_1)=S(w_2)e^{-2w_1^{-1}\cdot\tilde{\partial}_t}e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(-w_2),
$$

which, using  $(6.66)$ , can be rewritten as

<span id="page-36-4"></span>
$$
\frac{S(w_1)}{e^{-2w_2^{-1}\cdot\tilde{\partial}_t}S(w_1)} = \frac{S(w_2)}{e^{-2w_1^{-1}\cdot\tilde{\partial}_t}S(w_2)}.
$$
\n(6.69)

Taking logarithm of both sides of  $(6.69)$  and applying  $N(w_1)N(w_2)$  equation  $(6.42)$ for  $n = 1$  follows in the same way as in the proof of Lemma [6.5\(](#page-27-5)b).

#### <span id="page-37-0"></span>**REFERENCES**

- <span id="page-37-22"></span> $[Ad179]$  Adler, M. On a trace functional for formal pseudo differential operators and the symplectic structure of the Korteweg-de Vries type equations. Invent. Math. 50 (1978/79), no. 3, 219-248.
- <span id="page-37-20"></span>[BDSK09] Barakat A., De Sole A., Kac V., Poisson vertex algebras in the theory of Hamiltonian equations. Jpn. J. Math.  $4$  (2009), no. 2, 141-252.
- <span id="page-37-10"></span>[CDSKVvdL20] Carpentier S., De Sole A., Kac V.G., Valeri D., van de Leur J., p-reduced multicomponent KP hierarchy and classical W-algebras  $W(\mathfrak{gl}_N, p)$ , Comm. Math. Phys. 380 (2020), no. 2, 655-722.
- <span id="page-37-18"></span>[CW13] Chang L., Wu C.-Z., Tau function of the CKP hierarchy and nonlinearizable Virasoro symmetries, Nonlinearity 26 (2013), no. 9, 2577-2596.
- <span id="page-37-12"></span>[Che92] Cheng, Y., Constraints of the Kadomtsev-Petviashvili hierarchy, J. Math. Phys. 33 (1992), no. 11, 3774-3782.
- <span id="page-37-4"></span>[DJKM81] Date E., Kashiwara M., Jimbo M., Miwa T., KP hierarchies of orthogonal and symplectic type - transformation groups for soliton equations VI, Journal of the Physical Society of Jaan, 50 (1981), no.11, 3813-3818.
- <span id="page-37-5"></span>[DJKM83] Date E., Jimbo M., Kashiwara M., Miwa T., Transformation groups for soliton equations, Nonlinear integrable systems-classical theory and quantum theory (Kyoto, 1981), 39-119, World Sci. Publishing, Singapore, 1983.
- <span id="page-37-7"></span>[DSKV15] De Sole A., Kac V. G., Valeri D., Adler-Gelfand-Dickey approach to classical Walgebras within the theory of Poisson vertex algebras. Int. Math. Res. Not. IMRN (2015), 21, 11186-11235.
- <span id="page-37-21"></span>[DSKV16] De Sole A., Kac V. G., Valeri D., A new scheme of integrability for (bi)-Hamiltonian PDE, Comm. Math. Phys.  $347$  (2016), no.2, 449-488.
- <span id="page-37-8"></span>[DSKV16a] De Sole A., Kac V. G., Valeri D., Classical affine W-algebras for  $\mathfrak{gl}_N$  and associated integrable Hamiltonian hierarchies, Comm. Math. Phys.  $348$  (2016), no.1, 265-319.
- <span id="page-37-9"></span>[DSKV18] De Sole A., Kac V., Valeri D., Classical affine W-algebras and the associated integrable Hamiltonian hierarchies for classical Lie algebras, Comm. Math. Phys. 360 (2018), no.3, 851-918.
- <span id="page-37-17"></span>[Dic03] Dickey L.A., Soliton equations and Hamiltonian systems, Second edition, Advanced Series in Mathematical Physics 26, World Scientific Publishing Co., Inc., River Edge, NJ, 2003.
- <span id="page-37-26"></span>[DMH09] Dimakis A., Müller-Hoissen F., BKP and CKP revisited: the odd KP system. Inverse Problems 25 (2009), no. 4, 045001, 33 pp.
- <span id="page-37-6"></span>[DS85] Drinfeld V.G, Sokolov V.V, Lie algebras and equations of Korteweg-de Vries type, Soviet J. Math. 30 (1985), 1975-2036.
- <span id="page-37-2"></span>[GD76] Gelfand I.M., Dickey L.A., Fractional powers of operators, and Hamiltonian systems, Funkcional. Anal. i Prilozen. 10 (1976), no. 4, 13-29.
- [JM83] Jimbo M., Miwa T., Solitons and infinite-dimensional Lie algebras. Publ. Res. Inst. Math. Sci. 19 (1983), no. 3, 943-1001.
- <span id="page-37-23"></span>[Kaup80] Kaup, D.J. On the inverse scattering problem for cubic eigenvalue problems of the class  $\psi_{xxx} + 6Q\psi_x + 6R\psi = \lambda\psi$ . Stud. Appl. Math. 62 (1980), no. 3, 189-216.
- <span id="page-37-11"></span>[KSS91] Konopelchenko B., Sidorenko J., Strampp W., (1+1)-dimensional integrable systems as symmetry constraints of  $(2+1)$ -dimensional system, Phys. Lett. A 157 (1991), 17-21.
- <span id="page-37-13"></span>[KS92] Konopelchenko B., Strampp W., New reductions of the Kadomtsev-Petviashvili and twodimensional Toda lattice hierarchies via symmetry constraints, J. Math. Phys. 33 (1992), no. 11, 3676-3686.
- <span id="page-37-25"></span>[KN80] Krichever, I.M., Novikov, S.P. Holomorphic bundles over algebraic curves, and nonlinear equations. (Russian) Uspekhi Mat. Nauk 35 (1980), no. 6 (216), 47-68, 215.
- <span id="page-37-19"></span>[KZ20] Krichever I., Zabrodin A., Kadomtsev-Petviashvili turning points and CKP hierarchy, [arXiv:2012.04482.](https://arxiv.org/abs/2012.04482)
- <span id="page-37-1"></span>[Lax68] Lax P., Integrals of nonlinear equations of evolution and solitary waves. Comm. Pure Appl. Math. 21 (1968), 467-490.
- <span id="page-37-15"></span>[Ma81] Ma Y.-C., The resonant interaction among long and short waves, Wave motion 3 (1981), 257-267.
- <span id="page-37-3"></span>[Sat81] Sato M., Soliton equations as dynamical systems on an infinite dimensional Grassmann manifold, Proc. Symp., Kyoto 1981, RIMS Kokyuroku 439 (1981), 30-46.
- <span id="page-37-24"></span> $[SK74]$  Sawada K., Kotera T., A method for finding N-soliton solutions of the K.d.V. equation and  $K.d. V.$ -like equation. Progr. Theoret. Phys.  $51$  (1974), 1355-1367.
- <span id="page-37-16"></span>[Shi86] Shiota T., Characterization of Jacobian varieties in terms of soliton equations. Invent. Math. 83 (1986) 333-382.
- <span id="page-37-14"></span>[SS93] Sidorenko J., Strampp W., Multicomponent integrable reductions in the Kadomtsev-Petviashvili hierarchy, J. Math. Phys. 34 (1993), no. 4, 1429-1446.
- <span id="page-38-2"></span>[Sok84] Sokolov, V. V. Hamiltonian property of the Krichever-Novikov equation. (Russian) Dokl. Akad. Nauk SSSR 277 (1984), no. 1, 48-50.
- <span id="page-38-0"></span>[Y076] Yajima N., Oikawa M., Formation and interaction of sonic-Langmuir solitons - inverse scattering method, Progr. Theoret. Phys. 56 (1976), no. 6, 1719-1739.
- <span id="page-38-1"></span>[ZS74] Zakharov, V. E.; Shabat, A. B. A plan for integrating the nonlinear equations of mathematical physics by the method of the inverse scattering problem. I. (Russian) Funkcional. Anal. i Prilozen. 8 (1974), no. 3, 43-53.

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