

PSZ2 G091: A massive double cluster at $z \sim 0.822$ observed by the NIKA2 camera

E. Artis^{1,*}, *R. Adam*², *P. Ade*³, *H. Ajeddig*⁴, *P. André*⁴, *M. Arnaud*⁴, *H. Aussel*⁴, *I. Bartalucci*⁷, *A. Beelen*⁶, *A. Benoît*⁸, *S. Berta*⁹, *L. Bing*¹⁰, *O. Bourrion*¹, *M. Calvo*⁸, *A. Catalano*¹, *M. De Petris*¹¹, *F.-X. Désert*¹², *S. Doyle*³, *E. F. C. Driessen*⁹, *A. Ferragamo*¹¹, *A. Gomez*¹³, *J. Goupy*⁸, *F. Kéruzoré*¹, *C. Kramer*¹⁴, *B. Ladjelate*¹⁴, *G. Lagache*¹⁰, *S. Leclercq*⁹, *J.-F. Lestrade*¹⁵, *J.-F. Macías-Pérez*¹, *A. Maury*⁴, *P. Mauskopf*^{3,16}, *F. Mayet*¹, *A. Monfardini*⁸, *M. Muñoz-Echeverría*¹, *A. Paliwal*¹¹, *L. Perotto*¹, *G. Pisano*³, *E. Pointecouteau*⁵, *N. Ponthieu*¹², *G. W. Pratt*⁴, *V. Revéret*⁴, *A. J. Rigby*³, *A. Ritacco*^{6,17}, *C. Romero*¹⁸, *H. Roussel*¹⁹, *F. Ruppin*²¹, *K. Schuster*⁹, *S. Shu*²², *A. Sievers*¹⁴, *C. Tucker*³, and *G. Yepes*²⁰

¹Univ. Grenoble Alpes, CNRS, Grenoble INP, LPSC-IN2P3, 53, avenue des Martyrs, 38000 Grenoble, France

²LLR (Laboratoire Leprince-Ringuet), CNRS, École Polytechnique, Institut Polytechnique de Paris, Palaiseau, France

³School of Physics and Astronomy, Cardiff University, Queen's Buildings, The Parade, Cardiff, CF24 3AA, UK

⁴AIM, CEA, CNRS, Université Paris-Saclay, Université Paris Diderot, Sorbonne Paris Cité, 91191 Gif-sur-Yvette, France

⁵Univ. de Toulouse, UPS-OMP, Institut de Recherche en Astrophysique et Planétologie, Toulouse, France CNRS, IRAP, 9 Av. colonel Roche, BP 44346, F-31028 Toulouse cedex 4, France

⁶Institut d'Astrophysique Spatiale (IAS), CNRS, Université Paris Sud, Orsay, France

⁷INAF, IASF-Milano, Via A. Corti 12, 20133 Milano, Italy

⁸Institut Néel, CNRS, Université Grenoble Alpes, France

⁹Institut de RadioAstronomie Millimétrique (IRAM), Grenoble, France

¹⁰Aix Marseille Univ, CNRS, CNES, LAM (Laboratoire d'Astrophysique de Marseille), Marseille, France

¹¹Dipartimento di Fisica, Sapienza Università di Roma, Piazzale Aldo Moro 5, I-00185 Roma, Italy

¹²Univ. Grenoble Alpes, CNRS, IPAG, 38000 Grenoble, France

¹³Centro de Astrobiología (CSIC-INTA), Torrejón de Ardoz, 28850 Madrid, Spain

¹⁴Instituto de Radioastronomía Milimétrica (IRAM), Granada, Spain

¹⁵LERMA, Observatoire de Paris, PSL Research University, CNRS, Sorbonne Université, UPMC, 75014 Paris, France

¹⁶School of Earth and Space Exploration and Department of Physics, Arizona State University, Tempe, AZ 85287, USA

¹⁷Laboratoire de Physique de l'École Normale Supérieure, ENS, PSL Research University, CNRS, Sorbonne Université, Université de Paris, 75005 Paris, France

¹⁸Department of Physics and Astronomy, University of Pennsylvania, 209 South 33rd Street, Philadelphia, PA, 19104, USA

¹⁹Institut d'Astrophysique de Paris, CNRS (UMR7095), 98 bis boulevard Arago, 75014 Paris, France

²⁰Departamento de Física Teórica and CIAFF, Facultad de Ciencias, Modulo 8, Universidad Autónoma de Madrid, 28049 Madrid, Spain

*e-mail: emmanuel.artis@lpsc.in2p3.fr

²¹Kavli Institute for Astrophysics and Space Research, Massachusetts Institute of Technology, Cambridge, MA 02139, USA

²²Caltech, Pasadena, CA 91125, USA

Abstract. PSZ2 G091.83+26.11 is a massive galaxy cluster with $M_{500} = 7.43 \times 10^{14} M_{\odot}$ at $z = 0.822$. This object exhibits a complex morphology with a clear bimodality observed in X-rays. However, it was detected and analysed in the *Planck* sample as a single, spherical cluster following a universal profile [1]. This model can lead to miscalculations of thermodynamical quantities, like the pressure profile. As future multiwavelength cluster experiments will detect more and more objects at higher redshifts (where we expect the fraction of merging objects to be higher), it is crucial to quantify this systematic effect. In this work, we use high-resolution observations of PSZ2 G091.83+26.11 by the NIKA2 camera to integrate the morphological characteristics of the cluster in our modelling. This is achieved by fitting a two-halo model to the SZ image and then by reconstruction of the resulting projected pressure profile. We then compare these results with the spherical assumption.

1 Introduction

Galaxy clusters are the largest virialized objects in the universe. As such, they can constrain cosmological models, fundamental physics, and astrophysics [2, 3]. Precisely, the number of clusters per unit of mass and volume, modelled as the halo mass function (historically from [4]), constitutes a robust cosmological probe, measuring among others the matter density parameter Ω_m , the rms linear fluctuations on scales of $8 h^{-1} \text{Mpc}$, σ_8 , the dark energy equation of state parameters, as well as the sum of the mass of the neutrinos [5, 6]. Cluster counts have been used by many surveys like *Planck* [7], *DES* [8], and *XXL* [9] to cite a few examples.

However, the total mass of dark matter halos hosting clusters is not an observable quantity, and must be inferred from different physical phenomena [10], like the thermal Sunyaev-Zel'dovich effect (tSZ, [11]). As a consequence, astrophysical systematic effects are biasing our measurements [12], and their effects must be integrated in any cosmological analysis. Moreover, given the fact that the size of the catalogs will be increased by several orders of magnitude in the future, effects that are now neglected will play a crucial role [13, 14].

The Large program SZ of the NIKA2 experiment (LPSZ, [15]) aims at investigating these systematic effects, taking advantage of the spatial resolution and FoV of the NIKA2 camera. In this work, we focus on the impact of cluster morphology on the reconstruction of observable quantities in the case of the massive cluster PSZ2G091.

2 The case of PSZ2G091

Table 1: Characteristics of PSZ2 G091.83+26.11 as described in [16], with the observation length and SNR at SZ peak.

z	M_{500}	θ_{500}	$t_{\text{obs}}/t_{\text{LPSZ}}$	tSZ decrement peak
0.822	$7.43 \times 10^{14} M_{\odot}$	2.2 arcmin	2.5h/2.5h=1	14.9 σ

As part of the LPSZ, PSZ2G091 was observed in October, 2017, with an average elevation of 58.5° and an average atmospheric opacity at 225 GHz of 0.243. These conditions are standard

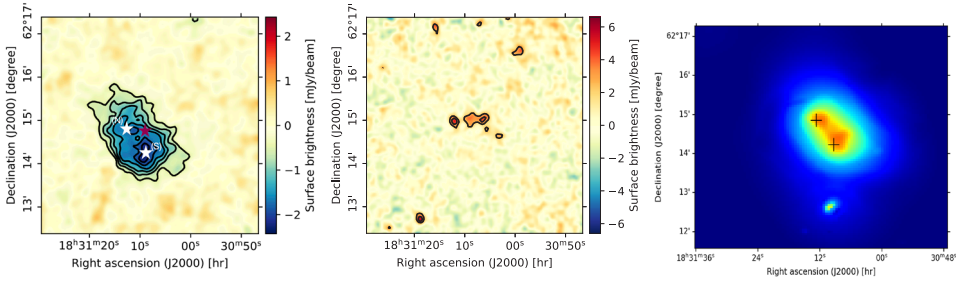


Figure 1. **Left:** Surface brightness at 2 mm from NIKA2 observations. The black contours represent the SNR levels starting at 3σ increasing by 2σ at each step. The northern (N) and southern (S) X-ray peaks are shown as white stars, the X-ray centroid being in purple. **Middle:** Surface brightness at 1 mm. The signal exhibits several point sources. Their treatment is explained in section 3. **Right:** X-ray surface brightness of PSZ2G091 from XMM-Newton. The cluster exhibits two peaks at the position $(\alpha_N, \delta_N) = (277.80, 62.24)$, and $(\alpha_S, \delta_S) = (277.79, 62.237)$, shown with the black crosses.

for observations at the IRAM 30 m telescope, at this season.

The data from NIKA2 were reduced following the pipeline described in [17]. In figure 1, we show the results of the data reduction at 1 and 2 mm. The cluster can clearly be seen in the 2 mm image, and it is clearly elongated in the NW-SW direction. There is a clear departure from sphericity, and a hint of bimodality later confirmed in the X-ray surface brightness. Conversely, in the 1 mm map, we only detect point sources, since the SZ signal is completely dominated by the noise as expected. The peaks in the X-ray map are in good agreement with the one observed in the NIKA2 2 mm map. This would imply the presence of two well-defined sub-halos in the first stages of a major merger.

3 Point sources treatment

The 1 mm map of PSZ2G091 exhibits 7 point sources (see figure 1). Their fluxes must be inferred at 2 mm, in order to fit the ICM signal. If external information on these sources is available, we can usually use SED models to infer a PDF of the flux at 2 mm for each source, and use it as a prior for the model reconstruction. However, in this case, no additional data was available. The point sources observed at this wavelength are either millimetric sources, or radio objects. For the latter, the flux expected at 2 mm is higher than the one seen in the 1 mm map. Given the value of the SZ signal in the 2 mm map, we were consequently able to identify the sources as millimetric sources. Thus we decided to use flat priors for the fluxes of the sources at 2 mm, with a conservative upper limit of $F_{2,\text{upper}} = F_1/3$. The fluxes of the point sources were then fitted jointly with the parameters of the profiles, using these flat priors.

4 Imaging analysis

4.1 Single spherical model

As a first test case, we consider a single spherical halo centred on the X-ray centroid coordinates. A forward modelling approach is incorporated in an MCMC sampling framework to

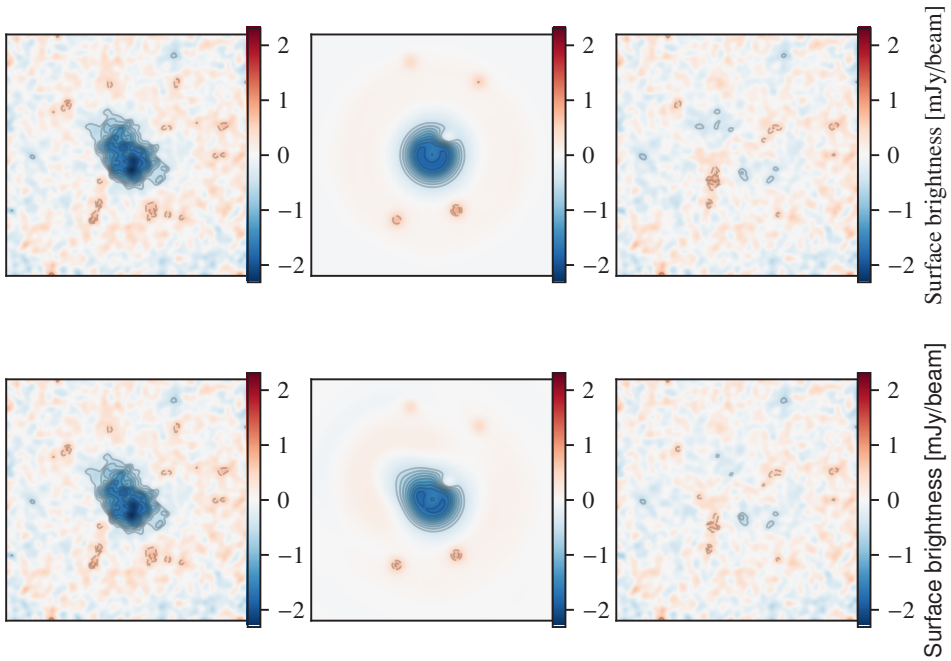


Figure 2. Single spherical model fit of the NIKA2 2 mm map for PSZ2G091 centered on the X-ray centroid coordinates (**top**), and 2-halo model (**bottom**). From **left** to **right**, we display the data, the single halo model, with the point sources treated as described in section 3, and the residuals. The maps are given in a $5' \times 5'$ area, and for display purposes, each map is smoothed with a gaussian kernel. The contours are showing the SNR level sets, starting at 3σ and increasing with a step of 1σ .

fit the parameters of the pressure profile, as well as the point sources, as described in section 3, with the collaboration software PANCO2. Instead of considering a generalised Navarro-Frenk-and-White profile [18], we use a power law model, where the pressure is divided into 6 bins, equally distributed in \log -space. Each one of these bins follow the identity

$$P(r) = P_i(r/r_i)^{-\alpha_i}. \quad (1)$$

The top row of figure 2 shows the results of the fitting procedure. Although returning reasonable residuals, the spherical symmetry clearly does not encapsulate the bimodal nature of the cluster. It is then required to improve our modelling.

4.2 Two-halo model

Instead of considering a single pressure profile, we jointly fit two halos at the positions of the X-ray peaks. The top row of figure 2 shows the results of the MCMC fitting. The results of the fits are shown on the bottom row of figure 2. It is clear that the two halo model yields a more realistic representation of the dynamical state of the cluster. Additionally, the residuals are slightly improved in the region of the northern subhalo. Of course, due to the non-spherical nature of this cluster, it is not possible to consider a radial pressure profile. However, in section 5, we describe how we recover an average radial profile for the two-halo model.

5 Pressure profile reconstruction

Our model takes into account the specific morphology of this cluster. However, the thermodynamical profiles are usually considered with the goal of reconstructing a mass profile, using the hydrostatic equilibrium equation. This requires the presence of a 1D pressure profile, when we previously fitted a 2D map. Thus, we use the following procedure to recover an

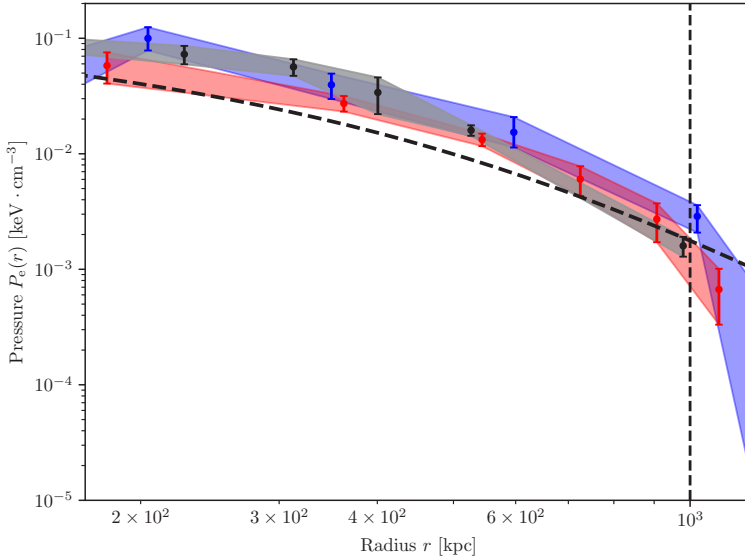


Figure 3. In red, composite pressure profile obtained by combining the northern (N) and southern (S) profiles, as described in equation 2. In blue, profile obtained with the single spherical model. In black, we represent the contours obtained from XMM X-ray data only. The black dashed line represents the universal profile from [19]. The vertical dashed line represents the radius R_{500} obtained from X-ray data assuming spherical symmetry. Although the single spherical model seems to slightly overestimate the pressure, both profile look compatible for the considered radius range.

average 1D profile. We consider both the profiles of the two halos fitted in section 4.2. With the assumption of the two sub-halos lying in the same plane perpendicular to the line of sight, we recover a mean radial pressure profile by integrating both the profiles in annuli centered around the X-ray centroid coordinates. In this framework, the pressure in the i -th bin reads:

$$P_i = \left(\int_{r_i}^{r_{i+1}} (P_N(r) + P_S(r)) 2\pi r dr \right) / (\pi(r_{i+1}^2 - r_i^2)), \tag{2}$$

where P_i is the value of the pressure, P_N and P_S are respectively the profiles of the northern and southern subhalos, r_i and r_{i+1} being the inner and outer radii of the annulus. This allows us to reconstruct the profile available in figure 3. The integration is performed up to the radius R_{500} that is found with the X-ray data. The 1D and average 2D models seem to be slightly compatible, which tends to show that even in the case of an highly disturbed cluster, the spherical model still yields robust results.

6 Conclusions

This analysis shows the challenges related to the complexity of cluster morphologies. In this work, we showed that for the case of the highly disturbed cluster PSZ2 G091, taking into account the merging state of the cluster yields results compatible with the spherical profile. This is a promising result, as the pressure profile impacts the $Y_{500} - M$ relation. To complete this analysis, we plan to perform a full thermodynamical analysis of this cluster, recovering the 2D maps of physical quantities like the temperature and the entropy. This a precondition to assess the impact of the morphology of this cluster on its full mass reconstruction, and generally on cosmological inference using clusters.

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