

A dynamical system approach to a class of radial weighted fully nonlinear equations

Liliane Maia, Gabrielle Nornberg & Filomena Pacella

To cite this article: Liliane Maia, Gabrielle Nornberg & Filomena Pacella (2021) A dynamical system approach to a class of radial weighted fully nonlinear equations, Communications in Partial Differential Equations, 46:4, 573-610, DOI: [10.1080/03605302.2020.1849281](https://doi.org/10.1080/03605302.2020.1849281)

To link to this article: <https://doi.org/10.1080/03605302.2020.1849281>



Published online: 10 Dec 2020.



Submit your article to this journal [↗](#)



Article views: 63



View related articles [↗](#)



View Crossmark data [↗](#)



A dynamical system approach to a class of radial weighted fully nonlinear equations

Liliane Maia^a , Gabrielle Nornberg^b , and Filomena Pacella^c 

^aUniversidade de Brasília, Brasília, Brazil; ^bInstituto de Ciências Matemáticas e de Computação, Universidade de São Paulo, São Carlos, Brazil; ^cSapienza Università di Roma, Roma, Italy

ABSTRACT

In this paper we study existence, nonexistence and classification of radial positive solutions of some weighted fully nonlinear equations involving Pucci extremal operators. Our results are entirely based on the analysis of the dynamics induced by an autonomous quadratic system which is obtained after a suitable transformation. This method allows to treat both regular and singular solutions in a unified way, without using energy arguments. In particular we recover known results on regular solutions for the fully nonlinear non weighted problem by alternative proofs. We also slightly improve the classification of the solutions for the extremal operator \mathcal{M}^- .

ARTICLE HISTORY

Received 23 June 2020
Accepted 1 November 2020

KEYWORDS

Critical exponents; dynamical system; fully nonlinear equations; regular and singular solutions

MSC2020

35J15; 35J60; 35B09; 34A34

1. Introduction and main results

In this paper we study positive radial solutions of the following class of fully nonlinear elliptic equations

$$\mathcal{M}_{\lambda, \Lambda}^{\pm}(D^2u) + |x|^a u^p = 0, \quad u > 0 \text{ in } \Omega, \quad (1.1)$$

where $a > -1$, $p > 1$, and \mathcal{M}^{\pm} are the Pucci's extremal operators which play an essential role in stochastic control theory and mean field games. Here, $0 < \lambda \leq \Lambda$ are the ellipticity constants, see Section 2 and [1] for their properties. The set $\Omega \in \mathbb{R}^N$, $N \geq 3$, is a radial domain such as \mathbb{R}^N , a ball B_R of radius $R > 0$ centered at the origin, the exterior of B_R , or an annulus. However, most of the time Ω will be either the whole space or a ball.

We deal with both regular and singular solutions u of (1.1) which are C^2 for $r > 0$. In the singular case Ω will be either $\mathbb{R}^N \setminus \{0\}$ or $B_R \setminus \{0\}$, and we assume the condition

$$\lim_{r \rightarrow 0} u(r) = +\infty, \quad r = |x|. \quad (1.2)$$

Finally, whenever Ω has a boundary, we prescribe the Dirichlet condition

$$u = 0 \text{ on } \partial\Omega, \quad \text{or} \quad u = 0 \text{ on } \partial\Omega \setminus \{0\} \text{ under (1.2)}. \quad (1.3)$$

Let us recall some previous results when $a = 0$. A general existence result in bounded domains Ω (not necessarily radial) was obtained in [2] under the condition

$$1 < p \leq \frac{\tilde{N}_+}{\tilde{N}_+ - 2} \text{ for } \mathcal{M}_{\lambda, \Lambda}^+, \quad 1 < p \leq \frac{\tilde{N}_-}{\tilde{N}_- - 2} \text{ for } \mathcal{M}_{\lambda, \Lambda}^- \quad (1.4)$$

if $\tilde{N}_+ > 2$, where \tilde{N}_\pm are the so called dimensional-like numbers

$$\tilde{N}_+ = \frac{\lambda}{\Lambda}(N-1) + 1, \quad \tilde{N}_- = \frac{\Lambda}{\lambda}(N-1) + 1. \quad (1.5)$$

The intervals in (1.4) describe the optimal range for existence of supersolutions to \mathcal{M}^\pm , as shown in [3]. Some extensions and related results can be found in [4, 5].

When the Pucci's operators reduce to the Laplacian (i.e. for $\lambda = \Lambda$), both exponents (1.4) are equal to $\frac{N}{N-2}$ which is known as Serrin exponent. They do not provide optimal bounds in terms of solutions of (1.1), as it is clear for instance by considering the semilinear case.

Nevertheless, as far as the radial setting is concerned, critical exponents which represent the threshold for the existence of solutions to (1.1) can be defined. They were introduced for $a=0$ by Felmer and Quaas in [6] in order to study existence and classification of radial positive solutions in \mathbb{R}^N . These are also the watershed for existence and nonexistence of positive solutions in the ball. Note that every positive solution in the ball when $a=0$ is radial, by [7], while this is not true in general for $a \neq 0$, even in the semilinear case.

When $\lambda = \Lambda$ the corresponding critical exponents are the same, both in radial and nonradial settings; see [8] for $a=0$, and [9] for $a \neq 0$. The identification of critical exponents in the nonradial case for fully nonlinear operators is still open.

Let us recall some preliminary definitions.

Definition 1.1. Let u be a radial solution of (1.1) for $\Omega = \mathbb{R}^N$. Set $r = |x|$ and $\alpha = \frac{2+a}{p-1}$. Then u is said to be:

- i. *fast decaying* if there exists $c > 0$ such that $\lim_{r \rightarrow \infty} r^{\tilde{N}-2} u(r) = c$, where \tilde{N} is either \tilde{N}_+ if the operator is \mathcal{M}^+ or \tilde{N}_- for \mathcal{M}^- , see (1.5);
- ii. *slow decaying* if there exists $c > 0$ such that $\lim_{r \rightarrow \infty} r^\alpha u(r) = c$;
- iii. *pseudo-slow decaying* if there exist constants $0 < c_1 < c_2$ such that

$$c_1 = \liminf_{r \rightarrow \infty} r^\alpha u(r) < \limsup_{r \rightarrow \infty} r^\alpha u(r) = c_2.$$

The definitions (i) and (ii) are classical from the theory of Lane-Emden equations. In turn (iii) was introduced in [6] and is peculiar of the fully nonlinear case. It corresponds to solutions oscillating at $+\infty$ by changing concavity infinitely many times.

Theorem 1.2 (Theorems 1.1 and 1.2 in [6]). *Assume $a=0$, $\tilde{N}_+ > 2$, and $\lambda < \Lambda$. Then there exist critical exponents p_+^* , p_-^* satisfying the bounds*

$$\max \left\{ \frac{\tilde{N}_+}{\tilde{N}_+ - 2}, \frac{N+2}{N-2} \right\} < p_+^* < \frac{\tilde{N}_+ + 2}{\tilde{N}_+ - 2} \quad \text{and} \quad \frac{\tilde{N}_- + 2}{\tilde{N}_- - 2} < p_-^* < \frac{N+2}{N-2},$$

such that the following holds for $\Omega = \mathbb{R}^N$:

- i. if $p \in (1, p_\pm^*)$ there is no nontrivial radial solution of (1.1);

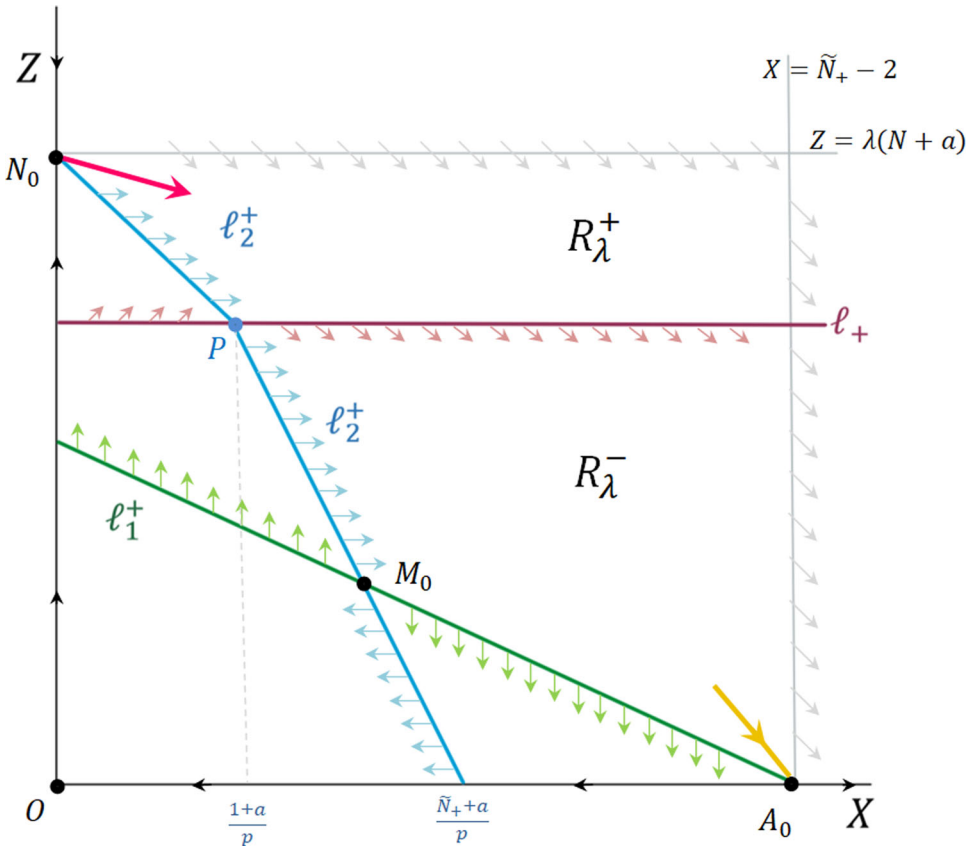


Figure 1. The flow behavior in $1Q$ for \mathcal{M}^+ when $p > p_+^{s,a}$.

- ii. if $p = p_+^*$ there exists a unique fast decaying radial solution to (1.1);
- iii. If $p > p_+^*$ there exists a unique radial solution of (1.1), which is either slow decaying or pseudo-slow decaying.

In (i) and (ii) uniqueness is meant up to scaling.

In addition, in the case of \mathcal{M}^+ the authors made precise the range of the exponent p for which pseudo-slow decaying solutions exist.

The existence of a critical exponent unveils an important feature of the Pucci’s operators. It reflects some intrinsic properties of these operators and induces concentration phenomena besides of energy invariance, see [10], as it happens in the classical semilinear case.

The proof of Theorem 1.2 in [6] is involved. It is a combination of the Emden-Fowler phase plane analysis and the Coffman Kolodner technique. The latter consists in differentiating the solution with respect to the exponent p , and then studying a related nonhomogeneous differential equation, from which they derive the behavior of the solutions for p on both right and left hand sides of p_+^* , as well as the uniqueness of the exponent p for which a fast decaying solution exists.

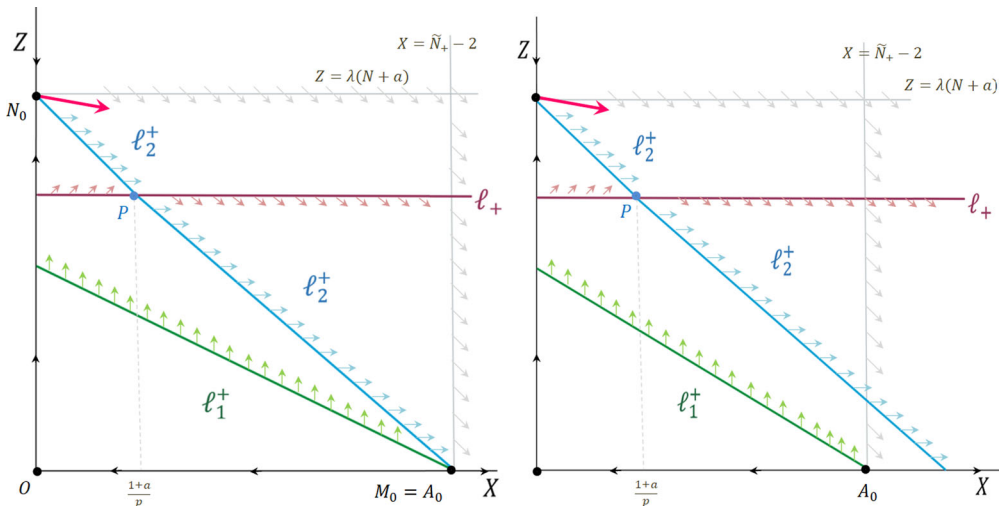


Figure 2. The flow behavior in 1Q for \mathcal{M}^+ when $p = p_+^{s,a}$ (LHS), and $p < p_+^{s,a}$ (RHS).

In this paper we study both regular and singular solutions of the more general weighted problem (1.1). For the regular ones we prove results similar to those of Theorem 1.2 but with different tools entirely based on a dynamical system approach.

Our solutions are understood in the classical sense out of 0 and they are of class C^1 up to 0, since $a > -1$, see Proposition 3.8. Before stating our results it is useful to fix some notations for relevant exponents, depending also on the number $a > -1$ which characterizes the weight in (1.1):

$$p_{\pm}^{p,a} = \frac{\tilde{N}_{\pm} + 2a + 2}{\tilde{N}_{\pm} - 2}, \quad p_{\pm}^{s,a} = \frac{\tilde{N}_{\pm} + a}{\tilde{N}_{\pm} - 2}, \quad p_{\Delta}^a = \frac{N + 2 + 2a}{N - 2}, \quad \alpha = \frac{2 + a}{p - 1}. \tag{1.6}$$

Theorem 1.3 (\mathcal{M}^+ regular solutions). Assume $\tilde{N}_+ > 2$, and $\lambda < \Lambda$. Then there exists a critical exponent p_{a+}^* such that

$$\max\{p_+^{s,a}, p_{\Delta}^a\} < p_{a+}^* < p_+^{p,a}, \tag{1.7}$$

and the following assertions hold:

- i. if $p \in (1, p_{a+}^*)$ there is no nontrivial radial solution of (1.1) in the whole \mathbb{R}^N , while for any $R > 0$ there exists a unique radial solution in the ball B_R ;
- ii. if $p = p_{a+}^*$ there exists a unique fast decaying radial solution of (1.1) in \mathbb{R}^N ;
- iii. if $p \in (p_{a+}^*, p_+^{p,a}]$ there is a unique pseudo-slow decaying radial solution to (1.1) in \mathbb{R}^N ;
- iv. if $p > p_+^{p,a}$ there exists a unique slow decaying radial solution of (1.1) in \mathbb{R}^N ;
- v. if $p > p_{a+}^*$ there is no nontrivial solution to (1.1), (1.3) when Ω is a ball.

In (ii)–(iv) uniqueness is meant up to scaling.

Theorem 1.4 (\mathcal{M}^- regular solutions). If $\lambda < \Lambda$, then there exists a critical exponent p_{a-}^* satisfying

$$p_{-}^{p,a} < p_{a-}^* < p_{\Delta}^a, \quad (1.8)$$

and there exists $\varepsilon > 0$ such that:

- i. if $p \in (1, p_{a-}^*)$ there is no nontrivial radial solution of (1.1) in the whole \mathbb{R}^N , while for any $R > 0$ there exists a unique radial solution of the Dirichlet problem (1.1), (1.3) in B_R ;
- ii. if $p = p_{a-}^*$ there exists a unique fast decaying radial solution of (1.1) in \mathbb{R}^N ;
- iii. if $p \in (p_{a-}^*, p_{\Delta}^a - \varepsilon]$ there is a unique pseudo-slow or slow decaying radial solution of (1.1) in \mathbb{R}^N ;
- iv. if $p > p_{\Delta}^a - \varepsilon$ there exists a unique slow decaying radial solution of (1.1) in \mathbb{R}^N ;
- v. if $p > p_{a-}^*$ there is no nontrivial solution to (1.1), (1.3) when Ω is a ball.

In (ii)–(iv) uniqueness is meant up to scaling.

In the \mathcal{M}^- case our result slightly improves the corresponding one of [6], for $a = 0$, by showing that for p near p_{Δ}^a only a slow decaying solution exists; cf. point (iv) of Theorem 1.4.

The proofs of the previous theorems rely entirely on a careful analysis of an autonomous quadratic dynamical system that we obtain after a suitable transformation, see Section 2. It was used in [11] to study the classical semilinear Lane-Emden system. Once the correspondence between the radial solution of (1.1) and the orbits of the dynamical systems (2.9) and (2.11) is made (see Section 3), all existence and classification results of Theorems 1.3 and 1.4 are derived by studying the stationary points and the flow lines of these systems. In particular, the uniqueness of the critical exponent and the behavior of the solutions, by varying the exponent p , are obtained as a direct consequence of the properties of the vector fields which define the dynamical systems.

Note that these systems are derived from the Pucci fully nonlinear equations and are piecewise C^1 . This, in particular, allows the presence of several periodic orbits which produce regular and singular solutions with different features like pseudo-slow decay or pseudo-blowing up behavior at infinity or at the origin.

One reason why our approach is quite simple is that the most relevant sets which determine the flow generated by (2.9) and (2.11) are just straight lines; see Figures 1 and 2. Moreover, the presence of the weight $|x|^a$ in (1.1) does not produce additional difficulties, while it could be complicated via the method of [6]. We stress that the usual change of variables which transforms Henon problems into non weighted ones for $\lambda = \Lambda$ (see for instance [9, 11, 12]) does not seem to work for Pucci's operators when $\lambda \neq \Lambda$ in order to achieve classification results. Finally we point out that our proofs do not involve any energy function.

On the other hand, by the same analysis of the dynamics induced by (2.9) and (2.11) we also get the classification of singular solutions of (1.1) in a punctured ball or in $\mathbb{R}^N \setminus \{0\}$. Before stating the results we present some definitions.

Definition 1.5. Let u be a radial solution of (1.1) and (1.2), with either $\Omega = \mathbb{R}^N \setminus \{0\}$ or $\Omega = B_R \setminus \{0\}$ for some $R > 0$. Then the singular solution u is said to be:

- i. $(\tilde{N} - 2)$ -blowing up if there exists $c > 0$ such that $\lim_{r \rightarrow 0} r^{\tilde{N}-2} u(r) = c$, where \tilde{N} is either \tilde{N}_+ if the operator is \mathcal{M}^+ or \tilde{N}_- for \mathcal{M}^- , see (1.5);
- ii. α -blowing up if there exists $c > 0$ such that $\lim_{r \rightarrow 0} r^\alpha u(r) = c$, with α as in (1.6);
- iii. pseudo-blowing up if there exist constants $0 < c_1 < c_2$ such that

$$c_1 = \liminf_{r \rightarrow 0} r^\alpha u(r) < \limsup_{r \rightarrow 0} r^\alpha u(r) = c_2.$$

We highlight that Definition 1.5 (iii) corresponds to a type of solutions which change concavity infinitely many times in a neighborhood of zero. The existence of such a type of singular solutions was already detected for a more general class of uniformly elliptic equations, for values of the exponent p close to the critical one, see [13, Section 6].

Remark 1.6. For all $p > p_+^{s,a}$ in the case of \mathcal{M}^+ , resp. $p > p_-^{s,a}$ for \mathcal{M}^- , the function $u_p(r) = C_p r^{-\alpha}$, C_p as in (3.8), is a singular solution of (1.1) in $\mathbb{R}^N \setminus \{0\}$. We call it *trivial singular solution*.

Theorem 1.7 (\mathcal{M}^+ singular solutions). *Assuming $\tilde{N}_+ > 2$ and $\lambda < \Lambda$, for (1.1)–(1.2) it holds:*

- i. for any $p \leq p_+^{s,a}$ there is no singular radial solution in $\mathbb{R}^N \setminus \{0\}$, while for each $R > 0$ there are infinitely many $(\tilde{N}_+ - 2)$ -blowing up radial solutions of (1.1)–(1.3) in $B_R \setminus \{0\}$;
- ii. if $p_+^{s,a} \leq p_\Delta^a$ then for any $p \in (p_+^{s,a}, p_\Delta^a]$ there is a unique α -blowing up radial solution in $\mathbb{R}^N \setminus \{0\}$ with fast decay at $+\infty$. Also, for any $R > 0$ there are infinitely many α -blowing up radial solutions of (1.1)–(1.3) in $B_R \setminus \{0\}$;
- iii. for each $p \in (p_\Delta^a, p_{a+}^*)$ there exists a unique singular radial solution in $\mathbb{R}^N \setminus \{0\}$ with fast decay at $+\infty$. Moreover, for any $R > 0$ there exist infinitely many singular radial solutions of the Dirichlet problem (1.1)–(1.3) in $B_R \setminus \{0\}$;
- iv. if $p = p_{a+}^*$ there exist infinitely many pseudo-blowing up radial solutions in $\mathbb{R}^N \setminus \{0\}$ with pseudo-slow decay at $+\infty$, and infinitely many α -blowing up in $\mathbb{R}^N \setminus \{0\}$ with pseudo-slow decay at $+\infty$. Also, there is no singular radial solution of (1.1)–(1.3) in $B_R \setminus \{0\}$;
- v. if $p \in (p_{a+}^*, p_+^{p,a})$ there are infinitely many α -blowing up radial solutions in $\mathbb{R}^N \setminus \{0\}$ with pseudo-slow decay at $+\infty$, and there is a pseudo-blowing up radial solution with pseudo-slow decay at $+\infty$. Further, there is no singular radial solution of (1.1)–(1.3) in $B_R \setminus \{0\}$;
- vi. if $p \in [p_+^{p,a}, +\infty)$ there are no nontrivial singular radial solutions, cf. Remark 1.6.

Here, uniqueness in $\mathbb{R}^N \setminus \{0\}$ is meant up to scaling.

Theorem 1.8 (\mathcal{M}^- singular solutions). *If $\lambda < \Lambda$, for the problem (1.1)–(1.2) we have:*

- i. if $p \leq p_-^{s,a}$ there is no singular radial solution in the whole $\mathbb{R}^N \setminus \{0\}$, while for any $R > 0$ there are infinitely many $(\tilde{N}_- - 2)$ -blowing up radial solutions of (1.1)–(1.3) in $B_R \setminus \{0\}$;

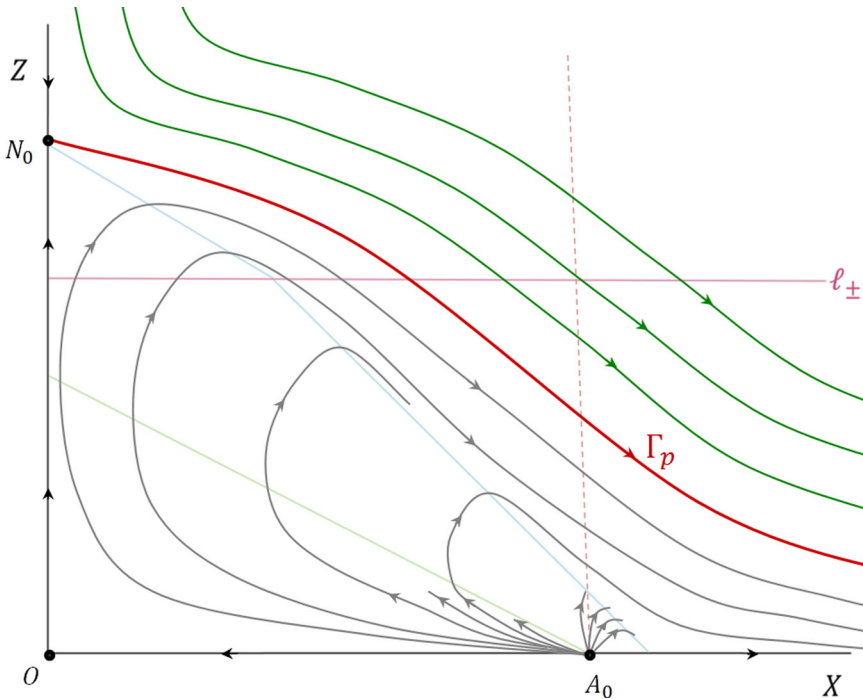


Figure 3. Case $p < p_{\pm}^{s,a}$ for \mathcal{M}^{\pm} ; A_0 is a source and M_0 belongs to the fourth quadrant. Below Γ_p trajectories corresponding to infinitely many $(\tilde{N}_{\pm} - 2)$ -blowing up solutions in a ball are shown, while above Γ_p there are orbits corresponding to solutions in an annulus.

- ii. for each $p \in (p_{-}^{s,a}, p_{-}^{p,a})$ there exists a unique α -blowing up radial solution in $\mathbb{R}^N \setminus \{0\}$ with fast decay at $+\infty$. Further, for any $R > 0$ there exist infinitely many α -blowing up radial solutions of the Dirichlet problem (1.1)–(1.3) in $B_R \setminus \{0\}$;
- iii. for any $p \in (p_{-}^{p,a}, p_{a-}^*)$ there are infinitely many pseudo-blowing up radial solutions in $\mathbb{R}^N \setminus \{0\}$. Among them there is a unique fast decaying, a pseudo-slow decaying, and infinitely many with slow decay at $+\infty$. Moreover, for each $R > 0$ there exist infinitely many pseudo-blowing up radial solutions of (1.1)–(1.3) in $B_R \setminus \{0\}$;
- iv. if $p = p_{a-}^*$ there exist infinitely many pseudo-blowing up radial solutions in $\mathbb{R}^N \setminus \{0\}$. Among them there are infinitely many with slow-decay at $+\infty$, and infinitely many pseudo-slow decaying at $+\infty$. Further, there is no singular radial solution of (1.1)–(1.3) in $B_R \setminus \{0\}$;
- v. there exists $\varepsilon > 0$ such that for $p \in [p_{\Delta}^a - \varepsilon, +\infty)$ no nontrivial singular radial solution exists.

Here, uniqueness in $\mathbb{R}^N \setminus \{0\}$ is meant up to scaling.

Our results on singular solutions are obtained by complementing the analysis of the flow lines of the dynamical systems (2.9) and (2.11). To the best of our knowledge they are the first global classification results on singular solutions found for

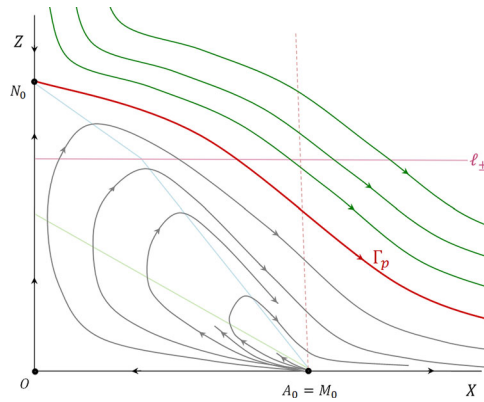


Figure 4. Case $p = p_{\pm}^{s,a}$: $p \in \mathcal{C}$ and $A_0 = M_0$ has infinitely many unstable directions. Below Γ_p are the orbits corresponding to infinitely many $(N_{\pm} - 2)$ -blowing up solutions in a ball.

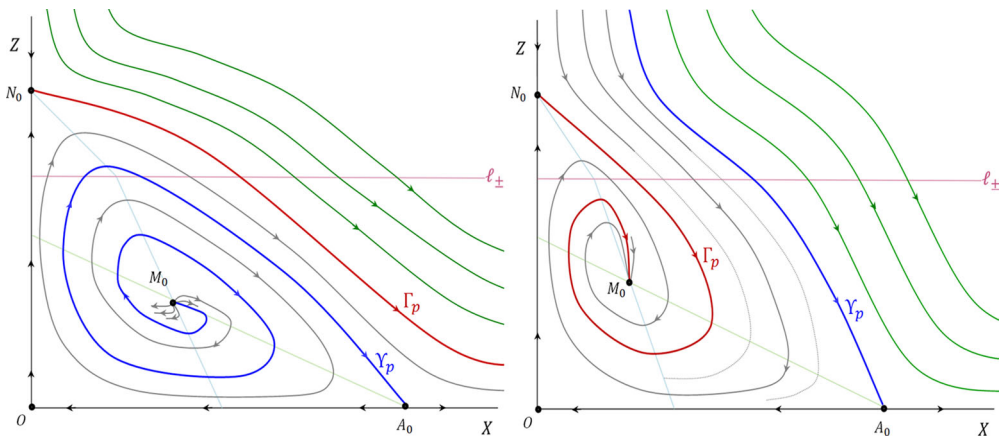


Figure 5. $p > p_{\pm}^{s,a}$, cases $p \in \mathcal{C}$ (LHS) and $p \in \mathcal{S}$ (RHS) without periodic orbits. For instance, this describes the range of p where Dulac criterion holds, for both operators \mathcal{M}^{\pm} .

this class of fully nonlinear equations. In [6, Remark 3.2] it is pointed out that, in the case $a = 0$, periodic orbits of the Endem-Fowler system would produce singular solutions, while in [13, Theorem 6.3] the existence of singular solutions is proved near the critical exponent.

For the critical exponents $p_{a\pm}^*$, our dynamical systems (2.9) and (2.11) furnish infinitely many periodic orbits. On the other hand, for $p_{\pm}^{p,a}$ infinitely many periodic orbits appear which do not correspond to C^2 solutions for $r > 0$, see Remark 4.10. For $p \in (p_{a-}^*, p_{\Delta}^a - \varepsilon)$ the existence of singular solutions cannot be deduced directly from the dynamical system approach.

Let us underline the fact that obtaining periodic orbits is in general a very difficult task in the theory of dynamical systems. Even in the very particular case of a polynomial autonomous system this question is not completely understood, see [14, 15].

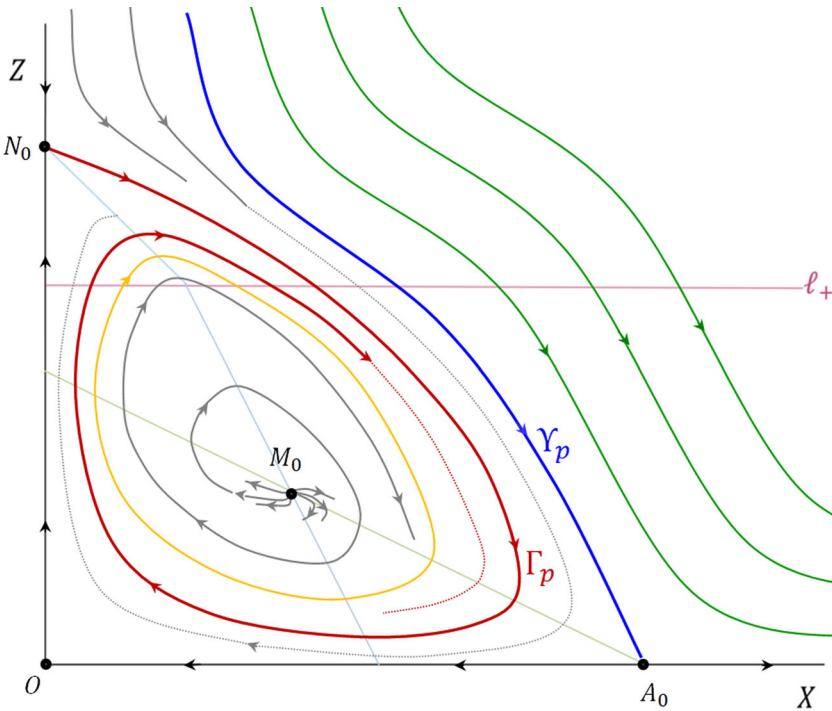


Figure 6. Case $p \in \mathcal{P}$ for $\mathcal{M}^+, p \in (p_{a^+}^*, p_+^{p,a})$. Here M_0 is a source. The orbits inside the displayed periodic orbit correspond to infinitely many α -blowing up solutions with pseudo-slow decay at $+\infty$. All trajectories above Γ_p correspond to solutions either in the exterior of a ball or in an annulus.

Finally, as a byproduct of the study of regular radial solutions of (1.1), either in \mathbb{R}^N or in a ball, we easily get the range of the exponents p for which a positive radial solution of the Dirichlet problem in the exterior of a ball does not exist. Indeed, we get the following result.

Theorem 1.9. *Let $p > 1$. Then there are no radial solutions of*

$$\mathcal{M}^+(D^2u) + |x|^a u^p = 0, \quad u > 0 \text{ in } \mathbb{R}^N \setminus B_R, \quad u = 0 \text{ on } \partial B_R \tag{1.9}$$

if $p \leq p_{a^\pm}^$ for each $R > 0$.*

In the case of $a = 0$, Theorem 1.9 has been recently proved in [16] with different arguments which rely both on the study of the second order ODE and on the analysis of the Emden-Fowler system. Their work presents a complete picture of existence and nonexistence of solutions for distinct intervals for the values of the parameter p . However, through our arguments we get their nonexistence result by a considerably shorter proof. Indeed, we will see in Sections 4 and 5 that the result of Theorem 1.9 becomes a straightforward consequence of the characterization of the critical exponents $p_{a^\pm}^*$ in terms of the associated quadratic system we consider. Let us point out that in [16] also the existence and classification of the solutions of (1.9) are provided when

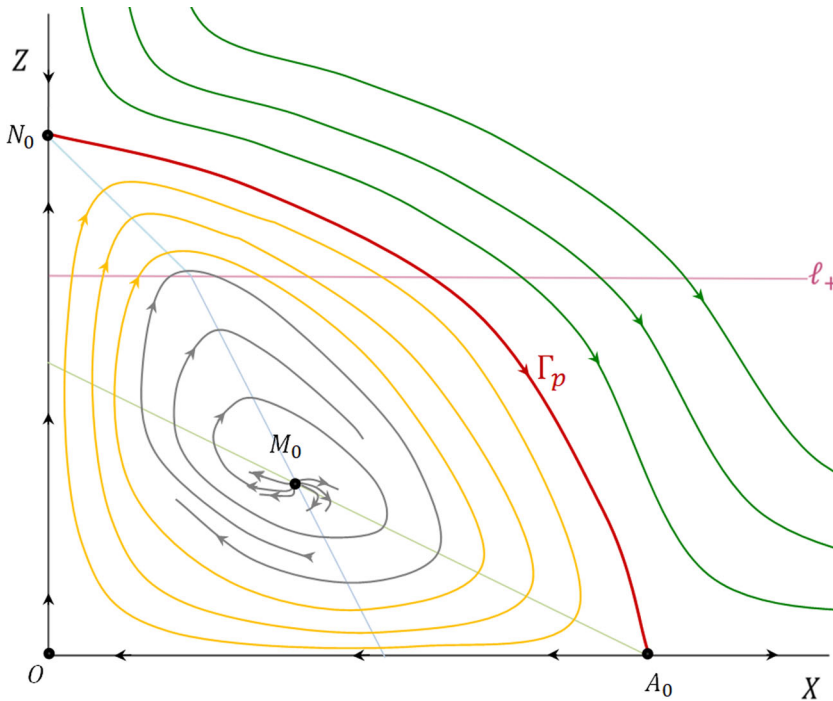


Figure 7. Case $p = p_{a_+}^* \in \mathcal{F}$. Here M_0 is a source and $\Gamma_p = \Upsilon_p$, see Remark 4.8. There are infinitely many α -blowing up solutions with pseudo-slow decay (inside the minimal periodic orbit), and infinitely many pseudo-blowing up solutions with pseudo-slow decay (periodic orbits). Moreover, there are no solutions in the exterior of a ball.

$a = 0$. Alternatively, this could be done through our methods. Since this is not the main goal of our research we just refer to Section 4 for further comments.

To conclude, we stress that another advantage of our approach is that it treats in a unified way several kind of solutions of (1.1). We refer the reader to Figures 3–8 where, for a given value of the exponent p , all the orbits of the system corresponding to different type of solutions of (1.1) are displayed simultaneously.

The paper is organized as follows. In Section 2 we write down the quadratic system associated to the problem (1.1) and study its intrinsic flow properties. In Section 3 we classify the different solutions of (1.1) in terms of orbits of the corresponding dynamical systems. Finally, Sections 4 and 5 are devoted to the proofs of the main results for the Pucci \mathcal{M}^+ and \mathcal{M}^- operators, respectively. In the Appendix we provide some details about the stationary points of the dynamical systems, for the reader convenience.

2. The associated dynamical system

In this section we define some new variables which allow to transform the radial fully nonlinear equations into a quadratic dynamical system.

We start by recalling that the Pucci’s extremal operators $\mathcal{M}_{\lambda, \Lambda}^\pm$, for $0 < \lambda \leq \Lambda$, are defined as

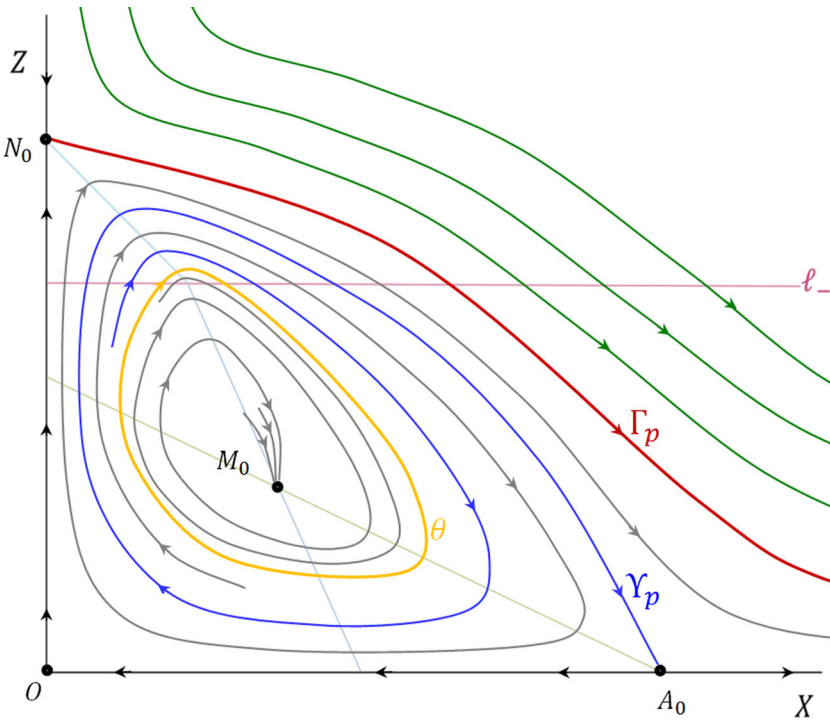


Figure 8. Case $p \in (p_{a-}^{p,a}, p_{a-}^*)$ for \mathcal{M}^- , M_0 is a sink. There are infinitely many pseudo-blowing up solutions: a unique fast decaying (given via the orbit Υ_p); a pseudo-slow decaying (periodic orbit θ); infinitely many in a ball (outside θ); infinitely many slow decaying (inside θ).

$$\mathcal{M}_{\lambda, \Lambda}^+(X) := \sup_{\lambda I \leq A \leq \Lambda I} \text{tr}(AX), \quad \mathcal{M}_{\lambda, \Lambda}^-(X) := \inf_{\lambda I \leq A \leq \Lambda I} \text{tr}(AX),$$

where A, X are $N \times N$ symmetric matrices, and I is the identity matrix. Equivalently, if we denote by $\{e_i\}_{1 \leq i \leq N}$ the eigenvalues of X , we can define the Pucci's operators as

$$\mathcal{M}_{\lambda, \Lambda}^+(X) = \Lambda \sum_{e_i > 0} e_i + \lambda \sum_{e_i < 0} e_i, \quad \mathcal{M}_{\lambda, \Lambda}^-(X) = \lambda \sum_{e_i > 0} e_i + \Lambda \sum_{e_i < 0} e_i. \tag{2.1}$$

From now on we will drop writing the parameters λ, Λ in the notations for the Pucci's operators.

In the case when u is a radial function, with an abuse of notation we set $u(|x|) = u(r)$ for $r = |x|$. If in addition u is C^2 , then the eigenvalues of the Hessian matrix D^2u are u'' which is simple, and $\frac{u'(r)}{r}$ with multiplicity $N - 1$. We then define the Lipschitz functions

$$m_+(s) = \begin{cases} \lambda s & \text{if } s \leq 0 \\ \Lambda s & \text{if } s > 0 \end{cases} \quad \text{and} \quad M_+(s) = \begin{cases} s/\lambda & \text{if } s \leq 0 \\ s/\Lambda & \text{if } s > 0; \end{cases} \tag{2.2}$$

$$m_-(s) = \begin{cases} \Lambda s & \text{if } s \leq 0 \\ \lambda s & \text{if } s > 0 \end{cases} \quad \text{and} \quad M_-(s) = \begin{cases} s/\Lambda & \text{if } s \leq 0 \\ s/\lambda & \text{if } s > 0. \end{cases} \tag{2.3}$$

The equations $\mathcal{M}^+(D^2u) + r^a u^p = 0$ and $\mathcal{M}^-(D^2u) + r^a u^p = 0$, for $r \neq 0$, in radial coordinates for positive solutions then become, respectively,

$$u'' = M_+(-r^{-1}(N - 1) m_+(u') - r^a u^p), \quad u > 0; (P_+)$$

$$u'' = M_-(-r^{-1}(N - 1) m_-(u') - r^a u^p), \quad u > 0, (P_-)$$

which are understood in the maximal interval where u is positive.

2.1. The new variables and the quadratic system

Let u be a positive solution of (P_+) or (P_-) . Thus we can define the new functions

$$X(t) = -\frac{ru'(r)}{u(r)}, \quad Z(t) = -\frac{r^{1+a} u^p(r)}{u'(r)} \quad \text{for } t = \ln(r), \tag{2.4}$$

whenever $r > 0$ is such that $u(r) \neq 0$ and $u'(r) \neq 0$.

We consider the phase plane $(X, Z) \in \mathbb{R}^2$. Since we are studying positive solutions, the points $(X(t), Z(t))$ belong to the first quadrant when $u' < 0$; or to the third quadrant when $u' > 0$. We denote the first and third quadrants by 1Q, 3Q respectively.

As a consequence of this monotonicity, the problems (P_+) and (P_-) become in 1Q:

$$\text{for } \mathcal{M}^+ : \quad u'' = M_+(-\lambda r^{-1}(N - 1)u' - r^a u^p), \quad u > 0 \quad \text{in 1Q,} \tag{2.5}$$

$$\text{for } \mathcal{M}^- : \quad u'' = M_-(-\Lambda r^{-1}(N - 1)u' - r^a u^p), \quad u > 0 \quad \text{in 1Q.} \tag{2.6}$$

On the other hand, since $u' > 0$ implies $u'' < 0$, one finds out in 3Q:

$$\text{for } \mathcal{M}^+ : \quad \lambda u'' = -\Lambda r^{-1}(N - 1)u' - r^a u^p, \quad u > 0 \quad \text{in 3Q,} \tag{2.7}$$

$$\text{for } \mathcal{M}^- : \quad \Lambda u'' = -\lambda r^{-1}(N - 1)u' - r^a u^p, \quad u > 0 \quad \text{in 3Q.} \tag{2.8}$$

In terms of the functions (2.4), we derive the following autonomous dynamical systems:

$$\text{in 1Q,} \quad \begin{cases} \dot{X} = X[X + 1 - M_+(\lambda(N - 1) - Z)], \\ \dot{Z} = Z[1 + a - pX + M_+(\lambda(N - 1) - Z)]; \end{cases} \tag{2.9}$$

$$\text{in 3Q,} \quad \dot{X} = X[X - (\tilde{N}_- - 2) + Z/\Lambda], \dot{Z} = Z[\tilde{N}_- + a - pX - Z/\Lambda], \tag{2.10}$$

corresponding to (2.5), (2.7) for \mathcal{M}^+ , where the dot $\dot{\cdot}$ stands for $\frac{d}{dt}$. Similarly one has

$$\text{in 1Q,} \quad \begin{cases} \dot{X} = X[X + 1 - M_-(\Lambda(N - 1) - Z)], \\ \dot{Z} = Z[1 + a - pX + M_-(\Lambda(N - 1) - Z)]; \end{cases} \tag{2.11}$$

$$\text{in 3Q,} \quad \dot{X} = X[X - (\tilde{N}_+ - 2) + Z/\lambda], \dot{Z} = Z[\tilde{N}_+ + a - pX - Z/\lambda], \tag{2.12}$$

associated to (2.6), (2.8) for \mathcal{M}^- .

We stress that (2.9), (2.11) correspond to positive, decreasing solutions of (P_+) , (P_-) . We will see in Section 3 that this holds for regular and singular solutions of (1.1) in \mathbb{R}^N or in a ball.

On the other hand, given any trajectory $\tau = (X, Z)$ of (2.9)-(2.12) either in 1Q or 3Q, we define

$$u(r) = r^{-\alpha} (XZ)^{\frac{1}{p-1}}(t), \quad \text{where } r = e^t. \tag{2.13}$$

Then we deduce

$$\begin{aligned} u'(r) &= -\alpha r^{-\alpha-1} (XZ)^{\frac{1}{p-1}}(t) + \frac{r^{-\alpha}}{p-1} (XZ)^{\frac{1}{p-1}-1}(t) \frac{\dot{X}Z + X\dot{Z}}{r} = -Xr^{-\alpha-1} (XZ)^{\frac{1}{p-1}}(t) \\ &= -\frac{X(t)u(r)}{r}, \end{aligned}$$

from which we recover (2.4). Since $X \in C^1$, then $u \in C^2$. From this, one immediately sees that u satisfies either (P_+) or (P_-) from the respective equations for \dot{X}, \dot{Z} in the dynamical system.

An important role in the study of our problem is played by the lines ℓ_{\pm} , defined by

$$\begin{aligned} \ell_+ &= \{(X, Z) : Z = \lambda(N - 1)\} \cap 1Q \text{ for } \mathcal{M}^+, \\ \ell_- &= \{(X, Z) : Z = \Lambda(N - 1)\} \cap 1Q \text{ for } \mathcal{M}^-. \end{aligned} \tag{2.14}$$

For each of the two systems (2.9) and (2.11) respectively, the lines ℓ_{\pm} splits 1Q into two regions, up and down:

$$R_{\lambda}^+ = \{(X, Z) : Z > \lambda(N - 1)\} \cap 1Q, \quad R_{\lambda}^- = \{(X, Z) : Z < \lambda(N - 1)\} \cap 1Q \text{ for } \mathcal{M}^+, \tag{2.15}$$

$$R_{\Lambda}^+ = \{(X, Z) : Z > \Lambda(N - 1)\} \cap 1Q, \quad R_{\Lambda}^- = \{(X, Z) : Z < \Lambda(N - 1)\} \cap 1Q \text{ for } \mathcal{M}^-. \tag{2.16}$$

In terms of (P_{\pm}) , ℓ_{\pm} is the line where a decreasing solution u changes concavity in the sense that, when $(X(t), Z(t)) \in R_{\lambda}^+$ (or R_{Λ}^+) then the corresponding solution u through the transformation (2.13) is concave, while for $(X(t), Z(t)) \in R_{\lambda}^-$ (or R_{Λ}^-), u is convex. Hence, these regions are essential to determine the precise expressions of (P_+) and (P_-) according to M_+ and M_- in (2.2), (2.3). For instance, when $u'' \leq 0$ and $u' < 0$, (P_+) becomes

$$-u''(r) - \frac{N-1}{r} u'(r) = r^a \frac{u^p(r)}{\lambda},$$

for which the left hand side is the standard radial Laplacian operator; while for $u'' > 0$ and $u' < 0$, (P_+) reads as

$$-u''(r) - \frac{\tilde{N}_+ - 1}{r} u'(r) = r^a \frac{u^p(r)}{\Lambda},$$

where, in turn, the LHS is the Laplacian in the noninteger dimension \tilde{N}_+ , see (1.5). Analogously one treats \mathcal{M}^- . Note that in 3Q we always obtain Laplacian operators in dimensions \tilde{N}_- (for \mathcal{M}^+) and \tilde{N}_+ (for \mathcal{M}^-), see (2.7), (2.8). Our system is then the union of equations driven by different Laplacian-like operators. This explains the difficulty in dealing with fully nonlinear operators.

We stress that Lane-Emden-Henon problems for Laplacian operators were already studied in [11] in terms of the dynamical system (2.9) in the case $\lambda = \Lambda = 1$ subject to the transformation (2.4).

At this stage it is worth observing that the systems (2.9) and (2.11) are continuous on ℓ_{\pm} . More than that, the right hand sides are locally Lipschitz functions of X, Z , so the

usual ODE theory applies. That is, one recovers existence, uniqueness, and continuity with respect to initial data as well as continuity with respect to the parameter p .

2.2. Stationary points and local analysis

We start this section investigating the sets where $\dot{X} = 0$ and $\dot{Z} = 0$. Let us focus our analysis on $1Q$, since the only stationary point on the boundary of $3Q$ is the origin. One writes the dynamical systems (2.9) and (2.11) in terms of the following ODE first order autonomous equation

$$\dot{x} = F(x), \quad \text{where } x = (X, Z), \quad F(x) := (f(x), g(x)). \quad (2.17)$$

with $\dot{x} = (\dot{X}, \dot{Z})$. For instance, in the case of the operator \mathcal{M}^+ , then f, g are given by

$$f(x) = \begin{cases} X \left(X - (N - 2) + \frac{Z}{\lambda} \right) & \text{in } R_\lambda^+ \\ X \left(X - (\tilde{N}_+ - 2) + \frac{Z}{\Lambda} \right) & \text{in } R_\lambda^- \end{cases}, \quad g(x) = \begin{cases} Z \left(N + a - pX - \frac{Z}{\lambda} \right) & \text{in } R_\lambda^+ \\ Z \left(\tilde{N}_+ + a - pX - \frac{Z}{\Lambda} \right) & \text{in } R_\lambda^- \end{cases}.$$

We first recall some standard definitions from the theory of dynamical systems.

Definition 2.1. A stationary point Q of (2.17) is a zero of the vector field F . If σ_1 and σ_2 are the eigenvalues of the Jacobian matrix $DF(Q)$, then Q is hyperbolic if both σ_1, σ_2 have nonzero real parts. If this is the case, Q is a *source* if $\text{Re}(\sigma_1), \text{Re}(\sigma_2) > 0$, and a *sink* if $\text{Re}(\sigma_1), \text{Re}(\sigma_2) < 0$; Q is a *saddle point* if $\text{Re}(\sigma_1) < 0 < \text{Re}(\sigma_2)$.

Next we recall an important result from the theory of dynamical systems which describes the local stable and unstable manifolds near saddle points of the system (2.17); see [15, theorems 9.29, 9.35]. Here the usual theory for autonomous planar systems applies since each stationary point Q possesses a neighborhood which is strictly contained in R_λ^+ or R_λ^- where the vector field F is C^1 .

Proposition 2.2. *Let Q be a saddle point of (2.17). Then the local stable (resp. unstable) manifold at Q is locally a C^1 graph over the stable (resp. unstable) line of the linearized vector field. In this case, if the linearized system has a stable line direction L , then there exists exactly two trajectories τ_1 and τ_2 arriving at Q which admit the same tangent at the point $Q = \bar{\tau}_1 \cap \bar{\tau}_2$ given by L . Analogously there are only two trajectories coming out from Q with the same property.*

We sometimes use the following notation to describe the limit of trajectories in the phase plane.

Definition 2.3 (α and ω limits). We call α -limit of the orbit τ , and we denote it by $\alpha(\tau)$, as the set of limit points of $\tau(t)$ as $t \rightarrow -\infty$. Similarly one defines $\omega(\tau)$ i.e. the ω -limit of τ at $+\infty$.

We observe that both X and Z axes are invariant by the flow. In particular, each quadrant is an invariant set for the dynamics. Moreover, let us keep in mind the following segments in the plane (X, Z) . For the system (2.9), we define

$$\ell_1^+ = \{(X, Z) : Z = \Lambda(\tilde{N}_+ - 2) - \Lambda X\} \cap 1Q \quad (2.18)$$

which is the set where $\dot{X} = 0$ and $X > 0$; also

$$\ell_2^+ = \ell_{2+}^+ \cup \ell_{2-}^+ \tag{2.19}$$

with $\ell_{2+}^+ = \{(X, Z) : Z = \lambda(N + a - pX)\} \cap R_\lambda^+$, $\ell_{2-}^+ = \{(X, Z) : Z = \Lambda(\tilde{N}_+ + a - pX)\} \cap R_\lambda^-$, which is the set where $\dot{Z} = 0$ and $Z > 0$; see [Figures 1 and 2](#).

Notice that ℓ_1^+ is a segment entirely contained in R_λ^- , since there are no other points in 1Q where $\dot{X} = 0$ in the interior of the region R_λ^+ . Moreover, (2.19) is the union of two segments which join at the point $(\frac{1+a}{p}, \lambda(N - 1)) \in \ell_+ \cap \bar{\ell}_2^+$, see [Figure 1](#). The analogous sets for \mathcal{M}^- are

$$\ell_1^- = \{(X, Z) : Z = \lambda(\tilde{N}_- - 2) - \lambda X\} \cap 1Q \tag{2.20}$$

which is the set where $\dot{X} = 0$ and $X > 0$ (contained in R_λ^-); and

$$\ell_2^- = \ell_{2+}^- \cup \ell_{2-}^- \tag{2.21}$$

with $\ell_{2+}^- = \{(X, Z) : Z = \Lambda(N_- + a - pX)\} \cap R_\Lambda^+$, $\ell_{2-}^- = \{(X, Z) : Z = \lambda(\tilde{N}_- + a - pX)\} \cap R_\Lambda^-$, which is the set where $\dot{Z} = 0$ and $Z > 0$.

Lemma 2.4. *The stationary points of the dynamical systems (2.9)–(2.12) are:*

for \mathcal{M}^+ : $O = (0, 0)$, $N_0 = (0, \lambda N + \lambda a)$, $A_0 = (\tilde{N}_+ - 2, 0)$, $M_0 = (X_0, Z_0)$,

where $X_0 = \alpha$, and $Z_0 = \Lambda(\tilde{N}_+ - p\alpha + a) = \Lambda(\tilde{N}_+ - 2 - \alpha)$, see [Figure 1](#);

for \mathcal{M}^- : $O = (0, 0)$, $N_0 = (0, \Lambda N + \Lambda a)$, $A_0 = (\tilde{N}_- - 2, 0)$, $M_0 = (X_0, Z_0)$,

where $X_0 = \alpha$ and $Z_0 = \lambda(\tilde{N}_- - p\alpha + a) = \lambda(\tilde{N}_- - 2 - \alpha)$.

Proof. We just show the \mathcal{M}^+ case. First notice that the system does not admit stationary points in 3Q nor on the line ℓ_+ . In the region \bar{R}_λ^+ we have already seen that $\dot{X} = 0$ implies $X = 0$, since ℓ_1^+ does not intersect R_λ^+ . By $\dot{Z} = 0$ we obtain $Z = \lambda(N + a - pX)$ since $Z \neq 0$ in R_λ^+ . Hence we reach the equilibrium point N_0 . In \bar{R}_λ^- , from $\dot{X} = 0$ we have either $X = 0$ or $Z = \Lambda(\tilde{N}_+ - 2 - X)$, while by $\dot{Z} = 0$ we deduce that either $Z = 0$ or $Z = \Lambda(\tilde{N}_+ + a - pX)$. Therefore we obtain the points O , A_0 , M_0 , and $(0, \Lambda(\tilde{N}_+ + a))$. However, the latter does not belong to \bar{R}_λ^- as long as $a > -1$. \square

Next we analyze the directions of the vector field F in (2.17) on the X, Z axes, on the concavity line ℓ_\pm , and also on ℓ_1^\pm, ℓ_2^\pm ; see (2.14), and (2.18)–(2.21).

Proposition 2.5. *The systems (2.9) and (2.11) enjoy the following properties (see [Figures 1 and 2](#)):*

- i. *Every trajectory of (2.9) in 1Q crosses the line ℓ_+ transversely except at the point $P = (\frac{1+a}{p}, \lambda(N - 1))$. Moreover, it passes from R_λ^+ to R_λ^- if $X > \frac{1+a}{p}$, while it moves from R_λ^- to R_λ^+ if $X < \frac{1+a}{p}$. The vector field at P always points to the right. A similar statement holds for \mathcal{M}^- via the system (2.11) considering respectively ℓ_- , $(\frac{1+a}{p}, \Lambda(N - 1))$, R_Λ^+ , R_Λ^- ;*
- ii. *The flow induced by (2.17) on the X axis points to the left for $X \in (0, \tilde{N}_\pm - 2)$, and to the right when $X > \tilde{N}_\pm - 2$. On the Z axis it moves up between O and N_0 , and down above N_0 ;*

- iii. The vector field F on the line ℓ_1^\pm is parallel to the Z axis whenever $X \neq \alpha$. It points up if $X < \alpha$, and down if $X > \alpha$. Further, on the set ℓ_2^\pm the vector field F is parallel to the Z axis for $X \neq \alpha$. It moves to the right if $X < \alpha$, and to the left if $X > \alpha$.

Proof. (1) We just observe that $\dot{X} = X(X + 1) > 0$, and $\dot{Z} = Z(1 + a - pX)$ on ℓ_\pm .

(2) For instance consider \mathcal{M}^+ . Since the X axis is contained in R_λ^- , then $\dot{X} = X(X - (\tilde{N}_+ - 2))$ which is positive for $X < \tilde{N}_+ - 2$ and negative for $X > \tilde{N}_+ - 2$. Now, $\dot{Z} = Z(N + a - Z/\lambda)$ in R_λ^+ is positive if $Z < \lambda(N + a)$ and negative for $Z > \lambda(N + a)$. On the other hand, $\dot{Z} = Z(\tilde{N}_+ + a - Z/\Lambda) > 0$ in R_λ^- , since $\Lambda(\tilde{N}_+ + a) > \lambda(N - 1)$ for $a > -1$.

(3) Notice that $\dot{Z} = (p - 1)Z(\alpha - X)$ on ℓ_1^\pm and $\dot{X} = (p - 1)X(\alpha - X)$ on ℓ_2^\pm . Both are positive quantities for $X < \alpha$, and negative when $X > \alpha$. □

Remark 2.6. It is not difficult to see that an orbit can only reach the point P in Proposition 2.5 (1) from R_λ^- . In fact by $\dot{X} > 0$ and the inverse function theorem, Z is a function of X on ℓ_+ . Next, an analysis of the continuous function $X \mapsto \partial_X^2 Z(X)$ at P shows that an orbit passing through P has a local maximum at this point; see the vector field at that point in Figures 1 and 2.

The next proposition gathers the crucial dynamics at each stationary point.

Proposition 2.7 (\mathcal{M}^\pm). *The following properties are verified for the systems (2.9) and (2.11),*

1. The origin O is a saddle point. The stable and unstable directions of the linearized system are the X and Z axes respectively;
2. N_0 is a saddle point. The tangent unstable direction is parallel to the line

$$Z = \frac{-p\lambda(N + a)}{N + 2 + 2a}X \text{ if the operator is } \mathcal{M}^+, \quad Z = \frac{-p\Lambda(N + a)}{N + 2 + 2a}X \text{ for } \mathcal{M}^-;$$

3. A_0 is a saddle point for $p > p_\pm^{s,a}$. The linear stable direction is parallel to the line

$$Z = \frac{-p(\tilde{N}_+ - 2) + 2 + a}{\tilde{N}_+ - 2}\Lambda X \text{ in the case of } \mathcal{M}^+, \quad Z = \frac{-p(\tilde{N}_- - 2) + 2 + a}{\tilde{N}_- - 2}\lambda X \text{ for } \mathcal{M}^-,$$

while the unstable tangent direction lies on the X axis. For $p < p_\pm^{s,a}$ A_0 is a source; at $p = p_\pm^{s,a}$ A_0 coincides with M_0 and belongs to the X axis. In this case, it is not a hyperbolic stationary point.

4. For $p < p_\pm^{s,a}$ M_0 belongs to the fourth quadrant. Also, $M_0 \in 1Q \iff p > p_\pm^{s,a}$ in which case:
 - i. M_0 is a source if $p_\pm^{s,a} < p < p_\pm^{p,a}$;
 - ii. M_0 is a sink for $p > p_\pm^{p,a}$;
 - iii. M_0 is a center at $p = p_\pm^{p,a}$.

The dynamics at each stationary point depends upon the linearization of the system (2.9). Since the point N_0 belongs to R_λ^+ where the system corresponds to the Henon equation for the standard Laplacian, we could just refer to [11] for the local analysis of

N_0 , as long as $p > p_{\pm}^{p,a}$. The other points O, N_0, A_0 instead belong to R_{λ}^- where the system now corresponds to the Henon equation for the Laplacian in dimension \tilde{N}_{\pm} . In this last case some variations with respect to [11] are needed. Since the classification of stationary points is the heart of our analysis and for reader's convenience, details are provided in the [Appendix](#).

On the other hand, in both cases, some deeper analysis is required when $p \leq p_{\pm}^{p,a}$. We treat this case in [Proposition 4.9](#) by using the dynamics of the system. For $p = p_{\pm}^{p,a}$ we refer to the Appendix, see [Proposition 6.1](#).

To finish the section, a local uniqueness result follows directly from [Propositions 2.2](#) and [2.7](#).

Proposition 2.8. *For every $p > 1$ there is a unique trajectory coming out from N_0 at $-\infty$, which we denote by Γ_p . Further, for $p > p_{\pm}^{s,a}$ there exists a unique trajectory arriving at A_0 at $+\infty$ that we denote by Υ_p . In terms of [Definition 2.3](#),*

$$\text{for all } p > 1, \quad \Gamma_p \text{ is such that } \alpha(\Gamma_p) = N_0; \tag{2.22}$$

$$\text{for all } p > p_{\pm}^{s,a}, \quad \Upsilon_p \text{ is such that } \omega(\Upsilon_p) = A_0. \tag{2.23}$$

Remark 2.9. Notice that these trajectories uniquely determine the global unstable and stable manifolds of the stationary points N_0 and A_0 respectively. In particular, by [Proposition 2.2](#) they are graphs of functions in a neighborhood of the stationary points in their respective ranges of p . The tangent directions at N_0, A_0 are displayed together in [Figure 1](#) for $p > p_{\pm}^{s,a}$. In fact, they both belong to the region where $\dot{X} > 0$ and $\dot{Z} < 0$ in their respective ranges of p .

2.3. Periodic orbits

In this section we continue investigating the limit sets of the trajectories. Let us see in which intervals of p the dynamical systems (2.9) and (2.11) allow the existence of periodic orbits.

The Poincaré-Bendixson theorem [18] for planar autonomous systems says that the only admissible ω and α limits of bounded trajectories are either a stationary point or a periodic orbit.

Observe that we have the following ordering for the exponents $p_{-}^{p,a} < p_{\Delta}^a < p_{+}^{p,a}$ defined in (1.6).

Proposition 2.10 (Dulac's criterion). *Let $\lambda < \Lambda$. For \mathcal{M}^+ there are no periodic orbits of (2.9) when $1 < p \leq p_{\Delta}^a$ or $p > p_{+}^{p,a}$. For \mathcal{M}^- no periodic orbits of (2.11) exist if $1 < p < p_{-}^{p,a}$ or $p \geq p_{\Delta}^a$. Moreover,*

- i. *there are no periodic orbits strictly contained in $R_{\lambda}^+ \cup \ell_+$ (resp. $(R_{\Lambda}^+ \cup \ell_-)$), for any $p > 1$;*
- ii. *periodic orbits contained in $R_{\lambda}^- \cup \ell_+$ (resp. $R_{\Lambda}^- \cup \ell_-$) are admissible only at $p = p_{\pm}^{p,a}$. Also, no periodic orbits at $p_{\pm}^{p,a}$ can cross the concavity line ℓ_{\pm} twice.*

Proof. Define $\varphi(X, Z) = X^\alpha Z^\beta$, where $\beta = \frac{3-p}{p-1}$ and α as in (1.6). Set $\Phi(X, Z) = \partial_X(\varphi f) + \partial_Z(\varphi g)$, with f and g defined in (2.17). For \mathcal{M}^+ we have

$$\Phi(X, Z) = \begin{cases} X^\alpha Z^\beta \left[\alpha \left(X - (N - 2) + \frac{Z}{\lambda} \right) + \beta \left(N + a - pX - \frac{Z}{\lambda} \right) + (2 - p)X + 2 - \frac{Z}{\lambda} \right] & \text{in } R_\lambda^+, \\ X^\alpha Z^\beta \left[\alpha \left(X - (\tilde{N}_+ - 2) + \frac{Z}{\Lambda} \right) + \beta \left(\tilde{N}_+ + a - pX - \frac{Z}{\Lambda} \right) + (2 - p)X + 2 - \frac{Z}{\Lambda} \right] & \text{in } R_\lambda^-, \\ \varphi(X, Z)(p - 1)^{-1} \left[-p(N - 2) + (N + 2 + 2a) \right] & \text{in } R_\lambda^+, \\ \varphi(X, Z)(p - 1)^{-1} \left[-p(\tilde{N}_+ - 2) + (\tilde{N}_+ + 2 + 2a) \right] & \text{in } R_\lambda^-. \end{cases}$$

Both expressions are positive if $1 < p < \min(p_+^{p,a}, p_\Delta^a) = p_\Delta^a$; and both are negative if $p > \max(p_+^{p,a}, p_\Delta^a) = p_+^{p,a}$. Then one concludes by the same argument as in the classical Bendixson–Dulac criterion, see also [17, Theorem 3.1]. Indeed, the vector field $F = (f, g)$ is Lipschitz continuous in (X, Z) , so Green’s area formula for the domain D enclosed by a periodic trajectory applies such as

$$\int_{\partial D} \varphi \{f \, dZ - g \, dX\} = \int_D \Phi(X, Z) \, dXdZ = \int_{R_\lambda^+ \cap D} \Phi(X, Z) \, dXdZ + \int_{R_\lambda^- \cap D} \Phi(X, Z) \, dXdZ. \tag{2.24}$$

The RHS is nonzero for $p \in (1, p_\Delta^a) \cup (p_+^{p,a}, \infty)$, but the LHS is zero because $dX = f dt, dZ = g dt$. Further, at p_Δ^a one has $\Phi = 0$ in R_λ^+ and so the first integral in (2.24) (in the RHS) vanishes, while the second one is positive. For \mathcal{M}^- the computations are similar by using that $\min(p_-^{p,a}, p_\Delta^a) = p_-^{p,a}$ and $\max(p_-^{p,a}, p_\Delta^a) = p_\Delta^a$.

Next we look at the interval $[p_\Delta^a, p_+^{p,a}]$ for \mathcal{M}^+ . Note that Poincaré-Bendixson theorem guarantees the existence of a stationary point in the domain D inside a periodic orbit. Since the only admissible stationary point in the interior of $1Q$ is $M_0 \in R_\lambda^-$ for $p > p_+^{s,a}$, while for $p \leq p_+^{s,a}$ M_0 is not an option (see Figure 2), then (i) follows.

To prove (ii) let us observe that if a periodic orbit is contained in $R_\lambda^- \cup \ell_+$ then by Proposition 2.5 (1) it may intersect the line ℓ_+ only at one point, namely the point P . Hence we can repeat the previous argument, neglecting the integral expression in R_λ^+ . Then we get that there are no periodic orbits in $R_\lambda^- \cup \ell_+$ for every $p \neq p_+^{p,a}$. To finish, if a periodic orbit existed which crossed twice the line ℓ_+ at $p_+^{p,a}$, then the first integral of (2.24) (in the RHS) would be positive, while the second one is equal to zero because $\Phi = 0$ in R_λ^- . The case for \mathcal{M}^- and $p_-^{p,a}$ is analogous. \square

Notably Dulac’s criterion brings out the critical exponents p_Δ^a and $p_\pm^{p,a}$. They correspond to critical exponents for the two Laplacian operators Δ_N and $\Delta_{\tilde{N}_\pm}$, in dimensions N and \tilde{N}_\pm .

Other limit cycles θ are admissible by the dynamical system as far as they cross ℓ_\pm twice. They do appear for \mathcal{M}^\pm as we shall see in Sections 4, 5. This happens because Dulac’s criterion is inconclusive in a whole interval when $\lambda < \Lambda$. Formally, the Pucci problem opens space for new periodic orbits in order to appropriately glue both Laplacian operators.

2.4. A Priori bounds

We prove ahead important bounds for trajectories of (2.9) or (2.11) which are defined for all t in intervals of type $(\hat{t}, +\infty)$ or $(-\infty, \hat{t})$.

By Poincaré-Bendixson theorem, if a trajectory of (2.9) or (2.11) does not converge to a stationary point neither to a periodic orbit, either forward or backward in time, then it necessarily blows up in that direction. In the next propositions we prove that a blow up may only occur in finite time. The first result is obtained in the first quadrant.

Proposition 2.11. *Let τ be a trajectory of (2.9) or (2.11) in 1Q, with $\tau(t) = (X(t), Z(t))$ defined for all $t \in (\hat{t}, +\infty)$, for some $\hat{t} \in \mathbb{R}$. Then*

$$X(t) < \tilde{N}_\pm - 2, \quad \text{for all } t \geq \hat{t}. \tag{2.25}$$

If instead, τ is defined for all $t \in (-\infty, \hat{t})$, for some $\hat{t} \in \mathbb{R}$, then

$$Z(t) < \lambda(N + a) \text{ in the case of } \mathcal{M}^+, \quad Z(t) < \Lambda(N + a) \text{ for } \mathcal{M}^-, \quad \text{for all } t \leq \hat{t}. \tag{2.26}$$

In particular, if a global trajectory is defined for all $t \in \mathbb{R}$ in 1Q then it remains inside the box $(0, \tilde{N}_+ - 2) \times (0, \lambda(N + a))$ in the case of \mathcal{M}^+ , while it stays in $(0, \tilde{N}_- - 2) \times (0, \Lambda(N + a))$ for \mathcal{M}^- .

Proof. Let us first prove (2.25) when the operator is \mathcal{M}^+ , $\tilde{N}_+ \leq N$. Arguing by contradiction we assume that for some $t_1 \geq \hat{t}$ we have $X(t_1) \geq \tilde{N}_+ - 2$. Notice that $\dot{X} > 0$ on the half line $L^+ = \{(X, Z) : X = \tilde{N}_+ - 2\} \cap 1Q$, see (2.18). Therefore $X(t) > X(t_1) \geq \tilde{N}_+ - 2$ for all $t > t_1$.

We claim that $X(t) \rightarrow +\infty$ as $t \rightarrow +\infty$. To see this, first notice that Z is bounded from t_1 on, since $\dot{Z} < 0$ to the right of L^+ , see (2.19). If we had $X(t) \leq C$ for some $C > 0$ for $t \geq t_1$, then τ would be a bounded trajectory from t_1 on. Then by Poincaré-Bendixson it should converge to a stationary point as $t \rightarrow +\infty$. Notice that periodic orbits are not allowed to the right of L^+ by the direction of the vector field, see Figure 1. This proves the claim, since no stationary points exist on the right of L^+ .

Thus, we can pick a time t_2 such that $X(t_2) > N - 2 \geq \tilde{N}_+ - 2$. Again by monotonicity, $X(t) > \tilde{N}_+ - 2$ for all $t \geq t_2$.

Then we have two cases: either the trajectory τ reaches the region R_λ^- for some $t_3 \geq t_2$, or it stays in R_λ^+ for all time. If the first holds, then $\tau(t)$ remains there for all $t \geq t_3$, since $\dot{Z} < 0$ to the right of L^+ , see Figure 1. Observe that the first equation in (2.9) yields $\frac{\dot{X}}{X[X - (\tilde{N}_+ - 2)]} \geq 1$. Moreover,

$$\frac{(\tilde{N}_+ - 2)\dot{X}}{X[X - (\tilde{N}_+ - 2)]} = \frac{\dot{X}}{X - (\tilde{N}_+ - 2)} - \frac{\dot{X}}{X} = \frac{d}{dt} \ln \left(\frac{X(t) - \tilde{N}_+ + 2}{X(t)} \right) \text{ for all } t \geq t_3. \tag{2.27}$$

Therefore, by integrating (2.27) in the interval $[t_3, t]$ we get

$$X(t) \geq \frac{\tilde{N}_+ - 2}{1 - ce^{(\tilde{N}_+ - 2)(t - t_3)}}, \text{ where } c = 1 - \frac{\tilde{N}_+ - 2}{X(t_3)} \in (0, 1), \tag{2.28}$$

and in particular X blows up in the finite time $t_1 = t_3 + \frac{\ln(1/c)}{N_+ - 2}$.

If instead τ stays in R_λ^+ from t_2 on, then the same computations developed with N in place of \tilde{N}_+ imply, using the first equation in (2.9), that X blows up in finite time. Both ways one gets a contradiction.

Let us now prove (2.26) for \mathcal{M}^+ . Notice that $\dot{Z} < 0$ in the region above the line $Z = \lambda(N + a)$ which is contained in R_λ^+ , see Figure 1. Now, if $Z = \lambda(N + a)$ occurs at some point for the orbit τ , then in particular there is some t_0 such that $Z > \lambda(N + a)$ for all $t \leq t_0$, thus $\dot{Z} \leq Z(N + a - Z/\lambda)$. In particular τ remains in the region R_λ^+ up to the time t_0 . Hence,

$$\frac{\lambda(N + a)\dot{Z}}{\lambda(N + a) - Z} = \frac{\dot{Z}}{Z} - \frac{\dot{Z}}{Z - \lambda(N + a)} = \frac{d}{dt} \left\{ \ln \left(\frac{Z(t)}{Z(t) - \lambda(N + a)} \right) \right\} \text{ for all } t \leq t_0.$$

integration in $[t, t_0]$ as before gives us that the trajectory blows up in finite time.

The proof of (2.25) and (2.26) for the operator \mathcal{M}^- is analogous if one uses $\tilde{N}_- \geq N$. □

Now, a similar argument of Proposition 2.11 allows to characterize all the orbits in 3Q.

Proposition 2.12. *Every orbit of (2.10) or (2.12) in 3Q blows up in finite time, backward and forward. The vector field in there always points to the right and down, with $\dot{X} > 0$ and $\dot{Z} < 0$.*

Proof. Recall that in 3Q we have $X, Z < 0$. Let us consider \mathcal{M}^+ . Hence, by the first equation in (2.10) one gets $\dot{X} \geq X(X - (\tilde{N}_+ - 2))$, which is positive. Similarly, by the second equation in (2.12) one figures out that $\dot{Z} \leq Z(\tilde{N}_+ + a - Z/\Lambda)$, which is now negative. Then integration as in (2.27), (2.28) gives us the result. For \mathcal{M}^- it is analogous. □

3. Classification of solutions

In this section we classify the solutions of the second order equations (P_+) and (P_-) and we show that this induces a classification of the orbits of the dynamical systems (2.9)-(2.12). We investigate three kinds of solutions to (P_+) and (P_-) : regular solutions, singular solutions, and exterior domain solutions. The denomination of the latter ones will be clarified in Section 3.3.

We begin by characterizing the blow ups admissible in the first quadrant.

Proposition 3.1 (Blow-up types in 1Q). *Let u be a positive solution of (P_+) or (P_-) in an interval (R_1, R_2) , $0 < R_1 < R_2$, and $\tau = (X, Z)$ be a corresponding trajectory of (2.9) or (2.11) lying in 1Q through the transformation (2.4). Then the following holds:*

- i. *there exists $r_1 \in (R_1, R_2)$ such that $u'(r_1) = 0 \iff$ there exists $t_1 \in \mathbb{R}$ such that $Z(t) \rightarrow +\infty$ as $t \rightarrow t_1^+$. In addition, $X(t) \rightarrow 0$ as $t \rightarrow t_1^+$;*
- ii. *there exists $r_2 \in (R_1, R_2)$ such that $u(r_2) = 0 \iff$ there exists $t_2 \in \mathbb{R}$ such that $X(t) \rightarrow +\infty$ as $t \rightarrow t_2^-$. Further, $Z(t) \rightarrow 0$ as $t \rightarrow t_2^-$.*

Moreover, no other blow-up types other than those of (i) and (ii) are admissible for τ in 1Q.

Proof. Let us first observe that u and u' can never be zero at the same point r_1 . Otherwise, by the uniqueness of the initial value problem we would have $u \equiv 0$ in a neighborhood of r_1 , which is not possible by the strong maximum principle.

(i) Assume that there exists $r_1 > 0$ such that $u'(r_1) = 0$. Thus $u(r_1) > 0$ and by (2.4) it is easy to deduce the limits of $X(t)$ and $Z(t)$ as $t \rightarrow t_1^+$, for $t_1 = \ln(r_1)$. Viceversa, if $Z(t) \rightarrow +\infty$ as $t \rightarrow t_1^+$, by (2.4) we immediately get $u'(r) \rightarrow 0$ as $r \rightarrow r_1 = e^{t_1}$, because u is continuous in (R_1, R_2) . This in turn gives that $X(t) \rightarrow 0$ as $t \rightarrow t_1^+$, and no other asymptote parallel to the Z axis is admissible.

(ii) Suppose that $u(r_2) = 0$ for some $r_2 > R_1$. Then $u'(r_2) < 0$ and by (2.4) we easily obtain the behavior of X and Z as $t \rightarrow t_2^-$, where $t_2 = \ln(r_2)$. Viceversa if $X(t) \rightarrow +\infty$ as $t \rightarrow t_2^-$ then necessarily $u(r) \rightarrow 0$ as $r \rightarrow r_2 = e^{t_2}$, because u' is continuous in (R_1, R_2) . Thus $Z(t) \rightarrow 0$ as $t \rightarrow t_2^-$ as before.

The arguments above also show that, in finite time, no other blow-ups are possible for τ in 1Q. Indeed, as soon as X or Z tends to infinity, then u or u' vanishes at a positive radius. Recall that a blow up in infinite time is not admissible by Proposition 2.11. □

Corollary 3.2. *Let u be a solution of (P_+) or (P_-) , and τ be a corresponding trajectory of (2.9) or (2.11) starting above the line ℓ_{\pm} in 1Q. Then u changes concavity at least once.*

Proof. Consider the \mathcal{M}^+ operator, for \mathcal{M}^- is the same. If u never changed concavity, then $\tau = (X, Z)$ would remain inside the region R_{λ}^+ for all time. By Lemma 2.4 and Proposition 2.10 there are no stationary points or periodic orbits in R_{λ}^+ . Recall that $\dot{X} > 0, \dot{Z} < 0$ in R_{λ}^+ , see Figure 1. Then τ must blow up at a finite forward time \hat{t} such that $X(t) \rightarrow +\infty$ and $Z(t) \rightarrow Z_1$ as $t \rightarrow \hat{t}$, for some $Z_1 > \lambda(N - 1) > 0$. But this blow-up is not admissible by Proposition 3.1. □

Remark 3.3 (Blow up in 3Q). Every orbit $\tau = (X, Z)$ of (2.10) or (2.12) in 3Q verifies $X(t) \rightarrow 0$ and $Z(t) \rightarrow -\infty$ as $t \rightarrow t_1^-$ for some $t_1 \in \mathbb{R}$ such that $r_1 = e^{t_1}$ and $u'(r_1) = 0$. Moreover, $X(t) \rightarrow -\infty$ and $Z(t) \rightarrow 0$ as $t \rightarrow t_3^+$ for some $t_3 \in \mathbb{R}$ where $r_3 = e^{t_3}$ and $u(r_3) = 0$.

3.1. Regular solutions

Let us consider the following initial value problem:

$$u'' = M_{\pm}(-r^{-1}(N - 1) m_{\pm}(u') - r^a|u|^{p-1}u), \quad u(0) = \gamma, \quad u'(0) = 0, \quad \gamma > 0, \quad (3.1)$$

where M_{\pm} and m_{\pm} are defined in (2.2), (2.3).

By regular solution we mean a solution $u = u_p$ of (3.1) which is twice differentiable for $r > 0$, and C^1 up to 0. We denote by R_p , with $R_p \leq +\infty$, the radius of the maximal interval $[0, R_p)$ where u is positive.

Hence, in such interval u is a solution of (P_{\pm}) . Obviously, if $R_p = +\infty$ then u corresponds to a radial positive solution of (1.1) for $\Omega = \mathbb{R}^N$. When $R_p < +\infty$ it gives a positive solution of the Dirichlet problem (1.1), (1.3) in the ball $\Omega = B_{R_p}$.

Remark 3.4. Given a regular positive solution $u = u_p$ in $[0, R_p)$ satisfying (3.1) for some $\gamma > 0$, then the rescaled function $v(r) = \tau u(\tau^{\frac{1}{2}}r)$, for α as in (1.6) and $\tau > 0$, is still a positive solution of the same equation in $[0, \tau^{-\frac{1}{2}}R_p)$ with initial value $v(0) = \tau\gamma$, see also [6, Lemma 2.3].

If u is defined in the whole interval $[0, +\infty)$, thus there is a family of regular solutions obtained via $v = v_\tau$, for all $\tau > 0$. In this case we say that u is *unique* up to scaling.

On the other hand, a solution in the ball of radius R_p automatically produces a solution for an arbitrary ball, by properly choosing the parameter $\tau > 0$.

Remark 3.5. Note that, by rescaling a fast decaying solution, we get infinitely many fast decaying solutions which give a different value for the constant c in Definition 1.1(i), see (3.7). The same happens for the $(\tilde{N} - 2)$ -blowing up solution in Definition 1.5(i). Instead it is easy to see that for the slow decaying solutions or α -blowing up solutions in definitions 1.1(ii) and 1.5(ii), the constant c is independent of the rescaling, see (3.8).

Now, using the transformation (2.4) our goal is to characterize the regular solutions of (P_+) or (P_-) as trajectories of the dynamical systems (2.9) and (2.11) in the first quadrant.

Proposition 3.6. *Let $u = u_p$ be any positive regular solution of (P_+) (resp. (P_-)). Then the corresponding trajectory belongs to 1Q and is the unique trajectory of (2.9) (resp. (2.11)) whose α -limit is N_0 .*

Proof. The proof is the same for both operators \mathcal{M}^\pm . The solution $u = u_p$ satisfies $\lim_{r \rightarrow 0} u(r) = u(0) = \gamma$ and $\lim_{r \rightarrow 0} u'(r) = u'(0) = 0$, for some $\gamma > 0$. In terms of the trajectory (X, Z) , by the definition of X in (2.4) we easily find

$$\lim_{t \rightarrow -\infty} X(t) = 0. \quad (3.2)$$

Moreover, in the simpler case when $a = 0$ we have

$$\lim_{t \rightarrow -\infty} Z(t) = \lim_{r \rightarrow 0} \frac{-ru^p(r)}{u'(r)} = -\gamma^p \lim_{r \rightarrow 0} \frac{r}{u'(r) - u'(0)} = -\frac{\gamma^p}{u''(0)} \in (0, +\infty),$$

since it is easier to check from the equation that $u''(0) < 0$. When $a \neq 0$ we need some other argument to show that $Z(t)$ has a finite limit as $t \rightarrow -\infty$. First let us show that

$$\text{there exists } R_1 > 0 \text{ such that } u' \neq 0 \text{ for all } r \in (0, R_1). \quad (3.3)$$

If this was not true, then there would exist a sequence of positive radii $r_n \rightarrow 0$ such that $u'(r_n) = 0$. By the mean value theorem this yields the existence a sequence $s_n \in (r_n, r_{n+1})$ such that $u''(s_n) = 0$. Thus, since u' cannot be identically zero in a neighborhood of 0 by the equation (P_+) , then u changes infinitely many times its concavity in a neighborhood of 0.

In terms of the dynamical system, say (2.9) for \mathcal{M}^+ , this means that the respective trajectory intersects the line ℓ_+ (see (2.14)) more than once as $t \rightarrow -\infty$. In particular it should pass from R_λ^+ to R_λ^- infinitely many times, which, by Proposition 2.5 (1), may

only occur at $X(s_n) > \frac{1+a}{p}$ for $a > -1$. This contradicts the fact that $X(s_n) \rightarrow 0$ for large n from (3.2).

By (3.3) we have that $Z(t)$ is well defined in some interval $(-\infty, \hat{t})$ so that (2.26) in Proposition 2.11 yields $Z(t) < \lambda(N + a)$ for all $t < \hat{t}$. Hence $\lim_{t \rightarrow -\infty} Z(t) < +\infty$. Note that the trajectory cannot belong to $3Q$ when blow-up in finite time occurs both backward and forward, by Proposition 2.12. Moreover, it cannot converge to O by Propositions 2.5 (2) and 2.7(1). Indeed, the unstable manifold at O is on the Z axis which cannot correspond to the solution u in any interval $(0, r)$. Hence it converges to N_0 , independently of the initial datum $\gamma > 0$.

Remark 3.7. Thus, by Propositions 3.6 and 2.8 one concludes that a regular solution u_p , corresponds to the unique trajectory Γ_p labeled in (2.22), for all $p > 1$. Here, Γ_p is defined in a maximal interval $[0, T_p), T_p = \ln R_p \leq +\infty$.

Note that the fact that Γ_p does not depend on the initial datum of u_p is not a surprise since we already observed that the change of initial datum is equivalent to rescaling the radius. This, in turn, is equivalent to shifting the time for the systems (2.9), (2.11), which does not produce any change in the trajectory since the system is autonomous.

We now prove the monotonicity and concavity properties of the regular solutions, deriving them directly by the dynamical systems (2.9) or (2.11), and not from the second order ODEs.

Proposition 3.8. *All regular solutions u of (P_+) or (P_-) are concave in an interval $(0, r_0)$ for some $r_0 > 0$ and change concavity at least once. Moreover, they are decreasing as long as they remain positive. Further, $u'(r) = O(-r^{1+a})$ and $u''(r) = O(-r^a)$ as $r \rightarrow 0$, for $a > -1$.*

Proof. Let us consider the \mathcal{M}^+ case; for \mathcal{M}^- is the same. By Proposition 3.6 the corresponding trajectory Γ_p starts at $-\infty$ from the stationary point N_0 and enters the region R_λ^+ which is above the concavity line ℓ_+ ; see Proposition 2.7 (item 2) and Remark 2.9. Then we immediately deduce that u is concave near $r=0$ and changes concavity at least once; see Corollary 3.2. Next, since $1Q$ is invariant by the flow and corresponds to positive decreasing solutions of (P_+) we get the monotonicity claim. For the asymptotics one computes

$$-\gamma^p \lim_{r \rightarrow 0} \frac{r^{1+a}}{u'(r)} = \lim_{t \rightarrow -\infty} Z(t) = \lambda(N + a), \text{ from which } \lim_{r \rightarrow 0} \frac{u''}{r^a} = -\frac{\gamma^p}{\lambda} \frac{a + 1}{N + a}. \quad \square$$

Given an exponent $p > 1$, for a regular solution u_p of (P_+) or (P_-) , which is positive in $[0, +\infty)$, Definition 1.1 holds according to its behavior at infinity. Taking into account that u_p is unique, up to rescaling, as in [6] we define the following sets:

$$\begin{aligned} \mathcal{F} &= \{p > 1 : u_p \text{ is fast decaying}\}; \\ \mathcal{S} &= \{p > 1 : u_p \text{ is slow decaying}\}; \\ \mathcal{P} &= \{p > 1 : u_p \text{ is pseudo-slow decaying}\}. \end{aligned} \tag{3.4}$$

Then we add the set

$$\mathcal{C} = \{p > 1 : (P_\pm) \text{ has a solution } u_p \text{ with } u_p(R_p) = 0\}, \tag{3.5}$$

where, as before, R_p is the radius of the maximal positivity interval for u_p . We characterize the previous sets in terms of the orbits of the dynamical systems (2.9) or (2.11).

Proposition 3.9. *In terms of the dynamical systems (2.9) or (2.11), the previous sets can be equivalently defined as follows:*

$$\begin{aligned} \mathcal{F} &= \{p > 1 : \omega(\Gamma_p) = A_0\}; & \mathcal{S} &= \{p > 1 : \omega(\Gamma_p) = M_0\}; \\ \mathcal{P} &= \{p > 1 : \omega(\Gamma_p) \text{ is a periodic orbit around } M_0\}; \\ \mathcal{C} &= \{p > 1 : \lim_{t \rightarrow T} X(t) = +\infty \text{ and } \lim_{t \rightarrow T} Z(t) = 0 \text{ for some } T > 0\}, \end{aligned} \tag{3.6}$$

where $\Gamma_p(t) = (X(t), Z(t))$ is as in Remark 3.7. In particular, $(1, +\infty) = \mathcal{C} \cup \mathcal{F} \cup \mathcal{P} \cup \mathcal{S}$.

Proof. The proof is the same for both operators \mathcal{M}^\pm , i.e. for both systems (2.9) and (2.11).

In the case of the sets \mathcal{F}, \mathcal{S} and \mathcal{P} , u_p (as in Proposition 3.6) is positive in $(0, +\infty)$ which implies that Γ_p is defined for all $t \in \mathbb{R}$. By Proposition 2.11 this trajectory is bounded, and so by Poincaré-Bendixson theorem it converges as $r \rightarrow +\infty$ either to a stationary point or to a periodic orbit. In the first case only A_0 and M_0 are admissible. Moreover, via the transformation (2.4),

$$\omega(\Gamma_p) = A_0 \iff \lim_{r \rightarrow +\infty} u(r)r^{\tilde{N}_\pm - 2} = C_{p,\gamma} \quad \text{for some } C_{p,\gamma} > 0, \tag{3.7}$$

and $u = u_p$ has fast decay at $+\infty$. On the other hand u is slow decaying at $+\infty$ when

$$\omega(\Gamma_p) = M_0 \iff \lim_{r \rightarrow +\infty} u(r)r^\alpha = C_p \quad \text{with } C_p = (X_0 Z_0)^{\frac{1}{p-1}}, \tag{3.8}$$

where $M_0 = (X_0, Z_0)$ is given explicitly in Lemma 2.4.

Indeed, (3.8) comes from the identity $X(t)Z(t) = r^{2+a}u^{p-1}$ for all $t \in \mathbb{R}$. On the other hand,

$$\lim_{t \rightarrow +\infty} X(t) = \tilde{N}_\pm - 2 \text{ is equivalent to } \frac{d}{dr} \ln(u(r)) = \frac{u'(r)}{u(r)} \sim -\frac{\tilde{N}_\pm - 2}{r} \text{ as } r \rightarrow +\infty.$$

Then, integration in $[r_0, r]$ for a fixed large r_0 implies (3.7) with $C_{p,\gamma} = u(r_0)r_0^{\tilde{N}_\pm - 2}$, $u = u_{p,\gamma}$. Now, by rescaling, the function $v = v_\tau$ in Remark 3.4 satisfies

$$\lim_{t \rightarrow +\infty} v(r)r^{\tilde{N}_\pm - 2} = \tau^{1 - \frac{\tilde{N}_\pm - 2}{\alpha}} C_{p,\gamma} \text{ under (3.7); } \lim_{r \rightarrow +\infty} v(r)r^\alpha = C_p \text{ under (3.8).}$$

Thus C_p is independent of the initial value $\gamma > 0$ in (3.1).

Finally, assume that $\omega(\Gamma_p)$ is a periodic orbit θ . Note that the region inside θ is bounded, and by Poincaré-Bendixson theorem it must contain M_0 . Using again $XZ = r^{2+a}u^{p-1}$ one defines

$$c_1^{p-1} := \inf_{t \in \mathbb{R}} \{X(t)Z(t) : (X, Z) \in \theta\}; \quad c_2^{p-1} = \sup_{t \in \mathbb{R}} \{X(t)Z(t) : (X, Z) \in \theta\}.$$

Therefore we deduce

$$\omega(\Gamma_p) = \theta \iff 0 < c_1 = \liminf_{r \rightarrow +\infty} u(r)r^\alpha < \limsup_{r \rightarrow +\infty} u(r)r^\alpha = c_2. \tag{3.9}$$

Now we consider the set \mathcal{C} . The corresponding trajectory Γ_p cannot be defined for all time since $u(R_p) = 0$. So it must blow up at the finite time $T_p = \ln(R_p)$ by [Proposition 3.1](#).

Viceversa if $\mathcal{F}, \mathcal{S}, \mathcal{P}, \mathcal{C}$ are defined in terms of the property of the trajectory Γ_p of the dynamical system then they give exactly the same sets as defined for u_p , by using [\(2.4\)](#) and [\(2.13\)](#), and arguing in a similar way.

Remark 3.10 (\mathcal{C} is open). When $p \in \mathcal{C}$, the trajectory Γ_p crosses the line $L = \{(X, Z) : X = \tilde{N}_+ - 2\}$ and next blows up in finite time. This property is preserved for p' close to p .

3.2. Singular solutions

As mentioned in Section 1, by singular solution we mean a radial solution u of [\(1.1\)](#) satisfying [\(1.2\)](#). Hence $u = u(|x|) = u(r)$ is a singular solution of (P_\pm) satisfying $\lim_{r \rightarrow 0} u(r) = +\infty$. It may be either positive for all $r \in (0, +\infty)$, or be equal to zero at a certain radius $R > 0$. In the latter case it produces a solution in $B_R \setminus \{0\}$.

In terms of the systems [\(2.9\)](#), [\(2.11\)](#), this means that the corresponding trajectory, say Σ_p , will be defined either in \mathbb{R} or in an interval $(-\infty, T)$, for some $T < \infty$. Under the latter, as in the characterization of \mathcal{C} in [\(3.6\)](#) we have that Σ_p blows up forward in a finite time $T < +\infty$. Otherwise, by [Proposition 2.11](#) the global trajectory Σ_p is contained in the box

$$\mathcal{Q}^+ = (0, \tilde{N}_+ - 2) \times (0, \lambda(N + a)) \text{ for } \mathcal{M}^+; \mathcal{Q}^- = (0, \tilde{N}_- - 2) \times (0, \Lambda(N + a)) \text{ for } \mathcal{M}^-.$$

Then the α and ω limits can be either a periodic orbit or a stationary point.

We point out that Σ_p cannot converge to N_0 , neither backward nor in forward time, because the stable direction at N_0 is the Z axis, while the unstable direction corresponds to the regular trajectory Γ_p , for all $p > 1$. So all possible α and ω limits of Σ_p are M_0 , A_0 , or a periodic orbit.

By the analysis of the stationary points M_0 and A_0 , and of the periodic orbits given in Section 2.3, the α and ω limits of Σ_p depend on the exponent p . Then a classification of the singular solutions, according to Definition 1.5 can be easily formulated in terms of the dynamical systems [\(2.9\)](#), [\(2.11\)](#), analogously to [Proposition 3.9](#). Obviously if Σ_p is defined in \mathbb{R} they are also classified according to the behavior at $+\infty$, as in Definition 1.1. Here we just emphasize that, as for the regular solutions, the so called pseudo-blowing up solutions, see Definition 1.5 (iii), may only occur at the values of p for which Σ_p has a periodic orbit as α -limit.

3.3. Exterior domain solutions

By exterior domain solution we mean a solution u of (P_+) or (P_-) defined in an interval $[\rho_0, \rho]$, for $\rho_0 \in (0, +\infty)$ and $\rho \leq +\infty$, and verifying the Dirichlet condition

$u(\rho_0) = 0$. Fixing $\rho_0 = 1$, then u satisfies the initial value problem

$$u(1) = 0, \quad u'(1) = \delta, \quad \text{for some } \delta > 0. \quad (3.10)$$

The equations (P_+) , (P_-) , together with (3.10) were studied in [16] and [18]. It was shown that for any $p > 1$ and for each $\delta > 0$ there exists a unique solution $u = u_\delta$ defined in a maximal interval $(1, \rho_\delta)$ where u is positive, $1 < \rho_\delta \leq +\infty$. Moreover, there exists a unique $\mu_\delta \in (1, \rho_\delta)$ such that

$$u'_\delta(r) > 0 \text{ for } r < \mu_\delta, \quad u'_\delta(\mu_\delta) = 0, \quad u'_\delta(r) < 0 \text{ for } r > \mu_\delta. \quad (3.11)$$

If $\rho_\delta = +\infty$ we get a positive radial solution in the exterior of the ball B_1 . In the second case, a positive solution in the annulus $(1, \rho_\delta)$ is produced. Note that equations (P_+) , (P_-) together with (3.10) are not invariant by rescaling.

Now we would like to describe the trajectories of the dynamical systems (2.9)-(2.11) which correspond to u_δ through the variables X, Z in (2.4).

Proposition 3.11. *Let $p > 1$ and $u_\delta = u_{\delta,p}$ be a positive solution of (P_+) (resp. (P_-)) satisfying (3.10). Then there exists a unique trajectory $\Xi_{\delta,p}$ in $1Q$ for the system (2.9) (resp. (2.11)) which blows up backward in a finite time t_δ . More precisely, if $\Xi_{\delta,p}(t) = (X(t), Z(t))$ then*

$$\lim_{t \rightarrow t_\delta^+} Z(t) = +\infty \text{ and } \lim_{t \rightarrow t_\delta^+} X(t) = 0, \quad (3.12)$$

where $t_\delta = \ln(\mu_\delta)$, μ_δ given in (3.11). The trajectory $\Xi_{\delta,p}$ corresponds, after the transformation (2.4) to the restriction of u_δ to the interval $I_\delta = (\mu_\delta, \rho_\delta)$, $\rho_\delta \leq +\infty$.

Proof. The proof works for both operators \mathcal{M}^\pm . We fix p and δ . By (3.11), u_δ is positive and decreasing in the interval I_δ . Then, after (2.4), to u_δ corresponds a unique trajectory $\Xi_\delta = \Xi_{\delta,p}$ contained to $1Q$ and is defined for all $t \in (t_\delta, \ln(\rho_\delta))$. Thus, by Proposition 3.1 we get (3.12). \square

When $\rho_\delta = +\infty$ we can classify the solutions accordingly to their behavior at $+\infty$, i.e. u_δ is fast, slow, or pseudo-slow decaying via the limits (i)-(iii) as $r \rightarrow +\infty$ in Definition 1.1.

4. The \mathcal{M}^+ case

In this section we study the solutions of the equations involving the Pucci \mathcal{M}^+ operator. Hence we refer to the dimension-like parameter \tilde{N}_+ and the relevant exponents for \mathcal{M}^+ defined in (1.5) and (1.6), as well as their ordering: $\max\{p_+^{s,a}, p_\Delta^a\} \leq p_+^{p,a}$.

4.1. Some properties of regular trajectories

We first consider the case of a regular solution of (P_+) whose corresponding trajectory for the system (2.9) will be denoted by $\Gamma_p = \Gamma_p(t)$ as in Section 3.1. We also keep the other notations already introduced, in particular for the sets $\mathcal{F}, \mathcal{S}, \mathcal{P}, \mathcal{C}$ defined in (3.4), (3.5), and Proposition 3.9.

Lemma 4.1. *For any $p > 1$, with $\Gamma_p = (X_p, Z_p)$, we have:*

- i. if Γ_p reaches the line ℓ_1^+ (see (2.18)) at some t_0 with $X_p(t_0) \geq \alpha$, then $p \in \mathcal{S} \cup \mathcal{P}$, i.e. the corresponding solutions u_p of (P_+) are either slow decaying or pseudo-slow decaying. In the latter case Γ_p crosses ℓ_1^+ and ℓ_2^+ (see (2.19)) infinitely many times;
- ii. if Γ_p does not intersect the line ℓ_1^+ , then it intersects the concavity line ℓ_+ exactly once. Moreover, $\dot{X}_p > 0$ and $\dot{Z}_p < 0$ for all time. In particular this happens for $p \in \mathcal{F} \cup \mathcal{C}$.

Proof. We recall that Γ_p starts at $-\infty$ from the stationary point N_0 and must cross the concavity line ℓ_+ at least once, see Proposition 3.8.

(i) If Γ_p reaches ℓ_1^+ for $X_p(t_0) = \alpha$ then clearly $\lim_{t \rightarrow +\infty} \Gamma_p(t) = M_0$, whenever M_0 belongs to 1Q (see Figure 1). If instead $X_p(t_0) > \alpha$, by taking into account Proposition 2.5 (3) (see again Figure 1) and that Γ_p cannot self intersect, we have that Γ_p is contained in a bounded region from which it cannot leave. Thus, by Poincaré-Bendixon theorem the ω -limit of Γ_p is either M_0 or a periodic orbit θ which contains M_0 in its interior. In the latter case Γ_p goes around θ clockwise according to the direction of the vector field, intersecting ℓ_1^+ and ℓ_2^+ infinitely many times.

(ii) If Γ_p does not intersect ℓ_1^+ then it cannot turn back and cross the concavity line ℓ_+ another time because of the direction of the flow. Moreover, it can neither intersect nor be tangent to the line ℓ_2^+ where $\dot{Z} = 0$, since a C^1 trajectory of (2.9) may only intersect the line ℓ_2^+ transversely by passing from left to right, see Proposition 2.5 (3) and Figure 1. This fact and item (i) conclude the final assertion. □

The next proposition is crucial to study the behavior of Γ_p for different values of p .

Proposition 4.2. *Assume that $p_1 \in \mathcal{F} \cup \mathcal{C}$, and let Γ_{p_2} be any regular trajectory with $p_2 \neq p_1$. Then Γ_{p_1} and Γ_{p_2} can never intersect.*

Proof. Both Γ_{p_1} and Γ_{p_2} have their α -limit at the stationary point N_0 which is a saddle point. By Proposition 2.7(2) the tangent unstable directions for Γ_{p_1} and Γ_{p_2} at N_0 are given, respectively, by

$$Z = -\frac{p_1 \lambda(N+a)}{N+2+2a} X \quad \text{and} \quad Z = -\frac{p_2 \lambda(N+a)}{N+2+2a} X. \tag{4.1}$$

Assume by contradiction that $\Gamma_{p_1}(t) = (X_1(t), Z_1(t))$ and $\Gamma_{p_2}(t) = (X_2(t), Z_2(t))$ intersect. Let us denote by Q the first intersection point. Since the dynamical system (2.9) is autonomous, one may assume that the intersection happens at the same time for both trajectories, i.e. $Q = (X_1(t_0), Z_1(t_0)) = (X_2(t_0), Z_2(t_0))$.

To fix the ideas assume $p_1 < p_2$. Then, by (4.1), at least in a neighborhood of N_0 , Γ_{p_1} is above Γ_{p_2} because $X \rightarrow 0^+$ (from the right). Moreover, from (2.9) and Lemma 4.1(ii) we have

$$\dot{X}_1(t_0) = \dot{X}_2(t_0) > 0, \quad \dot{Z}_2(t_0) < \dot{Z}_1(t_0) < 0, \tag{4.2}$$

since only \dot{Z} depends on p . In particular Γ_{p_1} remains above Γ_{p_2} after intersecting. Thus the two trajectories must have the same tangent at the point Q , which is not possible by (4.2). The case $p_2 < p_1$ is analogous. □

From the previous results we immediately get that a fast decaying solution can exist for only one value of p .

Corollary 4.3. *There exists at most one p in the interval $(p_+^{s,a}, +\infty)$ such that $p \in \mathcal{F}$.*

Proof. Assume by contradiction that $p_1, p_2 \in \mathcal{F}$ for some $p_+^{s,a} < p_1 < p_2$. This means that the corresponding trajectories Γ_{p_1} and Γ_{p_2} both come out from N_0 at $-\infty$ and converge to A_0 at $+\infty$. We have already observed by (4.1) that Γ_{p_1} stays above Γ_{p_2} in a neighborhood of N_0 .

On the other hand, since A_0 is a saddle point for $p > p_+^{s,a}$, looking at the linear stable directions given by Proposition 2.7(3), we have that Γ_{p_1} and Γ_{p_2} arrive at A_0 with a reversed order; i.e. Γ_{p_2} is above Γ_{p_1} . This is because $X \rightarrow (\tilde{N}_+ - 2)$ from the left.

Hence, Γ_{p_1} and Γ_{p_2} should intersect, but this is not possible by Proposition 4.2. □

Another important consequence of Proposition 4.2 is the following result.

Corollary 4.4. *Let $p_0 \in \mathcal{F}$, $p_0 > p_+^{s,a}$, then $p \in \mathcal{C}$ for $p_+^{s,a} < p < p_0$, and $p \in \mathcal{P} \cup \mathcal{S}$ for $p > p_0$.*

Proof. If $p_0 \in \mathcal{F}$ then Γ_{p_0} cannot intersect any other regular orbit Γ_p for $p \neq p_0$, by Proposition 4.2. This means that if Γ_p is above or below Γ_{p_0} in a neighborhood of N_0 , it remains so for all time. Moreover, $p \notin \mathcal{F}$ for $p \neq p_0$ by Corollary 4.3.

Thus, if $p < p_0$, Γ_p lies above Γ_{p_0} and so cannot converge to $M_0 = M_0(p)$ neither to a periodic orbit around it, since the line ℓ_1^+ is below Γ_{p_0} . Notice that ℓ_1^+ does not depend on p . So $p \in \mathcal{C}$.

On the other hand, if $p > p_0$ then Γ_p lies below Γ_{p_0} and therefore cannot cross the line $L = \{(X, Z) : X = \tilde{N}_+ - 2\}$ in order to blow up in finite time. Hence $p \notin \mathcal{C}$ and so must be in $\mathcal{P} \cup \mathcal{S}$. □

4.2. The critical exponent

Our goal here is to define and characterize the critical exponent which will be proved to have all properties listed in Theorem 1.3.

We start by showing that \mathcal{S} and \mathcal{C} contain the intervals $(p_+^{p,a}, +\infty)$ and $(1, p_\Delta)$ respectively.

Proposition 4.5. *If $p > p_+^{p,a}$ then $p \in \mathcal{S}$.*

Proof. In case $p > p_+^{p,a}$, by Proposition 2.10 we know that there are no periodic orbits of the system (2.9), hence $p \notin \mathcal{P}$. Moreover, M_0 is a sink by Proposition 2.7. Let us show that $p \notin \mathcal{C} \cup \mathcal{F}$.

If $p \in \mathcal{C}$, then Γ_p crosses the line $L := \{(X, Z) : X = \tilde{N}_+ - 2\}$, and blows up in finite time. Then the region D enclosed by Γ_p , L and the X, Z axes form a bounded domain from which an orbit can only leave D forward in time through L . Thus, an orbit arriving at $M_0 \in D$ cannot go anywhere in backward time, giving a contradiction with Poincaré-Bendixon theorem.

If instead $p \in \mathcal{F}$, then the bounded set whose boundary is given by Γ_p together with the X and Z axes, is invariant and contains M_0 . Again the orbits arriving at M_0 cannot escape in backward time. Therefore $p \in \mathcal{S}$. □

Proposition 4.6. For $p \in (1, \max\{p_+^{s,a}, p_\Delta^a\})$ it holds that $p \in \mathcal{C}$.

Proof. If $1 < p < p_+^{s,a}$, then A_0 is a source and $M_0 \notin 1Q$. In particular there are no periodic orbits contained in $1Q$. Hence $p \notin \mathcal{F} \cup \mathcal{S} \cup \mathcal{P}$, so if $\max\{p_+^{s,a}, p_\Delta^a\} = p_+^{s,a}$ the proof is complete.

Assume $p_+^{s,a} < p_\Delta^a$. Then, at $p = p_+^{s,a}$ no periodic orbits are allowed by **Proposition 2.10**, whose proof also shows nonexistence of homoclinics at $A_0 = M_0$ (i.e. orbits τ with $\omega(\tau) = \alpha(\tau) = A_0$). Therefore, if we had $\omega(\Gamma_p) = A_0$ then the orbits which come out from A_0 could not go anywhere. Alternatively, see **Proposition 4.9**. In particular, $p_+^{s,a} \notin \mathcal{F} \cup \mathcal{S} \cup \mathcal{P}$.

Finally, if $p_+^{s,a} < p < p_\Delta^a$, then M_0 is a source by **Proposition 2.7**. The trajectory Γ_p cannot be bounded, otherwise it could only converge to A_0 as $t \rightarrow +\infty$. As in the proof of **Proposition 4.5** this would produce a contradiction, because the region D enclosed by Γ_p , and the X, Z axes would be an invariant set from which any trajectory issued from M_0 cannot exit. In any case Γ_p blows up in finite time, and so $p \in \mathcal{C}$. \square

By **Propositions 4.6** and **4.5** we have that the set \mathcal{C} is nonempty and bounded from above. Therefore we define

$$p_{a+}^* = \sup \mathcal{C} \tag{4.3}$$

and obviously $p_{a+}^* \in [\max\{p_+^{s,a}, p_\Delta^a\}, p_+^{p,a}]$. From now on we refer to p_{a+}^* as the critical exponent for the Pucci operator \mathcal{M}^+ with weight $|x|^a$. The next result characterizes p_{a+}^* .

Theorem 4.7. The number p_{a+}^* defined in (4.3) belongs to \mathcal{F} . Thus it is the only exponent in the equation (P_+) for which there exists a unique, up to scaling, fast decaying solution.

Moreover, if $\lambda < \Lambda$, then $p_{a+}^* \neq p_\Delta^a$, $p_{a+}^* \neq p_+^{p,a}$, and (1.7) holds. Further, $\mathcal{P} = (p_{a+}^*, p_+^{p,a}]$, and for any $p \in \mathcal{P}$ the corresponding trajectory Γ_p crosses the concavity line ℓ_+ infinitely many times.

Proof. First, $p_{a+}^* \notin \mathcal{C}$ i.e. \mathcal{C} does not have a maximum because \mathcal{C} is open, see **Remark 3.10**. By **Proposition 2.7** we know that M_0 is a source for every $p \in [p_\Delta^a, p_+^{p,a})$; and M_0 is a center at $p = p_+^{p,a}$. Whence $p \notin \mathcal{S}$ for all $p \in [p_\Delta^a, p_+^{p,a}]$, and in particular $p_{a+}^* \notin \mathcal{S}$. On the other side, if $p_{a+}^* \in \mathcal{P}$ then $\Gamma_{p_{a+}^*}$ would cross the line ℓ_1^+ by **Lemma 4.1(ii)**. Thus, by continuity with respect to p , the trajectory Γ_p should also cross ℓ_1^+ for p close to p_{a+}^* . But every trajectory Γ_p for $p \in \mathcal{C}$ does not cross ℓ_1^+ , by **Lemma 4.1**. Therefore $p_{a+}^* \notin \mathcal{P}$.

Hence p_{a+}^* belongs to \mathcal{F} and the trajectory $\Gamma_{p_{a+}^*}$ together with the X and Z axes enclose a bounded invariant region D which contains M_0 in its interior. Since M_0 is a source for $p \in [p_\Delta^a, p_+^{p,a})$, and a center for $p = p_+^{p,a}$, the set D contains periodic orbits which cross the line ℓ_+ twice. Indeed, the flow is subjected to Poincaré-Bendixson theorem, see **Figure 7**. This implies that p_{a+}^* can be neither p_Δ^a nor $p_+^{p,a}$ if $\lambda < \Lambda$, by **Proposition 2.10**. In fact, at p_Δ^a there are no periodic orbits at all, while at $p_+^{p,a}$ no periodic orbits cross ℓ_+ twice.

Note that we obtain (1.7) as long as $p_\Delta^a \geq p_+^{s,a}$. If $p_\Delta^a < p_+^{s,a}$ we still need to prove that $p_{a+}^* \neq p_+^{s,a}$. For instance this follows by well known Liouville results in [3]. Alternatively, a proof of this fact is accomplished in Proposition 4.9 which give nonexistence of entire positive solutions for $p = p_+^{s,a}$.

Next, by Corollary 4.4 we get $(p_{a+}^*, p_+^{p,a}) = \mathcal{P}$, since we have already observed that $[p_\Delta^a, p_+^{p,a}] \cap \mathcal{S} = \emptyset$. By the definition of \mathcal{P} , the corresponding trajectory Γ_p goes around a periodic orbit θ . By Proposition 2.10, if $p < p_+^{p,a}$ then θ must necessarily intersect both R_λ^+ and R_λ^- , while for $p = p_+^{p,a}$ the maximal periodic orbit θ_0 does not intersect R_λ^+ .

We claim that θ_0 is tangent to ℓ_+ at the point $P = (\frac{1+a}{p}, \lambda(N-1)) \in \ell_+ \cap \ell_2^+$ when $p = p_+^{p,a}$. If this was not the case, then $\Gamma = \Gamma_{p_+^{p,a}}$ would belong to the region R_λ^- for all $t \in I = [T, +\infty)$ for some $T > 0$. Let us consider the restriction of Γ to I , namely τ . Since τ is a part of a trajectory for the Laplacian operator in dimension \tilde{N}_+ , we may follow τ backward in time as a trajectory of the respective Laplacian-like dynamical system. However, the characterization of $p_+^{p,a}$ as the critical exponent there immediately contradicts the existence of τ . Indeed, at the critical exponent only periodic trajectories are admissible around M_0 , see for instance the proof of Proposition 6.1.

Thus, in both cases Γ_p for $p \in \mathcal{P}$ must cross the concavity line ℓ_+ infinitely many times. □

Remark 4.8 ($\Gamma_p = \Upsilon_p$). The critical exponent p_{a+}^* is the unique value of p for which Γ_p and Υ_p coincide, see Proposition 2.8.

Proof of Theorem 1.3. One establishes the conclusion of Theorem 1.3 by combining (4.3), Corollary 4.4, Propositions 4.5 and 4.6, together with Theorem 4.7.

4.3. Singular and exterior domain solutions

Here we show how the analysis of the regular trajectories performed in the previous sections almost completely determines the behavior of the other orbits of the dynamical system (2.9).

Let us start by considering singular solutions. When $p \leq p_+^{s,a}$ we saw in Proposition 4.6 that $p \in \mathcal{C}$. On the other hand, there is not a unique trajectory arriving at the stationary point A_0 as in Proposition 2.8. Indeed, for $p < p_+^{s,a}$, A_0 is a source and M_0 belongs to the fourth quadrant, see Proposition 2.7. The case $p = p_+^{s,a}$ is a bit more involved. The point $A_0 = M_0$ is not a hyperbolic point, and we complement its local study in what follows.

Proposition 4.9. *At $p = p_+^{s,a}$, there exist infinitely many unstable orbits issued from $M_0 = A_0$ below the line ℓ_1^+ . They move clockwise and blow up in finite forward time.*

Proof. At $p = p_+^{s,a}$ the eigenvalues of $A_0 = M_0$ are $(\tilde{N}_+ - 2)$ and 0 . In particular, A_0 is not hyperbolic and Proposition 2.2 no longer applies. The linear direction corresponding to $(\tilde{N}_+ - 2)$ lies on the X axis, while the one corresponding to 0 coincides with the line ℓ_1^+ . However, through the flow analysis in Proposition 2.5 (3) (see Figure 2) it is easy to conclude that M_0 has infinitely many repulsive directions between these two

lines. In this case, the orbits are issued from A_0 , with respective tangent lines between the X axis and the line ℓ_1^+ .

To see this let us first observe that ℓ_1^+ and ℓ_2^+ intersect at A_0 . Then note that $\dot{X} > 0$ in the region above ℓ_1^+ . On the other hand, an orbit coming out from A_0 needs to increase its Z values, so staying below ℓ_2^+ . If it started between ℓ_1^+ and ℓ_2^+ , then it should initially decrease its X values, which gives a contradiction. Hence the only way to come out from A_0 is below the line ℓ_1^+ .

We have already deduced in the proof of Proposition 4.6 that periodic orbits at $p_+^{s,a}$ are not admissible if $p_+^{s,a} < p_\Delta^a$ by Dulac's criterion (Proposition 2.10). However, this is true even if $p_+^{s,a} \geq p_\Delta^a$ by the flow direction, see Figure 2. Indeed, the region $\dot{X}, \dot{Z} < 0$ does not intersect $1Q$. By the same reason, the trajectories near $M_0 = A_0$ move clockwise, by intersecting both lines ℓ_1^+ and ℓ_2^+ exactly once.

To conclude we infer that the behavior of the flow on the lines ℓ_1^+ and ℓ_2^+ does not allow any orbit to reach $A_0 = M_0$ in forward time. Assume on the contrary that there exists a homoclinic orbit τ with $\omega(\tau) = \alpha(\tau) = A_0$. In this case τ creates a bounded invariant region D such that any orbit inside D is also homoclinic, by Poincaré-Bendixson theorem. Fix a point $Q_0 \in R_\lambda^- \cap \ell_2^+ \cap D$, and consider the unique trajectory τ_0 passing through this point at time $t=0$. By construction, τ_0 lies entirely in the region R_λ^- . However, the proof of Proposition 2.10(ii), applied to the region D_0 enclosed by the trajectory τ_0 , yields a contradiction with the fact that $p_+^{s,a} \neq p_+^{p,a}$. \square

It is interesting to observe that when $p \leq p_+^{s,a}$ these results give a simple proof of some Liouville theorems in [3], concerning radial solutions. The same holds for \mathcal{M}^- , as we shall see in Section 5.

Remark 4.10. For all $p > p_+^{s,a}$, as already mentioned in Section 1, there exists a singular trivial solution given by $u_p = C_p r^{-\alpha}$, C_p as in (3.8). This corresponds to the stationary trajectory $\Sigma_p \equiv M_0$. Moreover, any periodic orbit of the dynamical system (2.9) which intersects the concavity line ℓ_+ twice corresponds to a classical pseudo-blowing up solution for the problem (P_+) . Instead, in the case $p = p_+^{p,a}$, periodic orbits around the center configuration of M_0 lying entirely in the region R_λ^- (see Proposition 6.1) cannot correspond to C^2 solutions when $r > 0$, since they oscillate between the two functions $c_1 r^{-\alpha}$ and $c_2 r^{-\alpha}$ indefinitely for some $0 < c_1 < c_2$, without never changing convexity. We stress that these two types of solutions originated from periodic orbits do exist in the case of the Laplacian in dimensions N and \tilde{N}_+ , for the critical exponents p_Δ^a and $p_+^{p,a}$ respectively, see Theorem 6.1(iii) in [19].

Lemma 4.11. *If $p > p_{a+}^*$ then Υ_p (see Proposition 2.8) blows up in finite backward time. In particular, Υ_p does not correspond to a singular solution for any $p > p_{a+}^*$.*

Proof. If $p \in (p_{a+}^*, p_+^{p,a}]$ we have $p \in \mathcal{P}$, and there exists a maximal periodic orbit θ_p around M_0 such that $\omega(\Gamma_p) = \theta_p$. If $\alpha(\Upsilon_p) = \theta_p$, then Γ_p and Υ_p would cross somewhere. Indeed, this comes from the fact that the stable linear tangent direction at A_0 is above the line ℓ_2^+ (see Figure 1), and the vector field on ℓ_2^+ points down for $X > \alpha$, by Proposition 2.5 (3). Obviously, crossings are not admissible by uniqueness of the ODE problem. On the other side, for $p > p_+^{p,a}$, M_0 is a sink and periodic orbits are not

allowed by [Proposition 2.10](#). Thus, in both cases, using [Proposition 2.11](#), we get that Υ_p blows up backward in finite time.

Proof of Theorem 1.7. (i) First we recall that for $p \in (1, p_+^{s,a}]$ there are no periodic orbits, by [Proposition 2.10](#) and the proof of [Proposition 4.9](#). Moreover, we know by [Corollary 4.4](#) and [Theorem 4.7](#) that the regular trajectory Γ_p blows up forward in finite time. Hence, Γ_p together with the line $L = \{(X, Z) : X = \tilde{N}_+ - 2\}$ and the X, Z axes, create a bounded region D from which an orbit of (2.9) may only leave through L .

Thus, if $p < p_+^{s,a}$, any trajectory issued from A_0 (which is a source by [Proposition 2.7](#)) crosses the line L and then blows up in finite time. If $p = p_+^{s,a}$ the same holds, by [Proposition 4.9](#). In both cases there are infinitely many such trajectories which correspond to singular solutions in an interval $(0, R)$, $R > 0$, see [Proposition 3.1](#). They are $(\tilde{N}_+ - 2)$ -blowing up, cf. (3.7) with the ω -limit exchanged by α -limit. Therefore there cannot be singular solutions in $\mathbb{R}^N \setminus \{0\}$ for this range of p .

(ii)-(iii) For $p \in (p_+^{s,a}, p_{a+}^*)$, as in (i), the trajectory Γ_p , the line L and the X, Z axes determine a bounded region D from which an orbit may only leave through L . Recall that for these values of p , A_0 is a saddle point and M_0 is a source, see [Proposition 2.7](#). Thus, the unique orbit Υ_p arriving at A_0 (see [Proposition 2.8](#)) can either converge to M_0 or to a periodic orbit around M_0 , backward in time. If $p \leq p_\Delta^a$ there are no periodic orbits ([Proposition 2.10](#)), so Υ_p corresponds to α -blowing up solution of (P_+) ; in particular this is the case for each $p \in (p_+^{s,a}, p_\Delta^a]$ if $p_+^{s,a} \leq p_\Delta^a$.

If in turn $p \in (p_\Delta^a, p_{a+}^*)$ there could be periodic orbits around M_0 , so that Υ_p corresponds to either a pseudo-blowing up or a α -blowing up solution of (P_+) .

All the other orbits coming out from M_0 or from a periodic orbit θ around M_0 (whenever such θ exists) must necessarily leave D by crossing the line L in forward time, and therefore blow up in finite time. This gives infinitely many singular solutions of (P_+) ; they are either α -blowing up or pseudo-blowing up in intervals $(0, R)$, $R > 0$.

If θ exists, we have in addition infinitely many orbits τ issued from M_0 and converging to a minimal periodic orbit (which is θ if the system has only one limit cycle). Each τ crosses infinitely many times the line ℓ_+ when $t \rightarrow +\infty$, and so corresponds to a α -blowing up solution of (P_+) with pseudo-slow decay at $+\infty$ as in (3.9). On the other hand, a periodic orbit itself in this range of p 's crosses ℓ_+ twice, so corresponds to a pseudo-blowing up solution to (P_+) , see [Remark 4.10](#); they are pseudo-slow decaying and change concavity infinitely many times both as $r \rightarrow 0$ and $r \rightarrow +\infty$.

(iv) When $p = p_{a+}^*$, the regular trajectory $\Gamma_{p_{a+}^*}$ together with the X and Z axes delimit an invariant set D containing M_0 . Since M_0 is a source, we have already seen that there exists a periodic orbit around M_0 ; say θ is the minimal one. We then infer that there exist infinitely many periodic orbits in the region $D \setminus \text{int}(\theta)$, at least in a neighborhood of ∂D , see [Figure 7](#). Indeed, the existence of a maximal periodic orbit θ_0 inside D would create a bounded region $D \setminus \text{int}(\theta_0)$ in which the orbits issued from θ_0 could not go anywhere, thus violating Poincaré-Bendixson theorem.

(v) When $p \in (p_{a+}^*, p_+^{p,a})$, we have that M_0 is a source and there exists a minimal periodic orbit θ around M_0 . In this case, θ crosses the line ℓ_+ twice since $p < p_+^{p,a}$, see [Proposition 2.10](#). Thus, all trajectories issued from M_0 converge to θ in forward time. These and the periodic orbits give us singular solutions as in the last part of the proof

of (iii). Finally, note that no singular solutions converge to A_0 by Lemma 4.11, so the assertion holds.

(vi) By Propositions 2.7, 2.10, and 6.1 we have that for $p = p_+^{p,a}$ the stationary point M_0 is a center while for $p > p_+^{p,a}$ M_0 is a sink without periodic orbits. These and the fact that A_0 is a saddle point, whose unstable manifold is the X axis, imply that no singular nontrivial solutions are admissible. \square

Finally we consider the case of exterior domain solutions, proving Theorem 1.9 for \mathcal{M}^+ . The proof for \mathcal{M}^- turns out to be the same.

In Section 3.3 we have observed that a solution u of (3.10) necessarily satisfies (3.11). Hence the corresponding trajectory Ξ_p blows up backward in finite time, see Proposition 3.11. Thus, to prove Theorem 1.9 it is enough to show that for $p \in (1, p_{a+}^*]$ there are no orbits of the dynamical system (2.9) defined in $(T, +\infty)$ for some $T > 0$ with this kind of blow-up behavior.

Proof of Theorem 1.9. By the definition and properties of the critical exponent in Sections 4.1, 4.2, we know that $p \in \mathcal{C}$ for $p \in (1, p_{a+}^*)$, while $p_{a+}^* \in \mathcal{F}$. In the first case the regular trajectory Γ_p together with the X and Z axes and the line $L = \{(X, Z) : X = \tilde{N}_+ - 2\}$ bound a region D from which any trajectory can only escape in forward time through L . In the second case Γ_p and the X and Z axes enclose a bounded invariant region D . In both cases the closure of D contains the points M_0 (for $p \geq p_+^{s,a}$) and A_0 .

By contradiction assume that a radial solution of (1.9) exists. Then, by Proposition 3.11 the corresponding trajectory Ξ_p is defined in an interval $(t_\delta, +\infty]$ for some $t_\delta > -\infty$, and blows at t_δ satisfying (3.12). Since it does not blow up in forward time, by Proposition 2.11 (see (2.25)) and Poincaré-Bendixson theorem the ω -limit $\omega(\Xi_p)$ is either M_0 (if $p > p_+^{s,a}$), or A_0 , or a periodic orbit around M_0 . In any case Ξ_p should cross Γ_p which is not possible.

Remark 4.12. It is proved in [16], when $a=0$, that for every $p \in (p_{a+}^*, +\infty)$ both a fast decaying and infinitely many slow or pseudo-slow decaying solutions of (1.9) exist. In terms of our quadratic system (2.9) this could be proved using Lemma 4.11 for the fast decaying solutions, or studying the trajectories arriving at M_0 or at a periodic orbit for the slow or pseudo-slow decaying solutions. However, since the proof of [16] easily extends to the case $a \neq 0$, we prefer to omit the details.

Note that with the analysis of the trajectories blowing up backward in finite time one can only get the existence of a solution u satisfying (3.11) at some radius $\mu > 0$. Then the solution should be continued (in 3Q) to reach a positive radius $\rho_0 > 1$ where $u(\rho_0) = 0$, so to verify the Dirichlet problem in the exterior of a ball. This is possible by using a shooting argument from μ , as done for instance in [16, proof of Theorem 6.1].

5. The \mathcal{M}^- case

In this section we analyze the complementary case for the operator \mathcal{M}^- . Recall its respective dimension-like parameter from (1.5) satisfying $\tilde{N}_- \geq N$. The main difference with the case of \mathcal{M}^+ is the reverse ordering of the exponents p_Δ^a and $p_-^{p,a}$ from (1.6),

with $p^{s,a} \leq p^{p,a} \leq p_\Delta^a$. Also, if $\lambda < \Lambda$ then the stationary point M_0 is a sink in the interval $(p_-^{p,a}, p_\Delta^a)$, see [Proposition 2.7](#) (4).

We start by pointing out that all properties stated for \mathcal{M}^+ in section 4.1 also hold for \mathcal{M}^- . In particular, one gets that the set \mathcal{F} possesses at most one point, which splits the interval $(1, +\infty)$ into two components \mathcal{C} and $\mathcal{P} \cup \mathcal{S}$ accordingly to [Corollary 4.4](#). Moreover, each of these components is nonempty, since one verifies, as in [Propositions 4.5](#) and 4.6, the following result.

Proposition 5.1. *If $p > p_\Delta^a$ then $p \in \mathcal{S}$, and for $p < p_-^{p,a}$ it holds that $p \in \mathcal{C}$.*

This allows to define, as in Section 4, the critical exponent p_{a-}^* as follows

$$p_{a-}^* = \sup \mathcal{C}.$$

Then, by [Proposition 5.1](#), $p_{a-}^* \in [p_-^{p,a}, p_\Delta^a]$ and, as for \mathcal{M}^+ , we call it the critical exponent for \mathcal{M}^- . Next we show that $p_{a-}^* \in \mathcal{F}$ and p_{a-}^* is in the interior of the previous interval.

Theorem 5.2. *The critical exponent p_{a-}^* belongs to \mathcal{F} . Thus it is the only exponent in the equation (P_-) for which there exists a unique, up to scaling, fast decaying solution.*

Moreover, if $\lambda < \Lambda$, then [\(1.8\)](#) holds and there exists $\varepsilon > 0$ such that $(p_\Delta^a - \varepsilon, +\infty) \subset \mathcal{S}$.

Proof. Obviously $p_{a-}^* \notin \mathcal{C}$ because \mathcal{C} is open, see [Remark 3.10](#). Moreover, p_{a-}^* cannot belong to \mathcal{P} ; otherwise $\Gamma_{p_{a-}^*}$ should cross the line ℓ_1^+ by [Lemma 4.1\(ii\)](#), while Γ_p for $p \in \mathcal{C}$ never does it.

Finally we show that $p_{a-}^* \notin \mathcal{S}$. In $\Gamma_{p_{a-}^*}$ deed, if this was the case then $p_{a-}^* > p_-^{p,a}$ because M_0 is a center at $p_-^{p,a}$, see [Proposition 6.1](#). Hence $M_0 = M_0(p)$ is a sink for every p in a neighborhood $I_\varepsilon = (p_{a-}^* - \varepsilon, p_{a-}^* + \varepsilon)$ for some $\varepsilon > 0$, by [Proposition 2.7](#) (4). Then there exists a maximal ball B_{η_p} centered at $M_0(p)$ with the property that any trajectory τ_p entering in B_{η_p} satisfies $\omega(\tau_p) = M_0(p)$, see [\[15\]](#). Since we are assuming that $p_{a-}^* \in \mathcal{S}$ then $\omega(\Gamma_{p_{a-}^*}) = M_0(p_{a-}^*)$. By the continuity of the dynamical system with respect to the parameter p , also $\omega(\Gamma_p) = M_0(p)$ for $p \in I_\varepsilon$ (up to diminishing ε). But this contradicts the definition of p_{a-}^* , since Γ_q blows up in finite time when $q \in \mathcal{C}$.

Hence, $p_{a-}^* \in \mathcal{F}$. The proof that p_{a-}^* cannot be $p_-^{p,a}$ nor p_Δ^a is the same as the one for \mathcal{M}^+ , see [Theorem 4.7](#). It relies on [Proposition 2.10](#) which states, in particular, that there are no periodic orbits of [\(2.11\)](#) for p_Δ^a . This also proves that the regular trajectory $\Gamma_{p_\Delta^a}$ converges to $M_0 = M_0(p_\Delta^a)$, so that $p_\Delta^a \in \mathcal{S}$. Next, a continuity argument as in the first part of this proof shows that $p \in (p_\Delta^a - \varepsilon, p_\Delta^a)$ also belongs to \mathcal{S} for sufficiently small $\varepsilon > 0$. Consequently, for such p 's it does not exist a periodic orbit around M_0 , and the proof is complete. □

Proof of Theorem 1.4. All previous results obtained for \mathcal{M}^- prove the theorem. In particular, the statements (iii)–(iv) follow from [Theorem 5.2](#). □

We finish the section with the proof of [Theorem 1.8](#) about singular solutions.

Proof of Theorem 1.8. It is enough to prove (iii) and (v), since the proof of the other items are the same as for [Theorem 1.7](#). Let us analyze the whole interval $p \in (p_-^{p,a}, p_\Delta^a]$.

Recall that M_0 is a sink whenever $p > p_-^{p,a}$. Therefore, there is no trajectory coming out from M_0 in the range $p \in (p_-^{p,a}, p_\Delta^a]$. In particular, $\alpha(\Upsilon_p)$ is never M_0 in this range of p .

For $p \in (p_-^{p,a}, p_{a-}^*)$ we have $p \in \mathcal{C}$. In this case the regular trajectory Γ_p and the line $L = \{(X, Z) : X = \tilde{N}_- - 2\}$, together with the X and Z axes, create a bounded region from which any orbit may only leave forward in time through L ; recall that the flow is going out on L , see [Figure 1](#). Therefore, Poincaré-Bendixson theorem implies the existence of a periodic orbit θ_p around M_0 such that $\alpha(\Upsilon_p) = \theta_p$. This immediately determines four types of nontrivial positive pseudo-blowing up solutions of (1.1)–(1.2) (in the case of \mathcal{M}^-):

- (1) a fast decaying solution corresponding to the trajectory Υ_p ;
- (2) solutions with slow decay, whose corresponding orbits lie inside a minimal periodic orbit θ_0 around M_0 ; here θ_0 crosses ℓ_- twice due to [Proposition 2.10](#);
- (3) solutions of the Dirichlet problem in $B_R \setminus \{0\}$, such that the corresponding orbits are issued from θ_p and blow up in finite forward time;
- (4) pseudo-slow decaying solutions, which correspond to the periodic orbits.

All of these singular solutions change concavity infinitely many times in a neighborhood of $r=0$. Further, there are infinitely many solutions of types (2) and (3); see [Figure 8](#). Thus, (iii) holds.

To prove (v), let us recall that at p_Δ^a no periodic orbits are admissible by [Proposition 2.10](#). Also, by [Theorem 5.2](#) there exists $\varepsilon > 0$ such that $p \in \mathcal{S}$ for all $p \in (p_\Delta - \varepsilon, +\infty)$. Now, arguing as in [Lemma 4.11](#) one sees that Υ_p blows up in finite backward time for $p > p_{a-}^*$, so (v) is proved.

Remark 5.3. In the case of \mathcal{M}^- the existence of singular solutions in the range $(p_{a-}^*, p_\Delta^a - \varepsilon)$ is not guaranteed, though solutions as in the cases (2)–(4) in the proof above are admissible.


Concerning the exterior domain solutions, we have already observed in [Section 4.3](#) that the proof of [Theorem 1.9](#) is the same for both operators \mathcal{M}^\pm .

Acknowledgments

Part of this work was done when the first and second authors were visiting Sapienza University of Rome. They would like to thank the warm hospitality and care provided by the people in the Math Department.

Funding

L. Maia was supported by FAPDF, CAPES, and CNPq grant 308378/2017-2. G. Nornberg was supported by FAPESP grants 2018/04000-9 and 2019/031019-9, São Paulo Research Foundation. F. Pacella was supported by INDAM-GNAMPA.

ORCIDLiliane Maia  <http://orcid.org/0000-0002-6163-1899>Gabrielle Nornberg  <http://orcid.org/0000-0002-0395-0250>Filomena Pacella  <http://orcid.org/0000-0002-1269-6530>**References**

- [1] Caffarelli, L. A., Cabré, X. (1995). *Fully Nonlinear Elliptic Equations, Volume 43 of American Mathematical Society Colloquium Publications*. Providence, RI: American Mathematical Society.
- [2] Quaas, A., Sirakov, B. (2006). Existence results for nonproper elliptic equations involving the Pucci operator. *Comm. Partial Differ. Equ.* 31(7):987–1003.
- [3] Cutri, A., Leoni, F. (2000). On the Liouville property for fully nonlinear equations. *Ann. L'Institut Henri Poincaré, Sect. (C)*. 17(2):219–245.
- [4] Armstrong, S. N., Sirakov, B. (2011). Sharp Liouville results for fully nonlinear equations with power-growth nonlinearities. *Ann. Sc. Norm. Super. Pisa Cl. Sci.* 10(3):711–728.
- [5] Armstrong, S. N., Sirakov, B., Smart, C. K. (2011). Fundamental solutions of homogeneous fully nonlinear elliptic equations. *Commun. Pure Appl. Math.* 64(6):737–777.
- [6] Felmer, P. L., Quaas, A. (2003). On critical exponents for the Pucci's extremal operators. *Ann. Inst. H. Poincaré Anal. Non Linéaire* 20(5):843–865.
- [7] Da Lio, F., Sirakov, B. (2007). Symmetry results for viscosity solutions of fully nonlinear uniformly elliptic equations. *J. Eur. Math. Soc. (JEMS)* 9(2):317–330.
- [8] Caffarelli, L. A., Gidas, B., Spruck, J. (1989). Asymptotic symmetry and local behavior of semilinear elliptic equations with critical Sobolev growth. *Commun. Pure Appl. Math.* 42(3):271–297.
- [9] Gladiali, F., Grossi, M., Neves, S. L. N. (2013). Nonradial solutions for the Hénon equation in \mathbb{R}^N . *Adv. Math.* 249:1–36.
- [10] Birindelli, I., Galise, G., Leoni, F., Pacella, F. (2018). Concentration and energy invariance for a class of fully nonlinear elliptic equations. *Calc. Var.* 57(6):1–22.
- [11] Bidaut-Veron, M. F., Giacomini, H. (2010). A new dynamical approach of Emden-Fowler equations and systems. *Adv. Differ. Equ.* 15(11–12):1033–1082.
- [12] Clément, P., de Figueiredo, D. G., Mitidieri, E. (1996). Quasilinear elliptic equations with critical exponents. *TMNA*. 7(1):133–170.
- [13] Felmer, P. L., Quaas, A. (2006). Critical exponents for uniformly elliptic extremal operators. *Indiana Univ. Math. J.* 55(2):593–629.
- [14] Chicone, C., Tian, J. H. (1982). On general properties of quadratic systems. *Am. Math. Monthly* 89(3):167–178.
- [15] Hale, J. K., Koçak, H. (1991). *Dynamics and Bifurcations, Volume 3 of Texts in Applied Mathematics*. New York: Springer-Verlag.
- [16] Galise, G., Iacopetti, A., Leoni, F. (2019). Liouville-type results in exterior domains for radial solutions of fully nonlinear equations. *J. Differ. Equ.* 266(2–3):1675–1697. To appear in
- [17] González-Melendez, O., Quaas, A. (2017). On critical exponents for Lane-Emden-Fowler-type equations with a singular extremal operator. *Ann. Matematica*. 196(2):599–615.
- [18] Galise, G., Leoni, F., Pacella, F. (2017). Existence results for fully nonlinear equations in radial domains. *Comm. Partial Differ. Equ.* 42(5):757–779.
- [19] Bidaut-Véron, M.-F. (1989). Local and global behavior of solutions of quasilinear equations of Emden-Fowler type. *Arch. Ration. Mech. Anal.* 107(4):293–324.

Appendix

A.1. Local study

In this section we detail the proof of Proposition 2.7. The linearization for $\mathcal{M}^+, \tilde{N} = \tilde{N}_+$, is

$$L(X, Z) = \begin{pmatrix} \partial_X f^+ & \partial_Z f^+ \\ \partial_X g^+ & \partial_Z g^+ \end{pmatrix} = \begin{pmatrix} 2X - (N - 2) + \frac{Z}{\lambda} & \frac{X}{\lambda} \\ -pZ & N + a - pX - \frac{2Z}{\lambda} \end{pmatrix} \text{ in } R_\lambda^+,$$

$$L(X, Z) = \begin{pmatrix} \partial_X f^- & \partial_Z f^- \\ \partial_X g^- & \partial_Z g^- \end{pmatrix} = \begin{pmatrix} 2X - (\tilde{N} - 2) + \frac{Z}{\Lambda} & \frac{X}{\Lambda} \\ -pZ & \tilde{N} + a - pX - \frac{2Z}{\Lambda} \end{pmatrix} \text{ in } R_\lambda^-.$$

For instance, at $N_0 = (0, \lambda(N + a))$ and $A_0 = (\tilde{N} - 2, 0)$ one has

$$L(N_0) = \begin{pmatrix} 2 + a & 0 \\ -p\lambda(N + a) & -N - a \end{pmatrix}, \quad L(A_0) = \begin{pmatrix} \tilde{N} - 2 & \frac{\tilde{N} - 2}{\Lambda} \\ 0 & \tilde{N} + a - p(\tilde{N} - 2) \end{pmatrix}.$$

The eigenvalues for N_0 are 2 and $-N - a$, while for A_0 are $\sigma_1 = \tilde{N}_\pm - 2$ and $\sigma_2 = \tilde{N}_\pm + a - p(\tilde{N} - 2)$. Recall that $M_0 = (X_0, Z_0)$, where $X_0 = \alpha = \frac{a+1}{p-1}$ and $Z_0 = \Lambda(\tilde{N} - p\alpha + a) = \Lambda(\tilde{N} - 2 - \alpha)$,

$$L(M_0) = \begin{pmatrix} \alpha & \frac{\alpha}{\Lambda} \\ -p\Lambda(\tilde{N} - p\alpha + a) & -(\tilde{N} - p\alpha + a) \end{pmatrix}.$$

In order to analyze the eigenvalues of $L(M_0)$ one needs to look at the roots of the equation

$$\sigma^2 + \sigma \left(\frac{Z_0}{\Lambda} - X_0 \right) + X_0(p - 1) \frac{Z_0}{\Lambda} = 0.$$

They are given by $2\sigma_\pm = X_0 - \frac{Z_0}{\Lambda} \pm \sqrt{\Delta}$, where $\Delta = \left(\frac{Z_0}{\Lambda} - X_0\right)^2 - 4(2 + a) \frac{Z_0}{\Lambda}$.

Note that $X_0 = \frac{Z_0}{\Lambda}$ is equivalent to $\alpha = \frac{\tilde{N}-2}{2}$, i.e. $p = p_+^{p,a}$. In this case $\text{Re}(\sigma_\pm) = 0$ and the roots are purely imaginary. Moreover, $X_0 > \frac{Z_0}{\Lambda} \iff p < p_+^{p,a}$, and $X_0 < \frac{Z_0}{\Lambda} \iff p > p_+^{p,a}$.

If $\text{Im}(\sigma_\pm) \neq 0$, this already determines the sign of $\text{Re}(\sigma_\pm)$. Assume then $\text{Im}(\sigma_\pm) = 0$ i.e. $\sigma_\pm \in \mathbb{R}$. Observe that $\Delta < \left(X_0 - \frac{Z_0}{\Lambda}\right)^2$ as far as M_0 stays in 1Q. This yields $\sigma_- > 0$ for $p_+^{s,a} < p < p_+^{p,a}$ (so $\sigma_\pm > 0$ and M_0 is a source); while $\sigma_+ < 0$ if $p > p_+^{p,a}$ (so $\sigma_\pm < 0$ and M_0 is a sink).

It is possible to prove that M_0 is a saddle point in the fourth quadrant when $1 < p < p_+^{s,a}$. However, this would correspond to solutions of the absorption problem $\mathcal{M}^+ u - u^p = 0$. See [11] in the case of the Laplacian operator.

A.2. Energy analysis

Let us consider the energy functional E of the operator \mathcal{M}^+ in the region R_λ^- , which is a slight variation of the energy of the Laplacian operator in dimension \tilde{N}_+ treated in [11],

$$E(t, X, Z) = e^{t(\tilde{N}_+ - 2 - 2\alpha)} X(XZ)^\alpha \left\{ \frac{X}{2} + \frac{Z}{\Lambda(p+1)} - \frac{\tilde{N}_+}{p+1} \right\} \text{ in } R_\lambda^- \cup \ell_+$$

understood as natural extension up to ℓ_+ . In terms of u , the energy functional E for \mathcal{M}^+ reads as

$$E(r) = E(r, u) = r^{\tilde{N}_+} \left(\frac{(u')^2}{2} + \frac{1}{\Lambda} \frac{r^a u^{p+1}}{p+1} \right) + \frac{\tilde{N}_+}{p+1} uu' r^{\tilde{N}_+-1} \quad \text{if } u'' \geq 0.$$

Of course these two expressions are equivalent after the transformation (2.4). Moreover,

$$E'(r) = r^{\tilde{N}_++a-1} (u')^2 \left(\frac{\tilde{N}_++a}{p+1} - \frac{\tilde{N}_+-2}{2} \right) \quad \text{if } u'' \geq 0,$$

and so the following monotonicity holds

$$\dot{E} < 0 \text{ if } p > p_+^{p,a}, \quad \dot{E} = 0 \text{ if } p = p_+^{p,a}, \quad \dot{E} > 0 \text{ if } p < p_+^{p,a} \quad \text{in } R_\lambda^- \cup \ell_+. \quad (\text{A.1})$$

Now we investigate the precise behavior of the trajectories close to M_0 at $p = p_\pm^p$. Here, $\lambda \leq \Lambda$ and the result gives an alternative proof in the case of the Laplacian operator $\lambda = \Lambda = 1$ in [19].

Proposition 6.1. M_0 is a center when $p = p_\pm^{p,a}$.

Proof. We present the proof for \mathcal{M}^+ ; for \mathcal{M}^- it is the same in light of Section 5. Let $\tau = (X, Z)$ be an orbit contained in $R_\lambda^- \cup \ell_+$. Let us show that τ is periodic. To simplify notation let $a = 0, \tilde{N} = \tilde{N}_+$. The energy of τ on the line ℓ_2^+ is given by

$$E|_{\ell_2^+ \cap R_\lambda^-} = E|_{\ell_2^+} (X) = -\Lambda^\alpha \tilde{N} X^{\alpha+2} (\tilde{N} - pX)^\alpha \quad \text{at } p = p_+^p = p_+^{p,0}.$$

Since the energy is a constant function of t when $p = p_+^p$ with $\frac{\alpha+2}{\alpha} = p$, then

$$(\tilde{N} - pX)X^p \equiv c > 0 \quad \text{on } \ell_2^+. \quad (\text{A.2})$$

Now we may translate the information from (A.2) in terms of the function h defined as

$$h(X) = (\tilde{N} - pX)X^p, \text{ for } X \in [1/p, \tilde{N}/p] \text{ where } p = p_+^p,$$

for which (A.2) represents its level curves. The domain $[1/p, \tilde{N}/p]$ entails the behavior of h in the respective interval delimited by ℓ_2^+ on R_λ^- , up to the boundary.

Let us analyze the function h ; it is positive at $1/p$, and equals to zero at \tilde{N}/p . Since

$$h'(X) = pX^p \left(\frac{\tilde{N}}{X} - 1 - p \right), \text{ with } \frac{\tilde{N}}{1+p} = \frac{\tilde{N}-2}{2} = \alpha, \quad p = p_+^p,$$

then h is increasing when $X < \alpha$, decreasing for $X > \alpha$, and it assumes the positive maximum value $h(\alpha) = \alpha^{p+1} := c_\infty$ at $X = \alpha$. Moreover, note that $h(1/p) = (\tilde{N}-1)(1/p)^p := c_1$.

Here h is a polynomial function which prescribes the value of the energy on $\ell_2^+ \cap R_\lambda^-$, namely $h = -E^{1/\alpha} \equiv c$. For any $k \in \mathbb{N}$ with $c_k \in [c_1, c_\infty)$, the line $h \equiv c_k$ intersects the graph of h at exactly two points X_1^k, X_2^k such that $X_1^k < \alpha < X_2^k$. Also, they satisfy

$$c_k = h(X_1^k) = h(X_2^k) \rightarrow h(\alpha) = c_\infty \text{ when } X_i^k \rightarrow \alpha \text{ as } k \rightarrow +\infty, \quad i = 1, 2. \quad (\text{A.3})$$

Furthermore, the line $h = c_\infty$ intersects the graph of h only once at the point $X = \alpha$.

In our phase plane context, this means that any trajectory τ contained in $R_\lambda^- \cup \ell_+$ bisects the line ℓ_2^+ at exactly two points $P_1 = (X_1, Z_1), P_2 = (X_2, Z_2)$, with $X_1 < \alpha, X_2 > \alpha$. By Proposition 2.5 (3) the flow moves horizontally on ℓ_2^+ , namely to the right for $X < \alpha$, and to the left when $X > \alpha$.

Observe that ℓ_2^+ is a transversal section to the flow, on which any trajectory approaching M_0 must pass across, either in the past or in the future. Hence, the trajectory τ has to be closed, by moving clockwise. Since this dynamics is realized for any trajectory contained on $R_\lambda^- \cup \ell_+$, and (A.3) holds, in particular any trajectory close to M_0 is periodic, so M_0 is a center. \square