Bubbling nodal solutions for a large perturbation of the Moser-Trudinger equation on planar domains

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March 5, 2019

Abstract

In this work we study the existence of nodal solutions for the problem

 $-\Delta u = \lambda u e^{u^2 + |u|^p}$ in $\Omega, u = 0$ on $\partial \Omega$,

where $\Omega \subseteq \mathbb{R}^2$ is a bounded smooth domain and $p \to 1^+$.

If Ω is ball, it is known that the case p = 1 defines a critical threshold between the existence and the non-existence of radially symmetric sign-changing solutions. In this work we construct a blowing-up family of nodal solutions to such problem as $p \to 1^+$, when Ω is an arbitrary domain and λ is small enough. As far as we know, this is the first construction of sign-changing solutions for a Moser-Trudinger critical equation on a non-symmetric domain.

1 Introduction

Let us consider the equation

$$\Delta u + \lambda u e^{u^2 + a|u|^p} = 0 \text{ in } \Omega, \ u = 0 \text{ on } \partial\Omega, \tag{1}$$

where Ω is a bounded smooth domain in \mathbb{R}^2 , λ is a positive parameter and the nonlinear term $h(u) := ue^{a|u|^p}$, with $a \in \mathbb{R}$ and $p \in [0, 2)$, is a lower-order perturbation of e^{u^2} according to the definition given by Adimurthi in [2].

The nonlinearity $f(u) = h(u)e^{u^2}$ is critical from the view point of the Trudinger imbedding. Indeed, in view of the Moser-Trudinger inequality (see [25, 29, 24])

$$\sup\left\{\int_{\Omega} e^{u^2} dx : u \in H^1_0(\Omega), \ \|u\|^2_{H^1_0(\Omega)} \le 4\pi\right\} < +\infty,$$
(2)

the functional

$$J_{\lambda}(u) := \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx - \lambda \int_{\Omega} F(u) dx, \ u \in H^1_0(\Omega),$$
(3)

where $F(t) = \int_{0}^{t} f(s)ds$, is well defined and its critical points are solutions to problem (1). Adimurthi in [2] proved that J_{λ} satisfies the Palais-Smale condition in the infinite energy range $(-\infty, 2\pi)$ but, as observed by Adimurthi and Prashant in [5], the critical nature of f(u) reflects in the failure of the Palais-Smale condition at the sequence of energy levels $2\pi k$ with $k \in \mathbb{N}$ (see also [7]).

In [2] Adimurthi proved the existence of a critical point of J_{λ} if the perturbation his large, i.e. $a \geq 0$, and if $0 < \lambda < \lambda_1(\Omega)$, where $\lambda_1(\Omega)$ is the first eigenvalue of $-\Delta$ with Dirichlet boundary condition ((see also [1])). Such a critical point is a positive solution to problem (1). Successively, Adimurthi and Prashant in [6] showed that the condition $a \geq 0$ is necessary to get a positive solution to (1). Indeed, they proved that if the perturbation h is small, i.e. a < 0, then there are no positive solutions to problem (1) when the domain Ω is a ball provided λ is small. The case a = 0 in a general domain Ω has been studied by Del Pino, Musso and Ruf [14] using a perturbative approach. Indeed they find multiplicity of positive solutions which blow-up in one or more points of Ω (depending on the geometry) as $\lambda \to 0$. We point out that a general qualitative analysis of blowing-up families of positive solutions to problem (1) has been obtained by Druet in [15] (see also [3, 17, 16]).

As far as it concerns the existence of sign-changing solutions, Adimurthi and Yadava in [8] proved that problem (1) has a nodal solution when λ is small if there is the further restriction p > 1 on the growth of the large perturbation h (i.e. a > 0). Actually, this condition turns out to be optimal for the existence of nodal radial solutions in a ball. Indeed Adimurthi and Yadava in [9] proved that if a > 0 and Ω is a ball, problem (1) does not have any radial sign-changing solution when λ is small and $p \in [0, 1]$. If one drops the radial requirement, Adimurthi and Yadava in [8] proved the existence of infinitely many sign-changing solutions in a ball whatever $\lambda > 0$ is. We point out that, in the case a = 0, the approach of Del Pino, Musso and Ruf [14] allows to find sign-changing solutions which blow-up positively and negatively at least at two different points in any domain Ω as $\lambda \to 0$ (even if this is not explicitly said in their work).

According to the previous discussion, it turns out that when a > 0 the case p = 1 defines a critical threshold for the existence of radial sign-changing solutions in the ball. Indeed, when $\Omega = B(0, 1)$, (1) has radially symmetric sign-changing solutions which blow-up as $p \to 1^+$. The precise behavior of such solutions was studied by Grossi and Naimen in [19]. Therefore, when a > 0, it is natural to ask whether it is possible to find sign-changing solutions to problem (1) on an arbitrary planar domain Ω which blow-up at one point in Ω as $p \to 1^+$.

In this paper we give a positive answer. More precisely, let us consider the problem

$$\begin{cases} -\Delta u = \lambda u e^{u^2 + |u|^{1+\varepsilon}} & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$
(4)

where ε is a positive small parameter. Set

$$f_{\varepsilon}(t) = te^{t^2 + |t|^{1+\varepsilon}}.$$
(5)

For a given $0 < \lambda < \lambda_1(\Omega)$, let u_0 be a positive solution of the problem

$$\begin{cases} -\Delta u_0 = \lambda f_0(u_0) & \text{in } \Omega, \\ u_0 > 0 & \text{in } \Omega, \\ u_0 = 0 & \text{on } \partial\Omega, \end{cases}$$
(6)

whose existence has been established by Adimurthi in [2]. We make the following assumptions:

(A1) u_0 is non-degenerate, i.e. there is no non-trivial solution $\varphi \in H^1_0(\Omega)$ of the equation

$$-\Delta \varphi = \lambda f_0'(u_0)\varphi \text{ in } \Omega, \ \varphi = 0 \text{ on } \partial\Omega.$$
(7)

(A2) u_0 has a C^1 -stable critical point $\xi_0 \in \Omega$ such that $u_0(\xi_0) > \frac{1}{2}$.

Then, we will show that (4) admits a family of sign-changing solutions which blow-up at ξ_0 with residual mass $-u_0$ as $\varepsilon \to 0$, namely:

Theorem 1.1 For $0 < \lambda < \lambda_1(\Omega)$, let u_0 be a solution of (6) such that (A1) and (A2) are satisfied. Let also ξ_0 be as in (A2). Then there exist $\varepsilon_0 > 0$ and a family $(u_{\varepsilon})_{0 < \varepsilon < \varepsilon_0}$ of sign-changing solutions to (4) such that:

- $\max_{\overline{B(\xi_0,r)}} u_{\varepsilon} \to +\infty \text{ as } \varepsilon \to 0, \text{ for any } 0 < r < d(\xi_0,\partial\Omega).$
- $u_{\varepsilon} \to -u_0$ weakly in $H_0^1(\Omega)$ and in $C^1(\overline{\Omega} \setminus \{\xi_0\})$.

Let us make some comments about assumptions (A1) and (A2).

- **Remark 1.2** The solution u_0 to problem (6) turns out to be non-degenerate when Ω is the ball as proved by Adimurthi, Karthik and Giacomoni in [4]. In a work in progress, Grossi and Naimen are going to prove that the solution is also non-degenerate when Ω is convex and symmetric (see [20]). Actually, we believe that the non-degeneracy condition holds true for most domains Ω and positive parameters λ . Indeed, one could use similar arguments to those used by Micheletti and Pistoia in [23] for a class of singularly perturbed equations.
 - We remind that ξ_0 is a C^1 -stable critical point of u_0 if the Brouwer degree $\deg(\nabla u_0, B(\xi_0, r), 0) \neq 0$. In particular, any strict local maximum point of u_0 is C^1 -stable. We point out that by Adimurthi and Druet [3] we can deduce that assumption (A2) holds true when the parameter λ is small enough.

• We strongly believe that the condition $u_0(\xi_0) > \frac{1}{2}$ is not purely technical, but it is necessary to build a solution which blows-up at ξ_0 . Indeed, we conjecture that, if $u_0(\xi_0) \leq \frac{1}{2}$, there does not exist any sign-changing solution which blows-up at ξ_0 with non-trivial residual mass u_0 as $\varepsilon \to 0$. We point out that, in a different setting, a similar condition was proved by Mancini and Thizy [22] for problem (1) on a ball with p = 1 and a < 0: in fact, they show that the value at the origin of the residual mass of any non-compact sequence of radially symmetric positive solutions must be equal to $-\frac{a}{2}$ (and we get $\frac{1}{2}$, when a = -1).

Actually, we can give a more precise description of the asymptotic behavior of the solution u_{ε} as $\varepsilon \to 0$, since it is build via a Lyapunov-Schmidt procedure. For $\delta, \mu > 0$, and $\xi \in \mathbb{R}^n$, let us consider the functions

$$U_{\delta,\mu,\xi}(x) = \log\left(\frac{8\mu^2\delta^2}{(\mu^2\delta^2 + |x - \xi|^2)^2}\right),$$
(8)

which describe the set of all the solutions to the Liouville equation

$$-\Delta U = e^U \quad \text{in } \mathbb{R}^2, \tag{9}$$

under the condition $e^U \in L^1(\mathbb{R}^2)$ (see [21, 12]). We further consider the projection $PU_{\delta,\mu,\xi} := (-\Delta)^{-1} e^{U_{\delta,\mu,\xi}}$, where $(-\Delta)^{-1} : L^2(\Omega) \to H^1_0(\Omega)$ is the inverse of $-\Delta$. Namely, $PU_{\delta,\mu,\xi}$ is defined as the unique solution to

$$\begin{cases} -\Delta P U_{\delta,\mu,\xi} = -\Delta U_{\delta,\mu,\xi} = e^{U_{\delta,\mu,\xi}} & \text{in } \Omega, \\ P U_{\delta,\mu,\xi} = 0 & \text{on } \partial\Omega. \end{cases}$$
(10)

Intuitively, we want to look for solutions of (4) that look like $\alpha PU_{\delta,\mu,\xi} - u_0$ for suitable choices of the parameters α, δ, μ, ξ . Unfortunately, in order to successfully perform Lyapunov-Schmidt reduction, a more precise ansatz is necessary and we are forced to replace u_0 with a better approximation of the solutions. First, the non-degeneracy assumption (A1) allows to find a positive solution $v_{\varepsilon} \in C^1(\overline{\Omega})$ of (4) such that

$$v_{\varepsilon} \to u_0 \qquad \text{in } C^1(\overline{\Omega}),$$

as $\varepsilon \to 0$. Then, we consider the function

$$V_{\varepsilon,\alpha,\xi} := v_{\varepsilon} + \alpha w_{\varepsilon,\xi} + \alpha^2 z_{\varepsilon,\xi},\tag{11}$$

where $\alpha \in (0,1)$ is a small positive parameter depending on ε, μ, ξ such that $\alpha \to 0$ as $\varepsilon \to 0$, and $w_{\varepsilon,\xi}$ and $z_{\varepsilon,\xi}$ are defined as the unique solutions to the couple of linear problems

$$\begin{cases} \Delta w_{\varepsilon,\xi} + \lambda f_{\varepsilon}'(v_{\varepsilon}) w_{\varepsilon,\xi} = 8\pi \lambda G_{\xi} f_{\varepsilon}'(v_{\varepsilon}) & \text{in } \Omega, \\ w_{\varepsilon,\xi} = 0 & \text{on } \partial\Omega, \end{cases}$$
(12)

and

$$\begin{cases} \Delta z_{\varepsilon,\xi} + \lambda f_{\varepsilon}'(v_{\varepsilon}) z_{\varepsilon,\xi} = \frac{\lambda}{2} f_{\varepsilon}''(-v_{\varepsilon}) (8\pi G_{\xi} - w_{\varepsilon})^2 & \text{in } \Omega, \\ z_{\varepsilon} = 0 & \text{on } \partial\Omega, \end{cases}$$
(13)

with G_{ξ} denoting the Green function of Ω with singularity at ξ , namely the distributional solution to

$$\begin{cases} -\Delta G_{\xi} = \delta_{\xi} & \text{in } \Omega, \\ G_{\xi} = 0 & \text{on } \partial\Omega. \end{cases}$$
(14)

Problems (12) and (13) are nothing but the linearization of problem (4) around the solution v_{ε} and the R.H.S.'s are the terms of the second order Taylor's expansion with respect to α of $f_{\varepsilon}(\alpha PU_{\delta,\mu,\xi} - V_{\varepsilon,\alpha,\xi})$ far away from the concentration point ξ (indeed $PU_{\delta,\mu,\xi} \sim 8\pi G_{\xi}$ because of (23)).

Theorem 1.1 follows at once by the following result:

Theorem 1.3 Let λ , u_0 , ξ_0 be as in Theorem 1.1. There exists $\epsilon_0 > 0$ and functions $\alpha, \delta, \mu : (0, \varepsilon_0) \to (0, +\infty), \xi : (0, \varepsilon_0) \to \Omega$ and $\varphi : (0, \varepsilon_0) \to H_0^1(\Omega)$ such that:

- $u_{\varepsilon} := \alpha(\varepsilon) P U_{\delta(\varepsilon), \mu(\varepsilon), \xi(\varepsilon)} V_{\varepsilon, \alpha(\varepsilon), \xi(\varepsilon)} + \varphi(\varepsilon)$ is a solution (4).
- $\alpha(\varepsilon) \to 0, \ \delta(\varepsilon) \to 0, \ \mu(\varepsilon) \to \sqrt{8}e^{-1}, \ \xi(\varepsilon) \to \xi_0, \ and \ u_{\varepsilon}(\xi(\varepsilon)) \to +\infty \ as \ \varepsilon \to 0.$
- $\|\varphi(\varepsilon)\|_{H^1_0(\Omega)} + \|\varphi(\varepsilon)\|_{L^{\infty}(\Omega)} = O(e^{-\frac{\log(2u_0(\xi_0))}{\varepsilon}}).$

Let us briefly sketch the main steps of the proof of Theorem 1.3. First, in Section 2, we choose $\alpha = \alpha(\varepsilon, \mu, \xi)$ and $\delta = \delta(\varepsilon, \mu, \xi)$ such that the function

$$\omega_{\varepsilon,\mu,\xi} := \alpha P U_{\delta,\mu,\xi} - V_{\varepsilon,\alpha,\xi} \tag{15}$$

is an approximate solution of (4). Then, we look for solutions of (4) of the form $\omega_{\varepsilon,\mu,\xi} + \varphi$ with $\varphi \in H_0^1(\Omega)$. Clearly, (4) can be written in terms of φ as

$$-\Delta\varphi - \lambda f_{\varepsilon}'(\omega_{\varepsilon,\mu,\xi})\varphi = R + N(\varphi), \qquad (16)$$

where the error term R is defined by

$$R = R_{\varepsilon,\mu,\xi} := \Delta \omega_{\varepsilon,\mu,\xi} + \lambda f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi}), \qquad (17)$$

and the higher order term N by

$$N(\varphi) = N_{\varepsilon,\mu,\xi}(\varphi) := \lambda \left(f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi} + \varphi) - f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi}) - f'_{\varepsilon}(\omega_{\varepsilon,\mu,\xi})\varphi \right).$$
(18)

Equivalently, introducing the linear operator

$$L\varphi = L_{\varepsilon,\mu,\xi}\varphi := \varphi - (-\Delta)^{-1} (\lambda f'(\omega_{\varepsilon,\mu,\xi})\varphi),$$
(19)

we need to solve

$$L\varphi = (-\Delta)^{-1} \left(R + N(\varphi) \right).$$
⁽²⁰⁾

A careful and delicate estimate of the error R will be given in Section 3. The behaviour of the operator L will be studied in Section 4. On the one hand, for functions supported away from a suitable schrinking neighborhood of ξ , we will show that L is close to the operator $L_1\varphi := \varphi - (-\Delta)^{-1}(\lambda f'_0(u_0)\varphi)$, which is invertible on $H^1_0(\Omega)$ because of the non-degeneracy assumption (A1). On the other hand, near the point ξ , L is close to the operator $L_0\varphi := \varphi - (-\Delta)^{-1}(e^{U_{\delta,\mu,\xi}}\varphi)$. This operator appears in the analysis of several critical problems in dimension 2 (see for example [10, 13, 18]) and its behavior is well known: although L_0 is not invertible, it is possible to find an approximate threedimensional kernel $K_{\delta,\mu,\xi}$ for L_0 by projecting on $H^1_0(\Omega)$ the three functions

$$Z_{0,\delta,\mu,\xi}(x) = \frac{\delta^2 \mu^2 - |x - \xi|^2}{|x - \xi|^2 + \delta^2 \mu^2}, \qquad Z_{i,\delta,\mu,\xi}(x) = \frac{2\delta\mu(x_i - \xi_i)}{|x - \xi|^2 + \delta^2 \mu^2}, \qquad i = 1, 2.$$

Such properties transfer to the operator L, which turns out to be invertible on the subspace $K_{\delta,\mu,\xi}^{\perp}$ orthogonal to $K_{\delta,\mu,\xi}$ in $H_0^1(\Omega)$. More precisely, denoting by π and π^{\perp} the projections of $H_0^1(\Omega)$ respectively on $K_{\delta,\mu,\xi}$ and $K_{\delta,\mu,\xi}^{\perp}$, we will show that $\pi^{\perp}L$ is invertible on $K_{\delta,\mu,\xi}^{\perp}$. Then, it is natural to split equation (20) as

$$\begin{cases} \varphi = (\pi^{\perp}L)^{-1}\pi^{\perp} (-\Delta)^{-1} (R + N(\varphi)), \\ \pi L\varphi = \pi (-\Delta)^{-1} (R + N(\varphi)). \end{cases}$$
(21)

The first equation of (21) will be solved in Section 5, where for any $\mu > 0$, ξ close to ξ_0 and any small $\varepsilon > 0$, we will find a solution $\varphi_{\varepsilon,\mu,\xi}$ via a contraction mapping argument on a sufficiently small ball in $K_{\delta,\mu,\xi}^{\perp} \cap L^{\infty}(\Omega)$. Then, recalling that dim $K_{\delta,\mu,\xi} = 3$ and using assumption (A2), we will show in Section 6 that it is possible to choose the three parameters $\mu = \mu(\varepsilon)$ and $\xi = \xi(\varepsilon) = (\xi_1(\varepsilon), \xi_2(\varepsilon))$ so that the second equation in (21) is also fullfilled. Clearly, for such choice of μ and ξ , the function $\varphi_{\varepsilon,\mu(\varepsilon),\xi(\varepsilon)}$ solves both the equations in (21) (or, equivalently (16) and (20)), and $u_{\varepsilon} := \omega_{\varepsilon,\mu(\varepsilon),\xi(\varepsilon)} + \varphi_{\varepsilon,\mu(\varepsilon),\xi(\varepsilon)}$ is a solution of (4).

It is important to point out that choice of the concentration point $\xi(\varepsilon)$ is extremely delicate since the scaling parameter δ turns out to be much smaller than the parameter α , whose powers control all the error terms. To overcome this difficulty, we introduce a new argument based on a precise Pohozaev-type identity. This allows us to bypass global a priori gradient estimates on the solution $\varphi_{\varepsilon,\mu,\xi}$, which are hard to obtain for Moser-Trudinger critical problems. Our argument requires a very precise ansatz of the approximate solution $\omega_{\varepsilon,\mu,\xi}$. In particular, the presence of the correction terms $w_{\varepsilon,\xi}$ and $z_{\varepsilon,\xi}$ in the expression of $V_{\varepsilon,\alpha,\xi}$ is not merely technical, but plays a crucial role both in the estimates of the error term R and in the choice of $\xi(\varepsilon)$.

2 Construction of the approximate solution

In this section we give the detailed construction of the approximate solution $\omega_{\varepsilon,\mu,\xi}$. Here and in the rest of the paper, we will assume that $(\mu,\xi) \in \mathcal{U} \times B(\xi_0,\sigma)$, where $\mathcal{U} \in \mathbb{R}^+$ is an open interval containing $\mu_0 := \sqrt{8}e^{-1}$, ξ_0 is as in the assumption (A2), and $0 < \sigma < \frac{1}{2}d(\xi_0,\partial\Omega)$. By (A2), we can also assume

$$\inf_{B(\xi_0,\sigma)} u_0(\xi) > \frac{1}{2}.$$
(22)

2.1 The main terms of the ansatz

Let us introduce the main property of the projection of the bubble $PU_{\delta,\mu,\xi}$ defined in (10), which gives the main term of the approximate solution close to the blow-up point ξ . Let $G_{\xi}(\cdot) = G(\cdot,\xi)$ be the Green's function of $-\Delta$ with Dirichlet boundary conditions introduced in (14) and let $H(\cdot,\xi)$ be its regular part, i.e.

$$H(x,\xi) := G_{\xi}(x) - \frac{1}{2\pi} \log \frac{1}{|x-\xi|}$$

Lemma 2.1 We have

$$PU_{\delta,\mu,\xi}(x) = U_{\delta,\mu,\xi}(x) - \log(8\mu^2\delta^2) + 8\pi H(x,\xi) + \psi_{\delta,\mu,\xi}(x),$$

where

$$\|\psi_{\delta,\mu,\xi}\|_{C^1(\overline{\Omega})} = O(\delta^2),$$

uniformly with respect to $\mu \in \mathcal{U}, \xi \in B(\xi_0, \sigma)$. In particular,

$$PU_{\delta,\mu,\xi} \to 8\pi G_{\xi} \text{ in } C^1_{loc}(\overline{\Omega} \setminus \{\xi\}).$$

$$(23)$$

Proof. See for example [11, Proosition 5.1].

Next, let us define the main term of the approximate solution in the whole domain as $\alpha PU_{\delta,\mu,\xi} - v_{\varepsilon}$ where α is a positive parameter approaching zero as $\varepsilon \to 0$ and v_{ε} is a non-degenerate solution to (4), whose existence is proved in the following lemma.

Lemma 2.2 Let λ and u_0 be as in Theorems 1.1 and 1.3. There exists $\varepsilon_0 > 0$, and a family of functions $(v_{\varepsilon})_{0 < \varepsilon < \varepsilon_0} \subseteq C^1(\overline{\Omega})$ such that:

- *i.* v_{ε} *is a non-degenerate weak solution of* (4) *for any* $\varepsilon \in (0, \varepsilon_0)$ *.*
- ii. $v_{\varepsilon} \to u_0$ in $C^1(\overline{\Omega})$ as $\varepsilon \to 0$.

iii. There exists c > 0 such that $v_{\varepsilon}(x) \ge cd(x, \partial \Omega)$ for any $x \in \Omega$, $\varepsilon \in (0, \varepsilon_0)$.

Proof. Let $F: (-1,1) \times H^1_0(\Omega) \to H^1_0(\Omega)$ be defined by

$$F(\varepsilon, u) = F_{\varepsilon}(u) := u - (-\Delta)^{-1} (\lambda f_{\varepsilon}(u)), \qquad (24)$$

where f_{ε} is defined as in (5). F is well defined because the Moser-Trudinger inequality (2) implies that $f_{\varepsilon}(u) \in L^{p}(\Omega)$ for any $1 \leq p < +\infty$ and $u \in H_{0}^{1}(\Omega)$. Moreover, it is a C^{1} -map and its partial derivative $DF_{\varepsilon}(u) : H_{0}^{1}(\Omega) \to H_{0}^{1}(\Omega)$ defined by

$$DF_{\varepsilon}(u)[\varphi] = \varphi - (-\Delta)^{-1} (\lambda f'_{\varepsilon}(u)\varphi)$$

is a Fredholm operator of index 0 (since the embedding $H_0^1(\Omega) \hookrightarrow L^p(\Omega)$ is compact).

Now, let u_0 be a non-degenerate weak solution of (6) such that (A1) holds true. In particular, $F_0(u_0) = 0$ and $DF_0(u_0)$ is invertible. Therefore, by the implicit function theorem, we can construct a C^1 curve $\varepsilon \mapsto v_{\varepsilon} \in H_0^1(\Omega)$, defined for $|\varepsilon| < \varepsilon_0$ such that $v_0 = u_0, F_{\varepsilon}(v_{\varepsilon}) = 0$, and $DF_{\varepsilon}(v_{\varepsilon})$ is invertible for $|\varepsilon| < \varepsilon_0$. Then *i*. holds.



Applying the Moser-Trudinger inequality (2) and standard elliptic estimates, we obtain ii.

Hopf's lemma and the compactness of $\partial\Omega$ give $\frac{\partial u_0}{\partial \nu} \leq -2c$ on $\partial\Omega$, for some c > 0. Then, for ε sufficiently small, we have $\frac{\partial v_{\varepsilon}}{\partial \nu} \leq -c$, which in turn gives $v_{\varepsilon}(x) \geq cd(x,\partial\Omega)$ for x in a neighborhood of $\partial\Omega$. Finally, since $v_{\varepsilon} \to u_0$ uniformly in $\overline{\Omega}$, and $u_0 > 0$ in Ω , we get *iii*.

2.2 The correction of the ansatz

We need to correct the ansatz in the whole domain by solving the following two linear problems (12) and (13):

$$\begin{cases} \Delta w_{\varepsilon,\xi} + \lambda f_{\varepsilon}'(v_{\varepsilon})w_{\varepsilon,\xi} = 8\pi\lambda G_{\xi}f_{\varepsilon}'(v_{\varepsilon}) & \text{in }\Omega, \\ w_{\varepsilon,\xi} = 0 & \text{on }\partial\Omega \end{cases}$$

and

$$\begin{cases} \Delta z_{\varepsilon,\xi} + \lambda f_{\varepsilon}'(v_{\varepsilon}) z_{\varepsilon,\xi} = \frac{\lambda}{2} f_{\varepsilon}''(-v_{\varepsilon}) (8\pi G_{\xi} - w_{\varepsilon})^2 & \text{in } \Omega, \\ z_{\varepsilon} = 0 & \text{on } \partial\Omega. \end{cases}$$

Lemma 2.3 For any $0 < \varepsilon < \varepsilon_0$ and any $\xi \in \Omega$, there exist $w_{\varepsilon,\xi}$, $z_{\varepsilon,\xi}$ such that (12) and (13) hold. Moreover, there exists C > 0 such that

$$\|w_{\varepsilon,\xi}\|_{C^1(\overline{\Omega})} + \|z_{\varepsilon,\xi}\|_{C^1(\overline{\Omega})} \le C \tag{25}$$

for $\varepsilon \in (0, \varepsilon_0), \xi \in \Omega$.

Proof. The existence of the solutions immediately follows from the non-degeneracy of the function v_{ε} proved in Lemma 2.2. Moreover, since for any $p \in [1, +\infty)$ one has

$$\sup_{\xi \in \Omega} \|G_{\xi}\|_{L^{p}(\Omega)} < +\infty \quad \text{and} \quad \sup_{0 < \varepsilon < \varepsilon_{0}} \|v_{\varepsilon}\|_{C^{1}(\overline{\Omega})} < +\infty,$$

(25) follows by standard elliptic estimates.

Finally, we introduce the corrected ansatz as

$$\omega_{\varepsilon,\mu,\xi} := \alpha P U_{\delta,\mu,\xi} - V_{\varepsilon,\alpha,\xi} \tag{26}$$

with

$$V_{\varepsilon,\alpha,\xi} := v_{\varepsilon} + \alpha w_{\varepsilon,\xi} + \alpha^2 z_{\varepsilon,\xi},\tag{27}$$

where v_{ε} is defined in Lemma 2.2 and $w_{\varepsilon,\xi}$ and $z_{\varepsilon,\xi}$ as in Lemma 2.3.

2.3 The choice of parameters

It will be necessary to choose the parameters $\alpha = \alpha(\varepsilon, \mu, \xi)$ and $\delta = \delta(\varepsilon, \mu, \xi)$ such that $\lambda f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi}) \sim \alpha e^{U_{\delta,\mu,\xi}}$ when $|x - \xi| \sim \delta$. We point out that one of the main difficulties in this problem is that this estimates holds true only at a very small scale.

Let us fix the values of α and δ according to the next lemma. The proof is based on the contraction mapping theorem and is postponed to the appendix.

Lemma 2.4 There exist $\varepsilon_0 > 0$ and functions $\alpha = \alpha(\varepsilon, \mu, \xi)$, $\beta = \beta(\varepsilon, \mu, \xi)$ and $\delta = \delta(\varepsilon, \mu, \xi)$, defined in $(0, \varepsilon_0) \times \mathcal{U} \times B(\xi_0, \sigma)$ and continuous with respect to μ and ξ , such that

$$\begin{cases} \lambda \beta e^{\beta^2 + \beta^{1+\varepsilon}} = \frac{\alpha}{\delta^2}, \\ 2\alpha\beta + \alpha\beta^{\varepsilon} + \varepsilon\alpha\beta^{\varepsilon} = 1, \\ \beta = 4\alpha \log \frac{1}{\delta} - V_{\varepsilon,\alpha,\xi}(\xi) + \alpha c_{\mu,\xi}, \end{cases}$$
(28)

where $c_{\mu,\xi} := -\log(8\mu^2) + 8\pi H(\xi,\xi)$ and $V_{\varepsilon,\alpha,\xi}$ is defined in (11). Moreover, as $\varepsilon \to 0$, we have that

$$\alpha(\varepsilon,\mu,\xi) = \frac{1}{2}e^{-\frac{\log(2u_0(\xi)) + o(1)}{\varepsilon}},\tag{29}$$

$$\beta(\varepsilon,\mu,\xi) = \frac{1}{2\alpha} - u_0(\xi) + o(1), \qquad (30)$$

$$\log \frac{1}{\delta(\varepsilon, \mu, \xi)} = \frac{1 + o(1)}{8\alpha^2},\tag{31}$$

where $o(1) \to 0$ as $\varepsilon \to 0$, uniformly for $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$.

Remark 2.5 Note that (29)-(31) and (22) give $\alpha(\varepsilon, \mu, \xi), \delta(\varepsilon, \mu, \xi) \to 0$ and $\beta(\varepsilon, \mu, \xi) \to +\infty$ as $\varepsilon \to 0$, uniformly for $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$.

From now on we let $\alpha = \alpha(\varepsilon, \mu, \xi)$, $\beta = \beta(\varepsilon, \mu, \xi)$ and $\delta = \delta(\varepsilon, \mu, \xi)$ be as in Lemma 2.4.

It will be convenient to work on the scaled domain $\frac{\Omega-\xi}{\delta} := \left\{\frac{x-\xi}{\delta}, x \in \Omega\right\}$. Note that we have the scaling relation

$$U_{\delta,\mu,\xi}(x) = \bar{U}_{\mu}\left(\frac{x-\xi}{\delta}\right) - 2\log\delta, \qquad (32)$$

where

$$\bar{U}_{\mu}(y) = U_{1,\mu,0}(y) = \log\left(\frac{8\mu^2}{(\mu^2 + |y|^2)^2}\right).$$
(33)

Lemma 2.6 As $\varepsilon \to 0$, we have

$$\omega_{\varepsilon,\mu,\xi}(\xi + \delta y) = \beta + \alpha \bar{U}_{\mu}(y) + O(\delta|y|) + O(\delta^2), \tag{34}$$

uniformly for $y \in B(0, \frac{\sigma}{\delta})$, $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$. Moreover, for any R > 0 it holds also true that

$$\lambda f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi})(\xi + \delta y) = \alpha e^{U_{\delta,\mu,\xi}(\xi + \delta y)} (1 + O(\alpha^2)), \tag{35}$$

as $\varepsilon \to 0$ uniformly for $y \in B(0, R)$, $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$.

Proof. Lemma 2.1 and the scaling relation (32) show that, as $\delta \to 0$, we have the following expansion uniformly for $\varepsilon \in (0, \varepsilon_0)$, $\mu \in \mathcal{U}, \xi \in B(\xi_0, \sigma)$ and $y \in B(0, \frac{\sigma}{\delta})$:

$$\omega_{\varepsilon,\mu,\xi}(\xi+\delta y) = \alpha \bar{U}_{\mu} + \underbrace{4\alpha \log \frac{1}{\delta} + \alpha c_{\mu,\xi} - V_{\varepsilon,\alpha,\mu}(\xi)}_{=\beta} + V_{\varepsilon,\alpha,\mu}(\xi) - V_{\varepsilon,\alpha,\xi}(\xi+\delta y) + \delta \pi \alpha (H(\xi+\delta y,\xi) - H(\xi,\xi)) + O(\delta^2).$$

By Lemmas 2.2 and 2.3, we know that $V_{\varepsilon,\alpha,\mu}$ is uniformly bounded in $C^1(\overline{\Omega})$. Thus

 $V_{\varepsilon,\alpha,\mu}(\xi + \delta y) = V_{\varepsilon,\alpha,\mu}(\xi) + O(\delta|y|).$

Similarly, since $H \in C^1(\overline{\Omega} \times B(\xi_0, \sigma))$, we have

$$H(\xi + \delta y, \xi) = H(\xi, \xi) + O(\delta|y|).$$

Then estimate (34) is proved.

Now, let us prove (35). Note that (29)-(31) yield $\beta = O(\frac{1}{\alpha}), \ \delta = O(e^{-\frac{1+o(1)}{8\alpha^2}})$, and $\beta^{\varepsilon} = 2u_0(\xi) + o(1) = O(1)$. For $|y| \leq R$, (34) implies

$$\omega_{\varepsilon,\mu,\xi}(\xi + \delta y) = \beta + \alpha U_{\mu}(y) + O(\delta).$$

In particular

$$\omega_{\varepsilon,\mu,\xi}(\xi+\delta y)^2 = \beta^2 + 2\alpha\beta\bar{U}_{\mu}(y) + O(\beta\delta), \qquad (36)$$

and

$$\omega_{\varepsilon,\mu,\xi}(\xi + \delta y)^{1+\varepsilon} = (\beta + \alpha \bar{U}_{\mu}(y) + O(\delta))(\beta + \alpha \bar{U}_{\mu}(y) + O(\delta))^{\varepsilon}
= (\beta + \alpha \bar{U}_{\mu}(y) + O(\delta))\beta^{\varepsilon} \left(1 + \frac{\alpha}{\beta}\bar{U}_{\mu}(y) + O(\alpha\delta)\right)^{\varepsilon}
= (\beta^{1+\varepsilon} + \alpha\beta^{\varepsilon}\bar{U}_{\mu}(y) + O(\delta)) \left(1 + \frac{\varepsilon\alpha}{\beta}\bar{U}_{\mu}(y) + O(\varepsilon\alpha^{4})\right)
= \beta^{1+\varepsilon} + \alpha\beta^{\varepsilon}\bar{U}_{\mu}(y) + \varepsilon\alpha\beta^{\varepsilon}\bar{U}_{\mu}(y) + O(\varepsilon\alpha^{3}).$$
(37)

Then, using (28) we get

$$\begin{split} \lambda f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi})(\xi+\delta y) &= \lambda \omega_{\varepsilon,\mu,\xi}(\xi+\delta y)e^{\omega_{\varepsilon,\mu,\xi}(\xi+\delta y)^2 + \omega_{\varepsilon,\mu,\xi}^{1+\varepsilon}(\xi+\delta y)} \\ &= \lambda \beta (1+O(\alpha^2))e^{\beta^2 + \beta^{1+\varepsilon} + (2\alpha\beta + \alpha\beta^{\varepsilon} + \alpha\varepsilon\beta^{\varepsilon})\bar{U}_{\mu}(y) + O(\alpha^2)} \\ &= \underbrace{\lambda \beta e^{\beta^2 + \beta^{1+\varepsilon}}}_{=\frac{\alpha}{\delta^2}} e^{\underbrace{(2\alpha\beta + \alpha\beta^{\varepsilon} + \alpha\varepsilon\beta^{\varepsilon})}_{=1}\bar{U}_{\mu}(y)} (1+O(\alpha^2))e^{O(\alpha^2)} \\ &= \frac{\alpha}{\delta^2} e^{\bar{U}_{\mu}(y)}(1+O(\alpha^2)) \\ &= \alpha e^{U_{\delta,\mu,\xi}(\xi+\delta y)}(1+O(\alpha^2)), \end{split}$$

which proves (35).

It is also useful to point out the following result which will be used in the next sections.

Remark 2.7 Lemma 2.1 and Lemma 2.4 give

$$0 \le \alpha P U_{\delta,\mu,\xi} \le \beta + u_0(\xi) + o(1),$$

and

$$-V_{\alpha,\varepsilon,\xi} \le \omega_{\varepsilon,\mu,\xi} \le \beta + o(1),$$

uniformly for $x \in \Omega$, $\varepsilon \in (0, \varepsilon_0)$, $\mu \in \mathcal{U}$, $\xi \in B(\xi_0, \sigma)$.

Notation: In order to simplify the notation, we will write U_{ε} , U, V_{ε} , ω_{ε} , w_{ε} and z_{ε} instead of $U_{\delta,\mu,\xi}$, \bar{U}_{μ} , $V_{\varepsilon,\alpha,\xi}$, $\omega_{\varepsilon,\mu,\xi}$, $w_{\varepsilon,\xi}$ and $z_{\varepsilon,\xi}$, without specifying explicitly the dependence on the parameters. It is important to point out that all the estimates of the next sections will be uniform with respect to $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$. This will allow us to choose freely the values of μ and ξ in Section 6. Consistently, the notation $O(f(x,\varepsilon,\alpha,\beta,\delta))$ and $o(f(x,\varepsilon,\alpha,\beta,\delta))$ will be used for quantities depending on ε,ξ,μ (and the parameters α,β,δ of Lemma 2.4) and satisfying respectively

$$|O(f(x,\varepsilon,\mu,\xi\alpha,\beta,\delta))| \le Cf(x,\varepsilon,\mu,\xi,\alpha,\beta,\delta)) \quad \text{and} \quad \frac{o(f(x,\varepsilon,\mu,\xi\alpha,\beta,\delta))}{f(x,\varepsilon,\mu,\xi,\alpha,\beta,\delta)} \to 0,$$

as $\varepsilon \to 0$, uniformly for $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$.

3 The estimate of the error term

In this section we give estimates for the error term R defined in (17)

$$R = R_{\varepsilon,\mu,\xi} := \Delta \omega_{\varepsilon,\mu,\xi} + \lambda f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi}).$$

It will be convenient to split Ω into four different regions:

$$\Omega = B(\xi, \rho_0) \cup \left(B(\xi, \rho_1) \setminus B(\xi, \rho_0) \right) \cup \left(B(\xi, \rho_2) \setminus B(\xi, \rho_1) \right) \cup \left(\Omega \setminus B(\xi, \rho_2) \right), \quad (38)$$

where $\rho_0 = \rho_0(\varepsilon, \mu, \xi)$, $\rho_1 = \rho_1(\varepsilon, \mu, \xi)$, $\rho_2 = \rho_2(\varepsilon, \mu, \xi)$, are defined by

$$\rho_0 = \delta e^{\frac{\varepsilon}{\alpha}}, \quad \rho_1 = e^{-\frac{u_0(\xi)}{2\alpha}} \quad \text{and} \quad \rho_2 = e^{-\frac{\varepsilon}{\alpha}}.$$
(39)

Note that

$$\delta \ll \rho_0 \ll \rho_1 \ll \rho_2 \ll 1, \qquad \text{as } \varepsilon \to 0,$$

by (29) and (31). Roughly speaking, we have to split the error into four parts: in $B(\xi, \rho_0)$ we have $\lambda f_{\varepsilon}(\omega_{\varepsilon}) = \alpha e^{U_{\varepsilon}}(1 + o(1))$ (see (35)) and we can use a blow-up argument to get a uniform weighted estimate on R. This estimate does not hold anymore in the set $\Omega \setminus B(\xi, \rho_0)$, which we further split into three parts: the region $\Omega \setminus B(\xi, \rho_2)$, where $\alpha G_{\xi} = O(\varepsilon)$ and a uniform estimate on R can be obtained via a Taylor expansion of $f_{\varepsilon}(\omega_{\varepsilon})$ (using that $\omega_{\varepsilon} = -V_{\varepsilon} + 8\pi\alpha G_{\xi} + o(\alpha^2)$), and the two annuli $B(\xi, \rho_1) \setminus B(\xi, \rho_0)$ and $B(\xi, \rho_2) \setminus B(\xi, \rho_1)$, where we give quite delicate integral estimates. The last two regions are treated separately since $\omega_{\varepsilon} \ge c_0 > 0$ in $B(\xi, \rho_1) \setminus B(\xi, \rho_0)$, while ω_{ε} changes sign in $B(\xi, \rho_2) \setminus B(\xi, \rho_1)$ (cfr. Lemma 3.2 and Lemma 3.11).

3.1 A uniform expansion in $B(\xi, \rho_1)$

In this section we give a more precise version of the expansions in (36)-(37).

Lemma 3.1 For any $\varepsilon \in (0,1)$ and $x \ge -1$, we have

$$|(1+x)^{1+\varepsilon} - 1 - (1+\varepsilon)x| \le \varepsilon x^2.$$

Proof. According to Bernoulli's inequality we have

$$(1+x)^{\varepsilon} \le 1 + \varepsilon x \tag{40}$$

and

$$(1+x)^{1+\varepsilon} \ge 1 + (1+\varepsilon)x. \tag{41}$$

Since $x \ge -1$, thanks to (40) we have that

$$(1+x)^{1+\varepsilon} \le (1+x)(1+\varepsilon x) = 1 + (1+\varepsilon)x + \varepsilon x^2.$$
(42)

Then, the conclusion follows from (41) and (42).

Lemma 3.2 Set $c_0 := \frac{1}{2} \inf_{\xi \in B(\xi_0, \sigma)} u_0(\xi)$. For $x \in B(\xi, \rho_1)$, we have that

$$\beta + \alpha \bar{U}\left(\frac{x-\xi}{\delta}\right) \ge c_0,\tag{43}$$

for sufficiently small ε . In particular, we have

$$c_0 \le \omega_{\varepsilon} \le \beta (1 + o(1)). \tag{44}$$

Proof. The definitons of \overline{U} and ρ_1 (see (33) and (39)), and (30)-(31) give

$$\beta + \alpha \bar{U}\left(\frac{x-\xi}{\delta}\right) \ge \beta + \alpha \bar{U}\left(\frac{\rho_1}{\delta}\right)$$
$$= \beta - 4\alpha \log \frac{\rho_1}{\delta} + o(1)$$
$$= u_0(\xi) + o(1),$$

which implies (43) for sufficiently small ε . To get (44), it is sufficient to apply Lemma 2.6 and Remark 2.7.

Lemma 3.3 For $x \in B(\xi, \rho_1)$, we have

$$\omega_{\varepsilon}^{2}(x) + \omega_{\varepsilon}^{1+\varepsilon}(x) = \beta^{2} + \beta^{1+\varepsilon} + \bar{U}\left(\frac{x-\xi}{\delta}\right) + \alpha^{2}\bar{U}^{2}\left(\frac{x-\xi}{\delta}\right) + O\left(\varepsilon\alpha^{3}\left(1+\bar{U}^{2}\left(\frac{x-\xi}{\delta}\right)\right)\right).$$

Proof. Set $y = \frac{x-\xi}{\delta} \in B(0, \frac{\rho_1}{\delta})$. Noting that $\overline{U}(y) = O(\alpha^{-2})$ and using Lemma 2.6, we get

$$\omega_{\varepsilon}^{2}(x) = \omega_{\varepsilon}^{2}(\xi + \delta y) = \left(\beta + \alpha \bar{U}(y) + O(\rho_{1})\right)^{2}$$
$$= \beta^{2} + 2\alpha\beta \bar{U}(y) + \alpha^{2} \bar{U}(y)^{2} + O(\beta\rho_{1}).$$

Similarly, since Lemma 3.2 gives $\frac{\alpha}{\beta}\overline{U}(y) \ge -1 + \frac{c_0}{\beta} \ge -1$, by Lemma 3.1 we infer

$$\begin{split} |\omega_{\varepsilon}|^{1+\varepsilon}(x) &= \beta^{1+\varepsilon} \left(1 + \frac{\alpha}{\beta} \bar{U}(y) + O(\alpha\rho_1) \right)^{1+\varepsilon} \\ &= \beta^{1+\varepsilon} \left(1 + (1+\varepsilon) \left(\frac{\alpha}{\beta} \bar{U}(y) + O(\alpha\rho_1) \right) + O\left(\varepsilon \left(\frac{\alpha}{\beta} \bar{U}(y) + O(\alpha\rho_1) \right)^2 \right) \right) \\ &= \beta^{1+\varepsilon} + (1+\varepsilon) \alpha \beta^{\varepsilon} \bar{U}(y) + O(\varepsilon \alpha^3 (1+\bar{U}^2(y))). \end{split}$$

Then the conclusion follows from the second equation in (28).

3.2 Expansions in $B(\xi, \rho_0)$

Let us now restrict our attention to the smaller ball $B(\xi, \rho_0)$. This allows to control the term $\alpha^2 \bar{U}^2$ appearing in the expansion of Lemma 3.3. Indeed, since $|\bar{U}(y)| = -4 \log |y| + O(1)$ as $|y| \to +\infty$, we have that

$$\bar{U}\left(\frac{x-\xi}{\delta}\right) = O\left(\frac{\varepsilon}{\alpha}\right) \quad \text{and} \quad \alpha^2 \bar{U}^2\left(\frac{x-\xi}{\delta}\right) = O(\varepsilon^2) \quad \text{for } x \in B(\xi,\rho_0).$$
(45)

Lemma 3.4 For $x \in B(\xi, \rho_0)$, we have

$$R(x) = \alpha^3 e^{U_{\varepsilon}(x)} \left(2\bar{U}\left(\frac{x-\xi}{\delta}\right) + \bar{U}^2\left(\frac{x-\xi}{\delta}\right) \right) + \alpha^4 e^{U_{\varepsilon}(x)} O\left(1 + \bar{U}^4\left(\frac{x-\xi}{\delta}\right)\right).$$

Proof. Set $y = \frac{x-\xi}{\delta}$. First by Lemma 2.6, Lemma 3.3, and (28)-(32), we get that

$$\begin{split} \lambda f_{\varepsilon}(\omega_{\varepsilon}(x)) &= \lambda \beta \left(1 + \frac{\alpha}{\beta} \bar{U}(y) + O(\alpha \rho_1) \right) e^{\omega_{\varepsilon}^2(x) + \omega_{\varepsilon}^{1+\varepsilon}(x)} \\ &= \frac{\alpha}{\delta^2} \left(1 + 2\alpha^2 \bar{U}(y) + O(\alpha^3(1 + |\bar{U}(y)|)) \right) e^{\bar{U}(y) + \alpha^2 \bar{U}^2(y) + O(\varepsilon \alpha^3(1 + \bar{U}^2(y)))} \\ &= \alpha e^{U_{\varepsilon}(x)} \left(1 + 2\alpha^2 \bar{U}(y) + O(\alpha^3(1 + |\bar{U}(y)|)) \right) e^{\alpha^2 \bar{U}^2(y) + O(\varepsilon \alpha^3(1 + \bar{U}^2(y)))}. \end{split}$$

Now, by (45), we can expand the last exponential term, and find

$$e^{\alpha^2 \bar{U}^2(y) + O(\varepsilon \alpha^3 (1 + \bar{U}^2(y)))} = 1 + \alpha^2 \bar{U}^2(y) + O(\varepsilon \alpha^3 (1 + \bar{U}^2(y))) + O(\alpha^4 (1 + \bar{U}^4(y)))$$

= 1 + \alpha^2 \bar{U}^2(y) + O(\varepsilon \alpha^3 (1 + \bar{U}^4(y))).

We can so conclude that

$$\lambda f_{\varepsilon}(\omega_{\varepsilon}(x)) = \alpha e^{U_{\varepsilon}(x)} + \alpha^3 e^{U_{\varepsilon}(x)} \left(2\bar{U}(y) + \bar{U}(y)^2 \right) + \alpha^4 e^{U_{\varepsilon}(x)} O(1 + \bar{U}^4(y)).$$
(46)

Moreover, by (10)-(13), and Lemmas 2.2-2.3 we have

$$\Delta\omega_{\varepsilon} = -\alpha e^{U_{\varepsilon}} + O(1) = -\alpha e^{U_{\varepsilon}} \left(1 + O(\alpha) e^{-U_{\varepsilon}} \right) = -\alpha e^{U_{\varepsilon}} (1 + o(\alpha^3)), \quad (47)$$

where in the last equality we used that

$$e^{-U_{\varepsilon}(x)} = \frac{(\delta^{2}\mu^{2} + |x - \xi|^{2})^{2}}{8\delta^{2}\mu^{2}} = O(\delta^{2}e^{\frac{4\varepsilon}{\alpha}}) = o(\alpha^{3}),$$

for $x \in B(\xi, \rho_0)$. Thanks to (46) and (47), we conclude that

$$R(x) = \alpha^3 e^{U_{\varepsilon}(x)} \left(2\bar{U}(y) + \bar{U}^2(y) \right) + \alpha^4 e^{U_{\varepsilon}(x)} O(1 + \bar{U}^4(y)).$$

As an immediate consequence of the previous lemma we obtain the estimate:

Corollary 3.5 We have that

$$R = O\left(\alpha^3 e^{U_{\varepsilon}} \left(\left(1 + \bar{U}^4 \left(\frac{\cdot - \xi}{\delta}\right)\right) \right)\right)$$

in $B(\xi, \rho_0)$.

3.3 Estimates on $B(\xi, \rho_1) \setminus B(\xi, \rho_0)$

In this region, it is diffcult to provide pointwise estimates of R because the term $\alpha^2 \overline{U}^2$ appearing in the expansion of Lemma 3.3 becomes very large. Then, we will look for integral estimates. Specifically we will show that R is (very) small in $L^p(B(\xi, \rho_1) \setminus B(\xi, \rho_0))$, for a suitable choice of $p = p(\alpha) > 1$, such that $p \to 1$ as $\varepsilon \to 0$, uniformly with respect to $\xi \in B(\xi_0, \sigma), \mu \in \mathcal{U}$.

Lemma 3.6 There exists $c_1 > 0$ such that

$$0 \le \lambda f_{\varepsilon}(\omega_{\varepsilon}) \le \alpha e^{U_{\varepsilon} + \alpha^2 (1 + c_1 \varepsilon \alpha) \bar{U}^2(\frac{\cdot - \xi}{\delta})},$$

in $B(\xi, \rho_1) \setminus B(\xi, \rho_0)$.

Proof. Since $0 \le \omega_{\varepsilon} \le \beta$ in $B(\xi, \rho_1) \setminus B(\xi, \rho_0)$, from Lemma 3.3 and (28) we get

$$\begin{split} \lambda f_{\varepsilon}(\omega_{\varepsilon}) &\leq \lambda \beta e^{\beta^2 + \beta^{1+\varepsilon} + \bar{U}(\frac{\cdots-\xi}{\delta}) + \alpha^2 \bar{U}^2(\frac{\cdots-\xi}{\delta})(1+O(\varepsilon\alpha))} \\ &= \frac{\alpha}{\delta^2} e^{\bar{U}(\frac{\cdots-\xi}{\delta}) + \alpha^2 \bar{U}^2(\frac{\cdots-\xi}{\delta})(1+O(\varepsilon\alpha))} \\ &= \alpha e^{U_{\varepsilon} + \alpha^2 \bar{U}^2(\frac{\cdots-\xi}{\delta})(1+O(\varepsilon\alpha))}. \end{split}$$

For c_1 as in Lemma 3.6, let us consider the function

$$\Gamma_{\varepsilon}(x) := e^{\bar{U}_{\varepsilon}(x) + \alpha^2 \bar{U}(\frac{x-\xi}{\delta})^2 (1+c_1 \varepsilon \alpha)}.$$
(48)

Lemma 3.7 Set $p := 1 + \alpha^2$. There exists $c_2 > 0$ such that

$$\|\Gamma_{\varepsilon}\|_{L^{p}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} = O\left(\alpha^{-1}e^{-\frac{c_{2}}{\sqrt{\alpha}}}\right)$$

Proof. First of all, we observe that for $q \in (\frac{1}{2}, +\infty)$, R > 0, one has

$$\int_{\mathbb{R}^2 \setminus B(0,R)} e^{q\overline{U}} dy \le \int_{\mathbb{R}^2 \setminus B(0,R)} \frac{(8\mu^2)^q}{|y|^{4q}} dy = \frac{\pi (8\mu^2)^q}{(2q-1)R^{4q-2}}.$$
(49)

For $x \in B(\xi, \rho_1) \setminus B(\xi, \rho_0)$, set $y = \frac{x-\xi}{\delta} \in B(0, \frac{\rho_1}{\delta}) \setminus B(0, \frac{\rho_0}{\delta})$. Clearly we have

$$\|\Gamma_{\varepsilon}\|_{L^{p}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} = \delta^{\frac{2-2p}{p}} \left(\int_{B(0,\frac{\rho_{1}}{\delta})\setminus B(0,\frac{\rho_{0}}{\delta})} e^{p\bar{U}(y)(1+\alpha^{2}\bar{U}(y)(1+c_{1}\varepsilon\alpha))} dy \right)^{\frac{1}{p}}.$$
 (50)

Set $\bar{\rho} = \delta e^{\frac{1}{\alpha^2}}$, so that $\rho_0 \ll \bar{\rho} \ll \rho_1$. For $\frac{\rho_0}{\delta} \le |y| \le \frac{\bar{\rho}}{\delta}$, we have

$$p\left(1+\alpha^2 \bar{U}(y)(1+\varepsilon c_1\alpha)\right) = 1 + O(\sqrt{\alpha}) \ge \frac{2}{3}$$

Then, for ε small enough, (49) yields

$$\int_{B(0,\frac{\bar{\rho}}{\delta})\setminus B(0,\frac{\rho_{0}}{\delta})} e^{p\bar{U}(y)\left(1+\alpha^{2}\bar{U}(y)(1+c_{1}\varepsilon\alpha)\right)} dy \leq \int_{\mathbb{R}^{2}\setminus B(0,\frac{\rho_{0}}{\delta})} e^{\frac{2}{3}\bar{U}(y)} dy = O\left(\left(\frac{\rho_{0,\varepsilon}}{\delta}\right)^{-\frac{2}{3}}\right) = O(e^{-\frac{2\varepsilon}{3\alpha}}).$$
(51)

For $\frac{\bar{\rho}}{\delta} \leq |y| \leq \frac{\rho_1}{\delta}$, by (30) and Lemma 3.2, we have

$$1 + \alpha^2 \bar{U}(y) \left(1 + c_1 \varepsilon \alpha\right) = 1 + \alpha (\beta + \alpha \bar{U}(y)) \left(1 + c_1 \varepsilon \alpha\right) - \alpha \beta \left(1 + c_1 \varepsilon \alpha\right)$$
$$\geq \frac{1}{2} + (c_0 + u_0(\xi))\alpha + o(\alpha)$$
$$\geq \frac{1}{2} + c_0 \alpha.$$

Hence, we get

$$\int_{B(0,\frac{\rho_{1}}{\delta})\setminus B(0,e^{\alpha^{-\frac{3}{2}}})} e^{p\bar{U}(y)\left(1+\alpha^{2}\bar{U}(y)(1+c_{1}\varepsilon\alpha)\right)} dy \leq \int_{\mathbb{R}^{2}\setminus B(0,e^{\alpha^{-\frac{3}{2}}})} e^{p\left(\frac{1}{2}+c_{0}\alpha\right)\bar{U}(y)} dy = O\left(\alpha^{-1}e^{-\frac{4c_{0}}{\sqrt{\alpha}}}\right).$$
(52)

Thus, by (50),(51),(52), we obtain

$$\|\Gamma_{\varepsilon}\|_{L^{p}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} = O\left(\delta^{\frac{2-2p}{p}}\alpha^{-\frac{1}{p}}e^{-\frac{4c_{0}}{p\sqrt{\alpha}}}\right)$$

Since (29)-(31) give

$$\delta^{\frac{2-2p}{p}} = \delta^{-\frac{2\alpha^2}{1+\alpha^2}} = O(1), \qquad \alpha^{\frac{1}{p}} = \alpha \alpha^{\frac{1-p}{p}} = \alpha(1+o(1)), \qquad e^{-\frac{4c_0}{p\sqrt{\alpha}}} = O(e^{-\frac{4c_0}{\sqrt{\alpha}}}),$$

we get the conclusion.

Lemma 3.8 Let p and c_2 be as in Lemma 3.7, then

$$\|R\|_{L^{p}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} = O(e^{-\frac{c_{2}}{\sqrt{\alpha}}}).$$

Proof. By Lemma 3.6 and Lemma 3.7 we get that

$$\|\lambda f_{\varepsilon}(\omega_{\varepsilon})\|_{L^{p}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} = O(e^{-\frac{c_{2}}{\sqrt{\alpha}}}).$$

On the other hand, we have

$$\Delta\omega_{\varepsilon}(x) = -\alpha e^{U_{\varepsilon}(y)} + O(1),$$

so that

$$\begin{split} \|\Delta\omega_{\varepsilon}\|_{L^{p}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} &\leq \alpha \|e^{U_{\varepsilon}}\|_{L^{p}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} + O(\rho_{1}^{\frac{2}{p}}) \\ &\leq \alpha \delta^{\frac{2-2p}{p}} \|e^{\bar{U}}\|_{L^{p}(\mathbb{R}^{2}\setminus B(0,\frac{\rho_{0}}{\delta}))} + O(\rho_{1}^{\frac{2}{p}}) \\ &= O\left(\frac{\alpha \delta^{2}}{\rho_{0}^{2}}\right) + O\left(\rho_{1}^{2}\right) \\ &= o(e^{-\frac{c_{2}}{\sqrt{\alpha}}}). \end{split}$$

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3.4 Estimates in $B(\xi, \rho_2) \setminus B(\xi, \rho_1)$

In $B(\xi, \rho_2) \setminus B(\xi, \rho_1)$ we can only say that ω_{ε} and R are uniformly bounded. Since ρ_2 is very small, we still get integral bounds for R.

Lemma 3.9 We have $\omega_{\varepsilon} = O(1)$ and R = O(1) in $\Omega \setminus B(\xi, \rho_1)$. In particular,

$$||R||_{L^2(B(\xi,\rho_2)\setminus B(\xi,\rho_1))} = O(\rho_2) = O(e^{-\frac{\varepsilon}{\alpha}}).$$

Proof. Let us recall that $\omega_{\varepsilon} = \alpha P U_{\varepsilon} - V_{\varepsilon}$ with $V_{\varepsilon} = V_{\varepsilon,\alpha,\xi}$ defined as in (11). According to Lemma 2.2 and Lemma 2.3, we have $V_{\varepsilon} = O(1)$ in Ω . Besides Lemma 2.1 gives

$$\alpha PU_{\varepsilon} = \alpha \log\left(\frac{1}{(\mu^2 \delta^2 + |x - \xi|^2)^2}\right) + O(\alpha) = O(\alpha \log \frac{1}{\rho_1}) + O(\alpha) = O(1),$$

for $x \in \Omega \setminus B(\xi, \rho_1)$. Then, $\omega_{\varepsilon} = O(1)$ and $f_{\varepsilon}(\omega_{\varepsilon}) = O(1)$ in $\Omega \setminus B(\xi, \rho_1)$. Similarly

$$\begin{split} \Delta \omega_{\varepsilon} &= -\alpha e^{U_{\varepsilon}} + O(1) \\ &= -\frac{\alpha \delta^2 \mu^2}{(\delta^2 \mu^2 + |x - \xi|^2)^2} + O(1) \\ &= O(\delta^2 \rho_1^{-4}) + O(1) = O(1). \end{split}$$

Therefore R = O(1).

3.5 Estimates in $\Omega \setminus B(\xi, \rho_2)$

In $\Omega \setminus B(\xi, \rho_2)$ we will use that $\omega_{\varepsilon} \sim 8\pi \alpha G_{\xi} - V_{\varepsilon}$. Our choice of V_{ε} will make R uniformly small, namely of order α^3 . Note further that the choice of ρ_2 gives $\alpha G_{\xi} = O(\varepsilon)$ on $\Omega \setminus B(\xi, \rho_2)$.

Lemma 3.10 As $\varepsilon \to 0$ we have

$$\|PU_{\varepsilon} - 8\pi G_{\xi}\|_{C^1(\overline{\Omega}\setminus B(\xi,\rho_2))} = O(\delta^2 \rho_2^{-3}).$$

Proof. By Lemma 2.1 we have

$$PU_{\varepsilon} = \log\left(\frac{1}{(\delta^{2}\mu^{2} + |x - \xi|^{2})^{2}}\right) + 8\pi H(x,\xi) + \psi_{\delta,\mu,\xi}$$

= $-4 \log|x - \xi| + 8\pi H(x,\xi) - 2 \log\left(1 + \frac{\delta^{2}\mu^{2}}{|x - \xi|^{2}}\right) + \psi_{\delta,\mu,\xi}$
= $8\pi G_{\xi}(x) - 2 \log\left(1 + \frac{\delta^{2}\mu^{2}}{|x - \xi|^{2}}\right) + \psi_{\delta,\mu,\xi}$

Since $\|\psi_{\delta,\mu,\xi}\|_{C^1(\overline{\Omega})} = O(\delta^2)$ as $\varepsilon \to 0$, it is sufficient to observe that

$$\|\log\left(1 + \frac{\delta^{2}\mu^{2}}{|\cdot -\xi|^{2}}\right)\|_{C^{1}(\overline{\Omega}\setminus B(\xi,\rho_{2}))} = O(\delta^{2}\rho_{2}^{-3})$$

Lemma 3.11 There exists a constant c > 0 such such that

$$\omega_{\varepsilon}(x) \le -c \, d(x, \partial \Omega) < 0.$$

for any $x \in \Omega \setminus B(\xi, \rho_2)$, provided ε is sufficiently small.

Proof. By Lemma 2.2, Lemma 2.3 and (11) we have

$$V_{\varepsilon}(x) \ge c(1+O(\alpha))d(x,\partial\Omega) \qquad \forall x \in \Omega,$$

for some c > 0. Then, Lemma 3.10 implies that

$$\omega_{\varepsilon}(x) \le -c(1+O(\alpha))d(x,\partial\Omega) \tag{53}$$

in a neighborhood of $\partial\Omega$. By definiton of ρ_2 , we have that $PU_{\varepsilon} = G_{\xi} + o(1) = O(\frac{\varepsilon}{\alpha})$ in $\Omega \setminus B(\xi, \rho_2)$. Then, using again Lemma 2.2 and Lemma 2.3, we get $\omega_{\varepsilon} = -u_0 + o(1)$ uniformly in $\Omega \setminus B(\xi, \rho_2)$. Since $u_0 > 0$ in Ω , this together with (53) yields the conclusion.

Lemma 3.12 In $\Omega \setminus B(\xi, \rho_2)$, we have $R = O(\alpha^3(1 + G_{\xi}^3))$. In particular,

$$||R||_{L^2(\Omega \setminus B(\xi,\rho_2))} = O(\alpha^3)$$

Proof. Since $v_{\varepsilon} > 0$ in Ω , $\omega_{\varepsilon} < 0$ in $\Omega \setminus B(\xi, \rho_2)$, and $f_{\varepsilon} \in C^3((-\infty, 0))$, for any $x \in \Omega \setminus B(\xi, \rho_2)$ we can find $\theta(x) \in [0, 1]$ such that

$$\begin{aligned} f_{\varepsilon}(\omega_{\varepsilon}) &= f_{\varepsilon}(-v_{\varepsilon} + \alpha P U_{\varepsilon} - \alpha w_{\varepsilon} - \alpha^{2} z_{\varepsilon}) \\ &= f_{\varepsilon}(-v_{\varepsilon}) + f_{\varepsilon}'(-v_{\varepsilon})(\alpha P U_{\varepsilon} - \alpha w_{\varepsilon} - \alpha^{2} z_{\varepsilon}) + \frac{1}{2} f_{\varepsilon}''(-v_{\varepsilon})(\alpha P U_{\varepsilon} - \alpha w_{\varepsilon} - \alpha^{2} z_{\varepsilon})^{2} \\ &+ \frac{1}{6} f'''(-v_{\varepsilon} + \theta(\alpha P U_{\varepsilon} - \alpha w_{\varepsilon} - \alpha^{2} z_{\varepsilon}))(\alpha P U_{\varepsilon} - \alpha w_{\varepsilon} - \alpha^{2} z_{\varepsilon})^{3} \end{aligned}$$

According to Lemma 2.3 and Lemma 3.10, we have

$$|z_{\varepsilon}| + |w_{\varepsilon}| = O(G_{\xi})$$
 and $\alpha PU_{\varepsilon} = 8\pi\alpha G_{\xi}(1 + o(\alpha^3)).$

Thus we get

$$f_{\varepsilon}(\omega_{\varepsilon}) = -f_{\varepsilon}(v_{\varepsilon}) + \alpha f_{\varepsilon}'(v_{\varepsilon})(8\pi G_{\xi} - w_{\varepsilon}) + \alpha^{2} \left(\frac{1}{2}f''(-v_{\varepsilon})(8\pi G_{\xi} - w_{\varepsilon})^{2} - f'(v_{\varepsilon})z_{\varepsilon}\right) + O(\alpha^{3}(1 + G_{\xi}^{3})) + O(\alpha^{3}|f'''(-v_{\varepsilon} + \theta(\alpha PU_{\delta,\mu} - \alpha w_{\varepsilon} - \alpha^{2}z_{\varepsilon}))|G_{\xi}^{3}).$$

A direct computation shows the existence of a constant C > 0 such that

$$|f_{\varepsilon}^{\prime\prime\prime\prime}(t)| \le C(|t|^{\varepsilon-1} + t^4)e^{t^2 + |t|^{1+\varepsilon}} \quad \forall t \neq 0.$$

Since $-v_{\varepsilon} + \theta(\alpha P U_{\varepsilon} - \alpha w_{\varepsilon} - \alpha^2 z_{\varepsilon}) = O(1)$ uniformly in $\Omega \setminus B(\xi, \rho_2)$, and since Lemma 3.10 implies $-v_{\varepsilon} + \theta(\alpha P U_{\varepsilon} + \alpha w_{\varepsilon} + \alpha^2 z_{\varepsilon}) \leq -cd(\cdot, \partial\Omega)$ in a neighborhood of $\partial\Omega$, we get

$$|f'''(-v_{\varepsilon} + \theta(\alpha P U_{\delta,\mu} - \alpha w_{\varepsilon} - \alpha^2 z_{\varepsilon}))| = O(1 + d(\cdot, \partial \Omega)^{\varepsilon - 1}).$$

Since $G_{\xi} = O(d(\cdot, \partial \Omega))$ near $\partial \Omega$, we deduce that

$$f_{\varepsilon}(\omega_{\varepsilon}) = -f_{\varepsilon}(v_{\varepsilon}) + \alpha f_{\varepsilon}'(v_{\varepsilon})(8\pi G_{\xi} - w_{\varepsilon}) + \alpha^2 \left(\frac{1}{2}f''(-v_{\varepsilon})(8\pi G_{\xi} - w_{\varepsilon})^2 - f_{\varepsilon}'(v_{\varepsilon})z_{\varepsilon}\right) + O(\alpha^3(1 + G_{\xi}^3)).$$

Since by construction we have $\Delta \omega_{\varepsilon} = -\alpha e^{U_{\varepsilon}} - \Delta v_{\varepsilon} - \alpha \Delta w_{\varepsilon} - \alpha^2 \Delta z_{\varepsilon}$, with v_{ε} , w_{ε} , z_{ε} solving (4) and (12)-(13), we conclude that

$$\begin{split} R &= -\alpha e^{U_{\varepsilon}} + O(\alpha^3 (1 + G_{\xi}^3)) \\ &= O(\delta^2 \rho_2^{-4}) + O(\alpha^3 (1 + G_{\xi}^3)) \\ &= O(\alpha^3 (1 + G_{\xi}^3)). \end{split}$$

3.6 The final estimate of the error in a mixed norm

We can summarize the estimates of the previous sections as follows: In $B(\xi, \rho_0)$, Corollary 3.5 gives $|R| \leq \alpha^3 j_{\varepsilon}$, where

$$j_{\varepsilon}(x) := e^{U_{\varepsilon}(x)} \left(1 + |\bar{U}\left(\frac{x-\xi}{\delta}\right)|^4 \right).$$
(54)

In $B(\xi, \rho_1) \setminus B(\xi, \rho_0)$, Lemma 3.8 shows that the norm of R in $L^{1+\alpha^2}$ is exponentially small in α .

Finally, in $\Omega \setminus B(\xi, \rho_1)$, Lemma 3.9 and Lemma 3.12 give L^2 estimates on R. This suggests to introduce the norm

$$\|f\|_{\varepsilon} := \|j_{\varepsilon}^{-1}f\|_{L^{\infty}(B(\xi,\rho_0))} + \frac{1}{\alpha^2} \|f\|_{L^{1+\alpha^2}(B(\xi,\rho_1)\setminus B(\xi,\rho_0))} + \|f\|_{L^2(\Omega\setminus B(\xi,\rho_1))}.$$
 (55)

The coefficient $\frac{1}{\alpha^2}$ is chosen in order to match the norm of $(-\Delta)^{-1}$ as a linear operator from $L^{1+\alpha^2}(B(\xi,\rho_1) \setminus B(\xi,\rho_0))$ into $L^{\infty}(B(\xi,\rho_1) \setminus B(\xi,\rho_0))$ (see Corollary B.4).

According to the estimates above we have:

Proposition 3.13 There exists $D_1 > 0$, $\varepsilon_0 > 0$ such that

$$||R||_{\varepsilon} \le D_1 \alpha^3$$

for any $\varepsilon \in (0, \varepsilon_0)$, $\mu \in \mathcal{U}$, $\xi \in B(\xi_0, \sigma)$.

We conclude this section by stating some simple properties of the norm $\|\cdot\|_{\varepsilon}$ and the weight j_{ε} .

Lemma 3.14 There exists a constant C > 0 such that

$$\|\cdot\|_{L^1(\Omega)} \le C \|\cdot\|_{\varepsilon}$$

for any $\varepsilon > 0$, $\mu \in \mathcal{U}$, $\xi \in B(\xi_0, \sigma)$.

Proof. Let $f: \Omega \to \mathbb{R}$ be a Lebesgue measurable function. Then

$$\|f\|_{L^1(B(\xi,\rho_0))} \le \|f\|_{\varepsilon} \int_{B(\xi,\rho_0)} j_{\varepsilon} dx = \|f\|_{\varepsilon} \int_{B(0,\frac{\rho_0}{\delta})} e^{\bar{U}} (1+\bar{U}^4) dy \le C \|f\|_{\varepsilon}.$$

By Hölder's inequality

$$\|f\|_{L^{1}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} \leq \|f\|_{L^{1+\alpha^{2}}(B(\xi,\rho_{1})\setminus B(\xi,\rho_{0}))} \rho_{1}^{\frac{2\alpha^{2}}{1+\alpha^{2}}} \leq C\|f\|_{\varepsilon},$$

and

$$\|f\|_{L^1(\Omega\setminus B(\xi,\rho_1))} \le \|f\|_{L^2(\Omega\setminus B(\xi,\rho_1))} |\Omega\setminus B(\xi,\rho_1)|^{\frac{1}{2}} \le C \|f\|_{\varepsilon}.$$

Hence, the conclusion follows.

Lemma 3.15 For any $\varepsilon > 0$ let ρ_{ε} , σ_{ε} be such that $\rho_2 \leq \sigma_{\varepsilon} \leq \sigma$ and $\delta \ll \rho_{\varepsilon} \leq \rho_0$ as $\varepsilon \to 0$. Let φ_{ε} of be the solution to

$$\begin{cases} -\Delta \varphi_{\varepsilon} = j_{\varepsilon} & \text{ in } B(\xi, \sigma_{\varepsilon}) \setminus B(\xi, \rho_{\varepsilon}), \\ \varphi_{\varepsilon} = 0 & \text{ on } \partial B(\xi, \sigma_{\varepsilon}) \setminus B(\xi, \rho_{\varepsilon}). \end{cases}$$

As $\varepsilon \to 0$, we have

$$\|\varphi_{\varepsilon}\|_{L^{\infty}(B(\xi,\sigma_{\varepsilon})\setminus B(\xi,\rho_{\varepsilon}))} = o(1)$$

Proof. Let us first note that there exists a constant c > 0, such that

$$\delta^2 j_{\varepsilon}(\xi + \delta \cdot) = e^{\bar{U}}(1 + \bar{U}^4) = \frac{8\mu^2}{(\mu^2 + |\cdot|^2)^2} \left(1 + \log^4\left(\frac{8\mu^2}{(\mu^2 + |\cdot|^2)^2}\right)\right) \le c \frac{\mu}{(\mu^2 + |\cdot|^2)^{\frac{3}{2}}}$$

in \mathbb{R}^2 . Then, by the maximum principle, we have

$$|\varphi_{\varepsilon}| \le c\psi\left(\frac{\cdot - \xi}{\delta}\right) \quad \text{in } B(\xi, \sigma_{\varepsilon}) \setminus B(\xi, \rho_{\varepsilon}), \tag{56}$$

where ψ satisfies

$$\begin{cases} -\Delta \psi = \frac{\mu}{(\mu^2 + |\cdot|^2)^{\frac{3}{2}}} & \text{in } A_{\varepsilon} := B(0, \frac{\sigma_{\varepsilon}}{\delta}) \setminus B(0, \frac{\rho_{\varepsilon}}{\delta}), \\ \psi = 0 & \text{on } \partial A_{\varepsilon}. \end{cases}$$

Since the function $W := -\log(\mu + \sqrt{|\cdot|^2 + \mu^2})$ satisfies $-\Delta W = \frac{\mu}{(\mu^2 + |\cdot|^2)^{\frac{3}{2}}}$, we have $\psi = a + b \log |\cdot| + W$,

for suitable constants $a, b \in \mathbb{R}$. Denoting $R_1 = \frac{\rho_{\varepsilon}}{\delta}$ and $R_2 = \frac{\sigma_{\varepsilon}}{\delta}$ one can verify that

$$a = \frac{W(R_2)\log R_1 - W(R_1)\log R_2}{\log R_2 - \log R_1} \quad \text{and} \quad b = \frac{W(R_1) - W(R_2)}{\log R_2 - \log R_1}.$$

Since

$$|W + \log|\cdot|| \le \frac{C\mu}{|\cdot|} = O\left(\frac{1}{R_1}\right),$$

uniformly in $\overline{A_{\varepsilon}}$, one has $a = O\left(\frac{\log R_2}{R_1(\log R_2 - \log R_1)}\right)$ and $b = 1 + O\left(\frac{1}{R_1(\log R_2 - \log R_1)}\right)$. Then

$$\psi = a + (b - 1) \log |\cdot| + O(\frac{1}{R_1})$$

= $O\left(\frac{1}{R_1} \frac{\log R_2}{\log R_2 - \log R_1}\right) + O\left(\frac{1}{R_1}\right)$
= $O\left(\frac{1}{R_1} \frac{1}{1 - \frac{\log R_1}{\log R_2}}\right) + O\left(\frac{1}{R_1}\right).$

Since

$$\frac{\log R_1}{\log R_2} = \frac{\log \frac{\rho_{\varepsilon}}{\delta}}{\log \sigma_{\varepsilon} - \log \delta} \le \frac{\log \frac{\rho_0}{\delta}}{\log \rho_2 - \log \delta} = O(\alpha),$$

we conclude that $\psi_{\mu} = O(\frac{1}{R_1}) = o(1)$, uniformly in A_{ε} . Then, the conclusion follows by (56).

4 The Linear Theory

Let us consider the linear operator

$$L\varphi = \varphi - (-\Delta)^{-1} (\lambda f_{\varepsilon}'(\omega_{\varepsilon})\varphi)$$

introduced in (19). In this section we give a priori estimates for the operator L and we prove its invertibility on a suitable subspace of $H_0^1(\Omega)$.

Lemma 4.1 The following expansions hold:

- 1. $\lambda f_{\varepsilon}'(\omega_{\varepsilon}) = e^{U_{\varepsilon}}(1 + O(\varepsilon^2))$ in $B(\xi, \rho_0)$. 2. $\lambda f_{\varepsilon}'(\omega_{\varepsilon}) = O(\Gamma_{\varepsilon})$ in $B(\xi, \rho_1)$, with Γ_{ε} as in (48). 3. $\lambda f_{\varepsilon}'(\omega_{\varepsilon}) = O(1)$ in $\Omega \setminus B(\xi, \rho_1)$.
- $4. \ \|\lambda f_{\varepsilon}'(\omega_{\varepsilon})\chi_{B(\xi,\rho_1)} e^{U_{\varepsilon}}\|_{\varepsilon} = o(1) \ \text{as} \ \varepsilon \to 0.$

Proof. For $x \in B(\xi, \rho_0)$, using (28)-(32), Lemma 3.3, (34), and (45), we have that

$$\lambda f_{\varepsilon}'(\omega_{\varepsilon}) = \lambda (1 + 2\omega_{\varepsilon}^{2} + (1 + \varepsilon)\omega_{\varepsilon}^{1 + \varepsilon})e^{\omega_{\varepsilon}^{2} + \omega_{\varepsilon}^{1 + \varepsilon}}$$
$$= \lambda \beta^{2} (2 + O(\alpha))e^{\beta^{2} + \beta^{1 + \varepsilon} + \bar{U}(\frac{\cdot - \xi}{\delta}) + O(\varepsilon^{2})}$$
$$= e^{U_{\varepsilon}} (1 + O(\varepsilon^{2})).$$

For $x \in B(\xi, \rho_1)$, using Remark 2.7, Lemma 3.3 we have

$$\begin{split} \lambda f_{\varepsilon}'(\omega_{\varepsilon}) &= \lambda (1 + 2\omega_{\varepsilon}^{2} + (1 + \varepsilon)\omega_{\varepsilon}^{1 + \varepsilon})e^{\omega_{\varepsilon}^{2} + \omega_{\varepsilon}^{1 + \varepsilon}} \\ &= \lambda \beta^{2} (2 + O(\alpha))e^{\beta^{2} + \beta^{1 + \varepsilon} + \bar{U}(\frac{\cdot}{\delta}) + \bar{U}(\frac{\cdot - \xi}{\delta})^{2} (1 + O(\varepsilon\alpha))} \\ &= O\left(\Gamma_{\varepsilon}\right). \end{split}$$

Claim 3 follows directly from Lemma 3.9. Finally, claim 4 follows by claims 1 and 2, using also Lemma 3.7 and the estimates

$$||e^{U_{\varepsilon}}||_{L^{1+\alpha}(B(\xi,\rho_1)\setminus B(\xi,\rho_0))} = o(1), ||e^{U_{\varepsilon}}||_{L^2(\Omega\setminus B(\xi,\rho_1))} = o(1).$$

According to Lemma 4.1, for $|x - \xi| \leq \rho_0$, L approaches the operator $L_0 \varphi := \varphi - (-\Delta)^{-1} (e^{U_{\varepsilon}} \varphi)$. Note that

$$\begin{split} L_0 \varphi &= 0 \quad \text{in } \Omega &\iff \quad -\Delta \varphi = e^{U_{\varepsilon}} \varphi \quad \text{in } \Omega \\ \iff \quad -\Delta \Phi = e^{\bar{U}} \Phi \quad \text{in } \frac{\Omega - \xi}{\delta}, \text{ where } \Phi = \varphi(\xi + \delta \cdot). \end{split}$$

Let us recall the following known fact about L_0 (see for example [10]).

Proposition 4.2 All bounded weak solutions of the problem

$$-\Delta \Phi = e^{\bar{U}} \Phi \quad in \ \mathbb{R}^2 \tag{57}$$

have the form

$$\Phi = c_0 Z_0 + c_1 Z_1 + c_2 Z_2,$$

where $c_0, c_1, c_2 \in \mathbb{R}$ and

$$Z_0(y) := \frac{\mu^2 - |y|^2}{\mu^2 + |y|^2}, \qquad Z_1(y) := \frac{2\mu y_1}{\mu^2 + |y|^2}, \qquad Z_2(y) := \frac{2\mu y_2}{\mu^2 + |y|^2}$$

Remark 4.3 The functions Z_0, Z_1, Z_2 are orthogonal in $D^{1,2}(\mathbb{R}^2)$, that is

$$\int_{\mathbb{R}^2} \nabla Z_i \cdot \nabla Z_j dy = \int_{\mathbb{R}^2} e^{\bar{U}} Z_i Z_j dy = \frac{8}{3} \pi \delta_{i,j}.$$
(58)

In the following we denote

$$Z_{i,\varepsilon}(x) := Z_i\left(\frac{x-\xi}{\delta}\right)$$
 and $PZ_{i,\varepsilon} = (-\Delta)^{-1}Z_{i,\varepsilon}, \quad i = 0, 1, 2$

Lemma 4.4 It holds true that

$$PZ_{0,\varepsilon} = Z_{0,\varepsilon} + 1 + O(\delta^2)$$
 and $PZ_{i,\varepsilon} = Z_{i,\varepsilon} + O(\delta), i = 1, 2,$

uniformly with respect to $\mu \in \mathcal{U}, \xi \in B(\xi_0, \sigma)$.

Proof. See for example Appendix A in [18].

Lemma 4.4 shows the smallness of $PZ_{i,\varepsilon} - Z_{i,\varepsilon}$ for i = 1, 2, but not for i = 0. For this reason, in many cases it is convenient to replace $PZ_{0,\varepsilon}$ with the function

$$\widetilde{Z}_{\varepsilon} := \begin{cases} Z_{0,\varepsilon} & \text{if } |x - \xi| \le \rho_0, \\ Z_{0,\varepsilon}(\rho_0)(\frac{\log \rho_1 - \log |x - \xi|}{\log \rho_1 - \log \rho_0}) & \text{if } \rho_0 \le |x - \xi| \le \rho_1, \\ 0 & \text{if } |x - \xi| \ge \rho_1. \end{cases}$$
(59)

Lemma 4.5 The function $\widetilde{Z}_{\varepsilon}$ satisfies the following properties:

- $\widetilde{Z}_{\varepsilon} \in H^1_0(\Omega)$ and $|\widetilde{Z}_{\varepsilon}| \leq 1$ in Ω .
- $\|\nabla(\widetilde{Z}_{\varepsilon}-Z_{0,\varepsilon})\|_{L^{2}(\Omega)} \to 0$, uniformly for $\mu \in \mathcal{U}$ and $\xi \in B(\xi_{0},\sigma)$.

Proof. The first property follows trivially from the definition. Moreover we have

$$\begin{split} \|\nabla(\widetilde{Z}_{\varepsilon} - Z_{0,\varepsilon})\|_{L^{2}(\Omega)}^{2} &\leq \frac{Z_{0,\varepsilon}(\rho_{0})^{2}}{(\log\rho_{1} - \log\rho_{0})^{2}} \int_{B(\xi,\rho_{1})\setminus B(\xi,\rho_{0})} \frac{1}{|x - \xi|^{2}} dx + \|\nabla Z_{0,\varepsilon}\|_{L^{2}(\Omega\setminus B(\xi,\rho_{0}))}^{2} \\ &\leq \frac{2\pi Z_{0,\varepsilon}(\rho_{0})^{2}}{\log\rho_{1} - \log\rho_{0}} + \|\nabla Z_{0}\|_{L^{2}(\mathbb{R}^{2}\setminus B(0,\frac{\rho_{0}}{\delta}))}^{2} \\ &= O(\alpha^{2}) + O(e^{-\frac{\varepsilon}{\alpha}}) \to 0, \end{split}$$

as $\varepsilon \to 0$.

We will denote by K_{ε} the subspace of $H_0^1(\Omega)$ spanned by $PZ_{i,\varepsilon}$, i = 0, 1, 2 and by K_{ε}^{\perp} the subspaces of $H_0^1(\Omega)$ orthogonal to K_{ε} , i.e.

$$K_{\varepsilon}^{\perp} = \left\{ u \in H_0^1(\Omega) : \int_{\Omega} \nabla P Z_{i,\varepsilon} \cdot \nabla u \, dx = \int_{\Omega} e^{U_{\varepsilon}} Z_{i,\varepsilon} u \, dx = 0, \ i = 0, 1, 2 \right\}.$$

Let π and π^{\perp} be the projections of $H_0^1(\Omega)$ respectively on K_{ε} and K_{ε}^{\perp} . Finally, we denote

$$Y_{\varepsilon} := \{ f \in L^1(\Omega) : \|f\|_{\varepsilon} < +\infty \}.$$

Proposition 4.6 There exist $\varepsilon_0 > 0$ and a constant $D_0 > 0$ such that

$$\|\varphi\|_{H^1_0(\Omega)} + \|\varphi\|_{L^\infty(\Omega)} \le D_0 \|h\|_{\varepsilon},\tag{60}$$

for any $\varepsilon \in (0, \varepsilon_0)$, $\mu \in \mathcal{U}$, $\xi \in B(\xi_0, \sigma)$, $h \in Y_{\varepsilon}$ and $\varphi \in K_{\varepsilon}^{\perp}$ satisfying

$$\pi^{\perp} \left\{ L \varphi - (-\Delta)^{-1} h \right\} = 0.$$
(61)

Proof. We assume by contradiction that there exists $\varepsilon_n \to 0$, $\mu_n \in \mathcal{U}$, $\xi_n \in B(\xi_0, \sigma)$, $h_n \in Y_{\varepsilon}$ and a solution $\varphi_n \in K_{\varepsilon_n}^{\perp}$ of (61) such that

$$\frac{\|\varphi_n\|_{H^1_0(\Omega)} + \|\varphi_n\|_{L^{\infty}(\Omega)}}{\|h_n\|_{\varepsilon_n}} \to +\infty.$$

Let $\delta_n, \alpha_n, \beta_n$ be the parameters in Lemma 2.4 corresponding to ε_n, μ_n and ξ_n . Let also $\rho_{0,n}, \rho_{1,n}, \rho_{2,n}$ be defined as in (39). We denote $\omega_n := \omega_{\varepsilon_n}, U_n := U_{\varepsilon_n}, Z_{i,n} := Z_{i,\varepsilon_n}$ and $f_n := f_{\varepsilon_n}$. W.l.o.g we can assume that $\|\varphi_n\|_{H_0^1(\Omega)} + \|\varphi_n\|_{L^{\infty}(\Omega)} = 1$ and $\|h_n\|_{\varepsilon_n} \to 0$. Since φ_n satisfies (61), there exist $c_{i,n} \in \mathbb{R}, i = 0, 1, 2$, such that

$$-\Delta\varphi_n - \lambda f'_n(\omega_n)\varphi_n = h_n + \sum_{i=0}^2 c_{i,n} e^{U_n} Z_{i,n}.$$
(62)

Step 1 We have $c_{i,n} \to 0$ as $n \to +\infty$, i = 0, 1, 2.

Let $\widetilde{Z}_n := \widetilde{Z}_{\varepsilon_n}$ be the function defined in (59). Testing equation (62) against \widetilde{Z}_n , we get

$$\sum_{j=0}^{2} c_{j,n} \int_{\Omega} e^{U_n} Z_{j,n} \widetilde{Z}_n dx = \int_{\Omega} \nabla \widetilde{Z}_n \cdot \nabla \varphi_n dx - \int_{\Omega} \lambda f'_n(\omega_n) \varphi_n \widetilde{Z}_n dx - \int_{\Omega} h_n \widetilde{Z}_n dx.$$
(63)

Since $\|\varphi_n\|_{H^1_0(\Omega)} \leq 1$ and $\varphi_n \in K^{\perp}_{\varepsilon_n}$, using Lemma 4.5 we get

$$\int_{\Omega} \nabla \widetilde{Z}_n \cdot \nabla \varphi_n dx = \int_{\Omega} \nabla Z_{0,n} \cdot \nabla \varphi_n dx + o(1) = \underbrace{\int_{\Omega} e^{U_n} Z_{0,n} \varphi_n dx}_{=0} + o(1) = o(1),$$

as $n \to +\infty$. By Lemma 4.1 and Lemma 3.7, we find

$$\int_{\Omega} \lambda f_n'(\omega_n) \varphi_n \widetilde{Z}_n dx = \int_{B(\xi_n, \rho_{0,n})} e^{U_n} \varphi_n Z_{0,n} dx + O(\varepsilon_n^2) + O\left(\|\Gamma_{\varepsilon}\|_{L^1(B(\xi_n, \rho_{1,n}) \setminus B(\xi_n, \rho_{0,n}))} \right)$$
$$= \underbrace{\int_{\Omega} e^{U_n} \varphi_n Z_{0,n} dx}_{=0} + o(1) = o(1).$$

Finally, Lemma 4.5 and Lemma 3.14 give

$$\left|\int_{\Omega} h_n \widetilde{Z}_n dx\right| \le \|h_n\|_{L^1(\Omega)} \le C \|h_n\|_{\varepsilon_n} = o(1).$$

Then (63) rewrites as

$$\sum_{j=0}^{2} c_{j,n} \int_{\Omega} e^{U_n} Z_{j,n} \widetilde{Z}_n dx = o(1).$$
(64)

With similar arguments, testing equation (62) against $PZ_{i,n}$ for i = 1, 2, we get that

$$\sum_{j=0}^{2} c_{j,n} \int_{\Omega} e^{U_n} Z_{j,n} P Z_{i,n} dx = -\int_{\Omega} \lambda f'_n(\omega_n) \varphi_n P Z_{i,n} dx - \int_{\Omega} h_n P Z_{i,n} dx$$

$$= \underbrace{\int_{\Omega} e^{U_n} \varphi_n Z_{i,n} dx}_{=0} + o(1) = o(1).$$
(65)

Note that, as in (58), we have

$$\begin{split} \int_{\Omega} e^{U_n} Z_{j,n} \widetilde{Z}_n dx &= \int_{B(\xi_n,\rho_{0,n})} e^{U_n} Z_{j,n} Z_{0,n} dx + O\left(\int_{\mathbb{R}^2 \setminus B(\xi_n,\rho_{0,n})} e^{U_n}\right) \\ &= \int_{B(0,\frac{\rho_{0,n}}{\delta_n})} e^{\overline{U}} Z_j Z_0 dy + o(1) \\ &= \frac{8}{3} \pi \delta_{0j} + o(1), \end{split}$$

for j = 0, 1, 2. Similarly

$$\int_{\Omega} e^{U_n} Z_{j,n} P Z_{i,n} dx = \int_{\Omega} e^{U_n} Z_{j,n} Z_{i,n} dx + o(1)$$
$$= \frac{8}{3} \pi \delta_{ij} + o(1),$$

for i = 1, 2, j = 0, 1, 2. Then, (63) and (64) rewrite as

$$\sum_{j=0}^{2} c_{j,n}(\delta_{ij} + o(1)) = o(1),$$

which implies the conclusion.

Step 2 If
$$\widetilde{h}_n := h_n + \left(\lambda f'_n(\omega_n)\chi_{B(\xi_n,\rho_{1,n})} - e^{U_n}\right)\varphi_n + \sum_{j=0}^2 c_{j,n}e^{U_n}Z_{j,n}$$
, then
 $-\Delta\varphi_n = e^{U_n}\varphi_n + \lambda f'_n(\omega_n)\chi_{\Omega\setminus B(\xi_n,\rho_{1,n})}\varphi_n + \widetilde{h}_n \quad in \ \Omega, \quad and \quad \|\widetilde{h}_n\|_{\varepsilon_n} \to 0.$ (66)

Since $||h_n||_{\varepsilon_n} \to 0$, $|Z_{i,n}| \le 1$, and $||\lambda f'_n(\omega_n)\chi_{B(\xi_n,\rho_{1,n})} - e^{U_n}||_{\varepsilon_n} \to 0$ by Lemma 4.1, it is sufficient to observe that $||e^{U_n}||_{\varepsilon_n} = O(1)$ and apply Step 1.

Step 3 There exists $\delta_n \ll \rho_n \leq \rho_{0,n}$ such that, up to a subsequence, $\|\varphi_n\|_{L^{\infty}(B(\xi_n,\rho_n))} \to 0$ as $n \to +\infty$.

Let us consider the sequence $\Phi_n(y) := \varphi_n(\xi_n + \delta_n y), y \in \frac{\Omega - \xi_n}{\delta_n}$. By (66) Φ_n satisfies

$$-\Delta \Phi_n = e^{\bar{U}} \Phi_n + \delta_n^2 \tilde{h}_n(\xi + \delta_n \cdot) \qquad \text{in } B\left(0, \frac{\rho_{1,n}}{\delta_n}\right).$$

We know that

$$\left|e^{\bar{U}(y)}\Phi_n(y)\right| \le e^{\bar{U}(y)} \le \frac{8}{\mu^2},$$

and, for $y \in B(0, \frac{\rho_{0,n}}{\delta_n})$, that

$$\delta_n^2 |\tilde{h}_n(\xi + \delta_n y)| \le \delta_n^2 j_{\varepsilon_n}(\xi + \delta_n y) \|\tilde{h}_n\|_{\varepsilon_n} = e^{\bar{U}(y)} (1 + |\bar{U}(y)|^4) \|\tilde{h}_n\|_{\varepsilon_n} \le C \|\tilde{h}_n\|_{\varepsilon_n} \to 0.$$

In particular Φ_n and $\Delta \Phi_n$ are uniformly bounded in $B(0, \frac{\rho_{0,n}}{\delta_n})$. By standard elliptic estimates, we can find $\Phi_0 \in C(\mathbb{R}^2) \cap H^1_{loc}(\mathbb{R}^2)$ and a sequence $R_n \to +\infty$, $R_n \leq \frac{\rho_{0,n}}{\delta_n}$, such that, up to a subsequence, $\|\Phi_n - \Phi_0\|_{L^{\infty}(B(0,R_n))} \to 0$. Moreover, $|\Phi_0| \leq 1$ and Φ_0 is a weak solution to

$$-\Delta \Phi_0 = e^{\bar{U}} \Phi_0 \quad \text{in } \mathbb{R}^2$$

According to Proposition 4.2, we must have $\Phi_0 = \kappa_0 Z_0 + \kappa_1 Z_1 + \kappa_2 Z_2$, for some $\kappa_i \in \mathbb{R}$, i = 0, 1, 2. Keeping in mind (58) and using that $e^{\bar{U}} \in L^1(\mathbb{R}^2)$, we obtain

$$\begin{aligned} 0 &= \int_{\Omega} e^{U_n} Z_{i,n} \phi_n \, dx = \int_{\frac{\Omega - \xi_n}{\delta_n}} e^{\bar{U}} Z_i \Phi_n \, dy \\ &= \int_{B(0,R_n)} e^{\bar{U}} Z_i \Phi_n \, dy + O\left(\int_{\mathbb{R}^2 \setminus B(0,R_n)} e^{\bar{U}} dy\right) \\ &\to \frac{8}{3} \pi \kappa_i, \end{aligned}$$

for i = 0, 1, 2. This implies $\kappa_i = 0, i = 0, 1, 2$. Then $\Phi_0 \equiv 0$ and we get the conclusion with $\rho_n = \delta_n R_n$.

Step 4 Up to a subsequence, $\xi_n \to \overline{\xi} \in \Omega$ and $\varphi_n \to 0$ in $L^{\infty}_{loc}(\Omega \setminus {\{\overline{\xi}\}})$, as $n \to \infty$.

We know that φ_n satisfies (66) in Ω . Since $|\varphi_n| \leq 1$, $||e^{U_n}||_{L^{\infty}(\Omega \setminus B(\xi_n, \rho_{1,n}))} \to 0$, $||h_n||_{L^2(\Omega \setminus B(\xi, \rho_{1,n}))} \to 0$, and $||f'_n(\omega_n)||_{L^{\infty}(\Omega \setminus B(\xi, \rho_{1,n}))} = O(1)$, by ellpitic estimates we find that φ_n is bounded in $C^{0,\gamma}_{loc}(\overline{\Omega} \setminus \{\overline{\xi}\})$, for some $\gamma \in (0, 1)$. Therefore, there exists $\varphi_0 \in C(\overline{\Omega}) \cap H^1_0(\Omega)$, such that $\varphi_n \to \varphi_0$ locally uniformly on $\overline{\Omega} \setminus \{\overline{\xi}\}$ and weakly in $H_0^1(\Omega)$. Noting that $\omega_n \to -u_0$ locally uniformly in $\overline{\Omega} \setminus \{\xi\}$ and that f'_n is even, we see that φ_0 satisfies $\Delta \varphi_0 + f'_0(u_0)\varphi_0$ in $\Omega \setminus \{\bar{\xi}\}$. Actually, since $\varphi_0, \Delta \varphi_0 \in L^{\infty}(\Omega), \varphi_0$ is a weak solution of $\Delta \varphi_0 + f'_0(u_0)\varphi_0 = 0$ in Ω . Then, the non-degeneracy of u_0 implies $\varphi_0 \equiv 0$.

Step 5 $\|\varphi_n\|_{L^{\infty}(\Omega)} \to 0.$

By Step 4, we can find a sequence $\sigma_n \geq \rho_{2,n}$ such that $\|\varphi_n\|_{L^{\infty}(\Omega \setminus B(\xi_n, \sigma_n))} \to 0$ as $n \to +\infty$, up to a subsequence. Then, it is sufficient to show that $\|\varphi_n\|_{L^{\infty}(A_n)} \to 0$, where $A_n := B(\xi_n, \sigma_n) \setminus B(\xi_n, \rho_n)$ and ρ_n is as in Step 3. We can split $\varphi_n = \varphi_n^{(0)} + \varphi_n^{(1)} + \varphi_n^{(2)} + \varphi_n^{(3)}$, where

$$\begin{cases} \Delta \varphi_n^{(0)} = 0 & \text{in } A_n, \\ \varphi_n^{(0)} = \varphi_n & \text{on } \partial A_n, \end{cases} \quad \text{and} \quad \begin{cases} -\Delta \varphi_n^{(i)} = f_{i,n} & \text{in } A_n, \\ \varphi_n^{(i)} = 0 & \text{on } \partial A_n, \end{cases} \quad \text{for } i = 1, 2, 3,$$

with

$$\begin{cases} f_{1,n} := e^{U_n} \varphi_n + \widetilde{h}_n \chi_{B(\xi_n,\rho_{0,n})}, \\ f_{2,n} := \widetilde{h}_n \chi_{B(\xi_n,\rho_{1,n}) \setminus B(\xi_n,\rho_{0,n})}, \\ f_{3,n} := \widetilde{h}_n \chi_{B(\xi_n,\sigma_n) \setminus B(\xi_n,\rho_{1,n})} + \lambda f'_n(\omega_n) \chi_{B(\xi_n,\sigma_n) \setminus B(\xi_n,\rho_{1,n})} \varphi_n. \end{cases}$$

By the maximum principle

$$\|\varphi_n^{(0)}\|_{L^{\infty}(A_n)} \le \|\varphi_n\|_{L^{\infty}(\partial A_n)} \to 0.$$

Since

$$|f_{1,n}| \le e^{U_n} + \|\widetilde{h}_n\|_{\varepsilon_n} j_{\varepsilon_n} \le j_{\varepsilon_n} (1+o(1)) \le 2j_{\varepsilon_n}$$

we get that $|\varphi_n^{(1)}| \leq 2\psi_n$, where ψ_n satisfies

$$\begin{cases} -\Delta \psi_n = j_{\varepsilon_n} & \text{in } A_n \\ \psi_n = 0 & \text{on } \partial A_n \end{cases}$$

Lemma 3.15 implies $\|\psi_n\|_{L^{\infty}(A_n)} \to 0$, hence $\|\varphi_n^{(1)}\|_{L^{\infty}(A_n)} \to 0$. Finally, since $|A_n|$ is uniformly bounded, elliptic estimates (see Corollaries B.3 and B.4) give

$$\|\varphi_n^{(2)}\|_{L^{\infty}(A_n)} \le \frac{C}{\alpha^2} \|f_{2,n}\|_{L^{1+\alpha^2}(A_n)} = \frac{C}{\alpha^2} \|\widetilde{h}_n\|_{L^{1+\alpha^2}(B(\xi_n,\rho_{1,n})\setminus B(\xi_n,\rho_{0,n}))} \le \|\widetilde{h}_n\|_{\varepsilon_n} \to 0,$$

and

$$\|\varphi_n^{(3)}\|_{L^{\infty}(A_n)} \le C \|f_{3,n}\|_{L^2(A_n)} = O(\|h_n\|_{\varepsilon_n}) + O(\sqrt{\sigma_n}) \to 0.$$

Step 6 Conclusion of the proof.

By Step 5, we have that $\|\varphi_n\|_{H^1_0(\Omega)} = 1 - \|\varphi_n\|_{L^{\infty}(\Omega)} \to 1$. But (66) gives

$$\begin{aligned} \|\varphi_n\|_{H^1_0(\Omega)}^2 &= \int_{\Omega} e^{U_n} \varphi_n^2 \, dx + \int_{\Omega \setminus B(\xi,\rho_{1,n})} \lambda f'_n(\omega_n) \varphi_n^2 \, dx + \int_{\Omega} \widetilde{h}_n \varphi_n \, dx \\ &= O(\|\varphi_n\|_{L^{\infty}(\Omega)}^2) + o(\|\varphi_n\|_{L^2(\Omega)}) \to 0. \end{aligned}$$

Then, we get a contadiction.

As a consequence we have that $\pi^{\perp}L$ is invertible on $K_{\varepsilon}^{\perp}.$

Corollary 4.7 $\pi^{\perp}L: K_{\varepsilon}^{\perp} \mapsto K_{\varepsilon}^{\perp}$ is invertible.

Proof. This follows by standard Fredholm theory. Indeed, for any $\varepsilon > 0$ the map $F(\varphi) := \pi^{\perp}(-\Delta)^{-1}(f'(\omega_{\varepsilon})\varphi)$ defines a compact operator on K_{ε}^{\perp} (in fact on $H_0^1(\Omega)$). Then $\pi^{\perp}L = Id_{K_{\varepsilon}^{\perp}} - F$ is a Fredholm operator of index 0. Proposition 4.6 implies that $\pi^{\perp}L$ is injective, hence it is invertible on K_{ε}^{\perp} .

5 The reduction to a finite dimensional problem

This section is devoted to reduce the problem to a finite dimensional one. More precisely, we prove:

Proposition 5.1 There exist $\varepsilon_0 > 0$ and a map $(\varepsilon, \mu, \xi) \to \varphi_{\varepsilon,\mu,\xi} \in K_{\varepsilon}^{\perp} \cap L^{\infty}(\Omega)$ defined in $(0, \varepsilon_0) \times \mathcal{U} \times B(\xi_0, \sigma)$ and continuous with respect to μ and ξ , such that for some D > 0

$$\|\varphi_{\varepsilon,\mu,\xi}\|_{H^1_0} + \|\varphi_{\varepsilon,\mu,\xi}\|_{L^{\infty}} \le D\alpha^3,\tag{67}$$

and

$$\pi^{\perp} \Big\{ L\varphi_{\varepsilon,\mu,\xi} - (-\Delta)^{-1} (R + N(\varphi_{\varepsilon,\mu,\xi})) \Big\} = 0,$$
(68)

where the linear operator L is defined in (19), the error term R is defined in (17) and the quadratic term N is defined in (18).

5.1 Estimates on $N(\varphi)$

For a function $\varphi \in H_0^1(\Omega) \cap L^{\infty}(\Omega)$, let $N(\varphi)$ be defined as in (18), i.e.

$$N(\varphi) = N_{\varepsilon,\mu,\xi}(\varphi) := \lambda \left(f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi} + \varphi) - f_{\varepsilon}(\omega_{\varepsilon,\mu,\xi}) - f'_{\varepsilon}(\omega_{\varepsilon,\mu,\xi})\varphi \right).$$

Let us estimate $||N(\varphi)||_{\varepsilon}$, where $||\cdot||_{\varepsilon}$ is defined as in (55). Let us define

$$\mathcal{B}_{\alpha} := \{ \varphi \in L^{\infty}(\Omega) : \|\varphi\|_{L^{\infty}(\Omega)} \le \alpha \}.$$
(69)

Lemma 5.2 There exists $D_2 > 0$ such that

$$\|N(\varphi_1) - N(\varphi_2)\|_{\varepsilon} \le D_2 \alpha^{-1} \left(\|\varphi_1\|_{L^{\infty}(\Omega)} + \|\varphi_2\|_{L^{\infty}(\Omega)} \right) \|\varphi_1 - \varphi_2\|_{L^{\infty}(\Omega)},$$

for any $\varphi_1, \varphi_2 \in \mathcal{B}_{\alpha}$.

Proof. First, for any $x \in \Omega$ we can find $\theta_1 = \theta_1(x) \in [0, 1]$ such that

$$N(\varphi_2) - N(\varphi_1) = \lambda \left(f_{\varepsilon}(\omega_{\varepsilon} + \varphi_2) - f_{\varepsilon}(\omega_{\varepsilon} + \varphi_1) - f'_{\varepsilon}(\omega_{\varepsilon})(\varphi_2 - \varphi_1) \right) \\ = \lambda \left(f'_{\varepsilon}(\omega_{\varepsilon} + \theta_1\varphi_2 + (1 - \theta_1)\varphi_1)(\varphi_2 - \varphi_1) - f'_{\varepsilon}(\omega_{\varepsilon})(\varphi_2 - \varphi_1) \right) \\ = \lambda \left(f'_{\varepsilon}(\omega_{\varepsilon} + \varphi_3) - f'_{\varepsilon}(\omega_{\varepsilon}) \right) (\varphi_2 - \varphi_1),$$

where $\varphi_3 := \theta_1 \varphi_2 + (1 - \theta_1) \varphi_1$. Furthermore, there exists $\theta_2 = \theta_2(x)$ such that

$$f'_{\varepsilon}(\omega_{\varepsilon}+arphi_3)=f'_{\varepsilon}(\omega_{\varepsilon})+f''_{\varepsilon}(\omega_{\varepsilon}+ heta_2arphi_3)arphi_3.$$

Thus, we obtain

$$|N(\varphi_1) - N(\varphi_2)| = \lambda |f_{\varepsilon}''(\omega_{\varepsilon} + \theta_2 \varphi_3)||\varphi_3||\varphi_1 - \varphi_2| \leq \lambda |f_{\varepsilon}''(\omega_{\varepsilon} + \theta_2 \varphi_3)| \left(\|\varphi_1\|_{L^{\infty}(\Omega)} + \|\varphi_2\|_{L^{\infty}(\Omega)} \right) \|\varphi_1 - \varphi_2\|_{L^{\infty}(\Omega)}.$$
(70)

Then, in order to conclude the proof, we shall bound $\|f_{\varepsilon}''(\omega_{\varepsilon} + \theta_2 \varphi_3)\|_{\varepsilon}$. Note that, there exists a universal constant $C_0 > 0$ such that

$$|f_{\varepsilon}''(t)| \le C_0(1+|t|^3)e^{t^2+|t|^{1+\varepsilon}}, \quad \forall t \in \mathbb{R}.$$

By Remark 2.7 we have $\omega_{\varepsilon} = O(\beta) = O(\alpha^{-1})$. Since $|\varphi_3| \le |\varphi_1| + |\varphi_2| \le 2\alpha$, we get

$$(\omega_{\varepsilon} + \theta_2 \varphi_3)^2 \le \omega_{\varepsilon}^2 + 2|\omega_{\varepsilon}||\varphi_3| + \varphi_3^2 = \omega_{\varepsilon}^2 + O(1).$$
(71)

By convexity, we also have

$$|\omega_{\varepsilon} + \theta_2 \varphi_3|^3 \le (|\omega_{\varepsilon}| + |\varphi_3|)^3 \le 4(|\omega_{\varepsilon}|^3 + |\varphi_3|^3) \le 4(|\omega_{\varepsilon}|^3 + \alpha^3).$$
(72)

In $B(\xi, \rho_1)$ we have $\omega_{\varepsilon} \ge c_0$ by Lemma 3.2, so that

$$(\omega_{\varepsilon} + \theta_2 \varphi_3)^{1+\varepsilon} \le \omega_{\varepsilon}^{1+\varepsilon} \left(1 + \frac{\alpha}{c_0}\right)^{1+\varepsilon} = \omega_{\varepsilon}^{1+\varepsilon} + O(1).$$
(73)

Clearly (71)-(73) yield the existence of a constant $C_1 > 0$ such that

$$|f_{\varepsilon}''(\omega_{\varepsilon}+\theta_{2}\varphi_{3})| \leq C_{1}\alpha^{-2}\omega_{\varepsilon}e^{\omega_{\varepsilon}^{2}+|\omega_{\varepsilon}|^{1+\varepsilon}} = C_{1}\alpha^{-2}f_{\varepsilon}(\omega_{\varepsilon}),$$

in $B(\xi, \rho_1)$. Arguing as in Lemma 3.4 (see (46)) we get

$$\lambda |f_{\varepsilon}''(\omega_{\varepsilon} + \theta_2 \varphi)| \le C \alpha^{-1} j_{\varepsilon} \quad \text{in } B(\xi, \rho_0).$$
(74)

Lemma 3.6 and Lemma 3.7 yield

$$\lambda \| f_{\varepsilon}''(\omega_{\varepsilon} + \theta_2 \varphi) \|_{L^{1+\alpha^2}(B(\xi,\rho_1) \setminus B(\xi,\rho_0))} = O(\alpha^{-2} e^{-\frac{c_2}{\sqrt{\alpha}}}).$$
(75)

Finally, thanks to Lemma 3.9, we know that

$$\lambda f_{\varepsilon}''(\omega_{\varepsilon} + \theta_2 \varphi_3) = O(1) \quad \text{in } \Omega \setminus B(\xi, \rho_1).$$
(76)

Thanks to (74)-(76) we infer

$$\lambda \| f_{\varepsilon}''(\omega_{\varepsilon} + \theta_2 \varphi_3) \|_{\varepsilon} = O(\alpha^{-1})$$

and the conclusion follows from (70).

Remark 5.3 Applying Lemma 5.2 with $\varphi_2 = 0$, we obtain that

$$\|N(\varphi)\|_{\varepsilon} \le D_2 \alpha^{-1} \|\varphi\|_{L^{\infty}(\Omega)}^2,$$

for any $\varphi \in \mathcal{B}_{\alpha}$.

Remark 5.4 The proof of Proposition 5.2 and Lemma 3.9 also shows that $\|N(\varphi)\|_{L^{\infty}(\Omega \setminus B(\xi,\rho_1))} \leq D_3 \|\varphi\|_{L^{\infty}(\Omega)}^2,$

for any $\varphi \in \mathcal{B}_{\alpha}$.

5.2 Proof of Proposition 5.1: a fixed point argument

Let us consider the operator

$$\mathcal{T} = \mathcal{T}_{\varepsilon,\mu,\xi} := (\pi^{\perp}L)^{-1}\pi^{\perp} \left[(-\Delta)^{-1} \left(N(\varphi) + R \right) \right]$$
(77)

on the space $X := K_{\varepsilon}^{\perp} \cap L^{\infty}(\Omega)$, which is a Banach space with respect to the norm

$$\|\cdot\|_X = \|\cdot\|_{H^1_0(\Omega)} + \|\cdot\|_{L^{\infty}(\Omega)}.$$

Let D_1 and D_0 be the constants defined in Proposition 3.13 and Proposition 4.6. Let us set

$$E_{\varepsilon} := \{ \varphi \in X : \|\varphi\|_X \le D_0(D_1 + 1)\alpha^3 \}.$$

Proposition 5.1 is an immediate consequence of the following result.

Proposition 5.5 There exists $\varepsilon_0 > 0$ such that, for any $\varepsilon \in (0, \varepsilon_0)$, $\mu \in \mathcal{U}$, $\xi \in B(\xi_0, \sigma)$, \mathcal{T} has a fixed point $\varphi_{\varepsilon,\mu,\xi} \in E_{\varepsilon}$, which depends continuously on μ and ξ .

Proof. Since E_{ε} is a closed subspace of X and \mathcal{T} depends continuously on μ and ξ , it is sufficient to verify that

- 1. \mathcal{T} maps E_{ε} into itself.
- 2. \mathcal{T} is a contraction, i.e. $||\mathcal{T}(\varphi_1) \mathcal{T}(\varphi_2)||_{H_0^1(\Omega)} \leq \theta ||\varphi_1 \varphi_2||_{H_0^1(\Omega)}$ for some positive constant $\theta < 1$ and for any $\varphi_1, \varphi_2 \in E_{\varepsilon}$.

Then the conclusion follows by the contraction mapping theorem.

Step 1 \mathcal{T} maps E_{ε} into itself.

Let us denote $C_0 := D_0(D_1 + 1)$. Take $\varphi \in E_{\varepsilon}$ and set

$$h(\varphi) := R + N(\phi).$$

If ε is small enough, we have that $\alpha^2 C_0 \leq 1$, so that $E_{\varepsilon} \subseteq \mathcal{B}_{\alpha}$ (see (69)). By Proposition 3.13 and Remark 5.3 we get

$$\begin{split} \|h(\varphi)\|_{\varepsilon} &\leq \|R\|_{\varepsilon} + \|N(\varphi)\|_{\varepsilon} \\ &\leq D_1 \alpha^3 + D_2 \alpha^{-1} \|\varphi\|_{L^{\infty}(\Omega)}^2 \\ &\leq D_1 \alpha^3 + C_0^2 D_2 \alpha^5, \end{split}$$

for any $\varphi \in E_{\varepsilon}$. Then, if we take ε small enough so that $C_0^2 D_2 \alpha^2 \leq 1$, we get that

$$\|h(\varphi)\|_{\varepsilon} \le (D_1 + 1)\alpha^3.$$

Since by definition

$$\pi^{\perp}L(\mathcal{T}(\varphi)) = \pi^{\perp}(-\Delta)^{-1}h(\varphi),$$

we have by Proposition 4.6 that

$$\|\mathcal{T}(\varphi)\|_X \le D_0 \|h(\varphi)\|_{\varepsilon} \le D_0 (D_1 + 1)\alpha^3,$$

that is $\mathcal{T}(\varphi) \in E_{\varepsilon}$.

Step 2 \mathcal{T} is a contraction mapping in E_{ε} .

Let us take ε small enough so that $D_0 D_2 C_0 \alpha^2 \leq \frac{1}{4}$ and $E_{\varepsilon} \subseteq \mathcal{B}_{\alpha}$. By Propositions 4.6 and 5.2 we have

$$\begin{split} \|T(\varphi_1) - T(\varphi_2)\|_X &\leq D_0 \|h(\varphi_1) - h(\varphi_2)\|_{\varepsilon} \\ &= D_0 \|N(\varphi_1) - N(\varphi_2)\|_{\varepsilon} \\ &\leq D_0 D_2 \alpha^{-1} (\|\varphi_1\|_{L^{\infty}(\Omega)} + \|\varphi_2\|_{L^{\infty}(\Omega)}) \|\varphi_1 - \varphi_2\|_{L^{\infty}(\Omega)} \\ &\leq 2C_0 D_0 D_2 \alpha^2 \|\varphi_1 - \varphi_2\|_{L^{\infty}(\Omega)} \\ &\leq \frac{1}{2} \|\varphi_1 - \varphi_2\|_{L^{\infty}(\Omega)}, \end{split}$$

for any $\varphi_1, \varphi_2 \in E_{\varepsilon}$. Then, \mathcal{T} is a contraction mapping on E_{ε} .

6 The reduced problem: proof of Theorem 1.3 completed

Let $\varphi_{\varepsilon} := \varphi_{\varepsilon,\mu,\xi}$ be as in Proposition 5.1. By (68), we can find $\kappa_{\varepsilon,i} = \kappa_{\varepsilon,i}(\mu,\xi)$, i = 0, 1, 2(which depend continuously on μ , and ξ), such that

$$-\Delta\varphi_{\varepsilon} = \lambda f_{\varepsilon}'(u_{\varepsilon})\varphi_{\varepsilon} + R + N(\varphi_{\varepsilon}) + \sum_{j=0}^{2} \kappa_{\varepsilon,j} e^{U_{\varepsilon}} Z_{\varepsilon,j}.$$
(78)

Equivalently, setting $u_{\varepsilon} := \omega_{\varepsilon} + \varphi_{\varepsilon}$,

$$-\Delta u_{\varepsilon} = \lambda f_{\varepsilon}(u_{\varepsilon}) + \sum_{j=0}^{2} \kappa_{\varepsilon,j} e^{U_{\varepsilon}} Z_{\varepsilon,j}.$$
(79)

Our aim is to find the parameter $\mu = \mu(\varepsilon)$ and the point $\xi = \xi(\varepsilon)$ so that the $\kappa_{\varepsilon,i}$'s are zero provided ε is small enough.

Proposition 6.1 It holds true that

$$\kappa_{0,\varepsilon} = 6\pi\alpha^3 \left(2 - \log\left(\frac{8}{\mu^2}\right) + o(1)\right),\tag{80}$$

and

$$\kappa_{i,\varepsilon} = -\kappa_{0,\varepsilon} a_{i,\varepsilon} + \frac{3\mu}{2} \delta \frac{\partial v_{\varepsilon}}{\partial x_i}(\xi) + O(\alpha \delta), \ i = 1, 2$$
(81)

as $\varepsilon \to 0$ uniformly with respect to $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$. Here, the $a_{i,\varepsilon}$'s are continuous functions of μ and ξ and $a_{i,\varepsilon} = O(\alpha^2)$ uniformly for $(\mu, \xi) \in \mathcal{U} \times B(\xi_0, \sigma)$.

Proof.

Step 1 Let us prove that

$$\kappa_{i,\varepsilon} = O(\alpha^3) \quad \text{for } i = 0, 1, 2$$
(82)

and

$$\|\varphi_{\varepsilon}\|_{C^1(\overline{\Omega}\setminus B(\xi_0, 2\sigma))} = O(\alpha^3).$$
(83)

First, since (67) gives $\|\phi\|_{L^{\infty}(\Omega)} = O(\alpha^3)$, Proposition 3.13, Lemma 3.14, Remark 5.3 and Lemma 4.1 yield

$$||R||_{L^1(\Omega)} = O(\alpha^3), \quad ||N(\varphi_{\varepsilon})||_{L^1(\Omega)} = O(\alpha^5), \quad ||\lambda f_{\varepsilon}'(\omega_{\varepsilon})\varphi_{\varepsilon}||_{L^1(\Omega)} = O(\alpha^3).$$

Recalling that

$$\int_{\Omega} e^{U_n} Z_{j,n} P Z_{i,n} dx = \frac{8}{3} \pi \delta_{ij} + O(\delta), \text{ for } i, j = 0, 1, 2,$$

by Lemma 4.4 and (58), we get (82) by testing equation (78) with $PZ_{i,n}$, i = 0, 1, 2. By Lemma 3.12, Remark 5.4, and Lemma 4.1, one has

$$\lambda f_{\varepsilon}'(\omega_{\varepsilon}) = O(1), \quad R = O(\alpha^3), \quad N(\varphi_{\varepsilon}) = O(\alpha^6),$$

uniformly in $\Omega \setminus B(\xi, \frac{\sigma}{2})$. Then

$$\|\Delta\varphi_{\varepsilon}\|_{L^{\infty}(\Omega\setminus B(\xi,\frac{\sigma}{2}))} + \|\varphi_{\varepsilon}\|_{L^{\infty}(\Omega)} = O(\alpha^{3}),$$

and we get (83) by standard elliptic estimates.

Step 2 Proof of (80).

Let $\widetilde{Z}_{\varepsilon}$ be the function defined in (59). We shall test equation (78) against $\widetilde{Z}_{\varepsilon}$. With the same arguments of the proof of Proposition 4.6 (Step 1), we obtain

$$\int_{\Omega} \nabla \varphi_{\varepsilon} \cdot \nabla \widetilde{Z}_{\varepsilon} \, dx = \int_{\Omega} \nabla \varphi_{\varepsilon} \cdot \nabla Z_{0,\varepsilon} \, dx + o(\|\varphi_{\varepsilon}\|_{H^{1}_{0}(\Omega)}) = o(\alpha^{3}).$$

Moreover

$$\int_{\Omega} \lambda f_{\varepsilon}'(\omega_{\varepsilon}) \varphi_{\varepsilon} \widetilde{Z}_{\varepsilon} dx = \int_{B(\xi,\rho_0)} e^{U_{\varepsilon}} Z_{0,\varepsilon} \varphi_{\varepsilon} dx + O(\varepsilon^2 \alpha^3) + O(\alpha^3 \|\Gamma_{\varepsilon}\|_{L^1(B(\xi,\rho_1) \setminus B(\xi,\rho_0))}) = o(\alpha^3),$$

and

$$\int_{\Omega} e^{U_n} Z_{j,\varepsilon} \widetilde{Z}_{\varepsilon} dx = \int_{\mathbb{R}^2} e^{\bar{U}} Z_j Z_0 dy + O\left(\int_{\mathbb{R}^2 \setminus B(0, \frac{\rho_0}{\delta})} e^{\bar{U}} dx\right)$$
$$= \frac{8}{3} \pi \delta_{ij} + O(\delta^2 \rho_0^{-2}).$$

By Lemma 3.4 and Lemma 3.8, we get

$$\begin{split} \int_{\Omega} R \widetilde{Z}_n dx &= \int_{B(\xi,\rho_0)} R Z_{0,n} dx + O(\|R\|_{L^1(B(\xi,\rho_1) \setminus B(\xi,\rho_0))}) \\ &= \alpha^3 \int_{B(0,\rho_0)} e^{\bar{U}} \left(2\bar{U} + \bar{U}^2 \right) Z_0 dy + O\left(\alpha^4 \int_{\mathbb{R}^2} e^{\bar{U}} (1 + \bar{U}^4) dy \right) + o(\alpha^4) \\ &= 16\pi \alpha^3 \left(\log\left(\frac{8}{\mu^2}\right) - 2 \right) + O(\alpha^4). \end{split}$$

Finally, we have that

$$\int_{\Omega} N(\varphi) \widetilde{Z}_{\varepsilon} dx = O(\|N(\varphi)\|_{\varepsilon}) = O(\alpha^5).$$

Then, testing (78) against $\widetilde{Z}_{\varepsilon}$ and using (82), one gets

$$0 = 16\pi\alpha^3 \left(\log\left(\frac{8}{\mu^2}\right) - 2 \right) + \frac{8}{3}\pi k_{0,\varepsilon} + o(\alpha^3),$$

from which we get (80).

Step 3 Let us prove

$$\sum_{j=0}^{2} \kappa_{j,\varepsilon} \int_{\Omega} e^{U_{\varepsilon}} Z_{j,\varepsilon} \frac{\partial u_{\varepsilon}}{\partial x_{i}} dx = -8\pi \alpha \frac{\partial v_{\varepsilon}}{\partial x_{i}} (\xi) + O(\alpha^{2}), \ i = 1, 2,$$
(84)

We multiply (79) and $\frac{\partial u_{\varepsilon}}{\partial x_i}$. Applying the Pohozaev identity (see e.g. [27, Proposition 2, Proof of Step 1]), we obtain

$$-\frac{1}{2}\int_{\partial\Omega}\frac{\partial u_{\varepsilon}}{\partial x_{i}}\frac{\partial u_{\varepsilon}}{\partial\nu}\nu_{i}\,d\sigma = \lambda\int_{\Omega}f_{\varepsilon}(u_{\varepsilon})\frac{\partial u_{\varepsilon}}{\partial x_{i}}dx + \sum_{j=0}^{2}\kappa_{j,\varepsilon}\int_{\Omega}e^{U_{n}}Z_{j,\varepsilon}\frac{\partial u_{\varepsilon}}{\partial x_{i}}dx_{i}.$$
 (85)

Since $u_{\varepsilon} = 0$ on $\partial \Omega$, the divergence theorem yields

$$\int_{\Omega} f_{\varepsilon}(u_{\varepsilon}) \frac{\partial u_{\varepsilon}}{\partial x_{i}} dx = \int_{\Omega} \frac{d}{dx_{i}} \left(\int_{0}^{u_{\varepsilon}(x)} f_{\varepsilon}(t) dt \right) dx$$

$$= \int_{\partial \Omega} \nu_{i} \left(\int_{0}^{u_{\varepsilon}(x)} f_{\varepsilon}(t) dt \right) d\sigma = 0.$$
(86)

By (83), the definition of u_{ε} and ω_{ε} , Lemma 2.3, Lemma 3.10, we have

$$\frac{\partial u_{\varepsilon}}{\partial \nu} = -\frac{\partial v_{\varepsilon}}{\partial \nu} + \alpha \frac{\partial}{\partial \nu} (8\pi G_{\xi} - w_{\varepsilon}) + O(\alpha^2)$$

on $\partial\Omega$. Thus, keeping in mind that $|\nabla v_{\varepsilon}|$, $|\nabla w_{\varepsilon}|$ and $|\nabla G_{\xi}|$ are uniformly bounded on $\partial\Omega$ (see Lemma (2.2) and (2.3)) and that $\frac{\partial u_{\varepsilon}}{\partial x_i} = \frac{\partial u_{\varepsilon}}{\partial \nu} \nu_i$, we obtain

$$\int_{\partial\Omega} \frac{\partial u_{\varepsilon}}{\partial x_{i}} \frac{\partial u_{\varepsilon}}{\partial \nu} \, d\sigma = \int_{\partial\Omega} \frac{\partial v_{\varepsilon}}{\partial x_{i}} \frac{\partial v_{\varepsilon}}{\partial \nu} \, d\sigma + 2\alpha \int_{\partial\Omega} \frac{\partial v_{\varepsilon}}{\partial x_{i}} \frac{\partial}{\partial \nu} (w_{\varepsilon} - 8\pi G_{\xi}) \, d\sigma + O(\alpha^{2}). \tag{87}$$

Applying the Pohozaev identity to v_{ε} and arguing as in (86), we get that

$$\int_{\partial\Omega} \frac{\partial v_{\varepsilon}}{\partial x_i} \frac{\partial v_{\varepsilon}}{\partial \nu} \, d\sigma = -2\lambda \int_{\Omega} f_{\varepsilon}(v_{\varepsilon}) \frac{\partial v_{\varepsilon}}{\partial x_i} dx = 0.$$
(88)

Integrating by parts and noting that $-\Delta \frac{\partial v_{\varepsilon}}{\partial x_i} = \lambda f'_{\varepsilon}(v_{\varepsilon}) \frac{\partial v_{\varepsilon}}{\partial x_i}$ in Ω , we get

$$\int_{\partial\Omega} \frac{\partial v_{\varepsilon}}{\partial x_{i}} \frac{\partial}{\partial \nu} (w_{\varepsilon} - 8\pi G_{\xi}) d\sigma = \int_{\Omega} \left(\frac{\partial v_{\varepsilon}}{\partial x_{i}} \Delta w_{\varepsilon} - w_{\varepsilon} \Delta \frac{\partial v_{\varepsilon}}{\partial x_{i}} \right) dx + 8\pi \frac{\partial v_{\varepsilon}}{\partial x_{i}} (\xi) + 8\pi \int_{\Omega} G_{\xi} \Delta \frac{\partial v_{\varepsilon}}{\partial x_{i}} dx = \int_{\Omega} \frac{\partial v_{\varepsilon}}{\partial x_{i}} \underbrace{\left(\Delta w_{\varepsilon} + \lambda f_{\varepsilon}'(v_{\varepsilon}) w_{\varepsilon} - 8\pi \lambda f_{\varepsilon}'(v_{\varepsilon}) G_{\xi} \right)}_{=0 \text{ by (12)}} dx + 8\pi \frac{\partial v_{\varepsilon}}{\partial x_{i}} (\xi).$$

This together with (87)-(88) gives

$$\frac{1}{2} \int_{\partial\Omega} \frac{\partial u_{\varepsilon}}{\partial x_i} \frac{\partial u_{\varepsilon}}{\partial \nu} d\sigma = 8\pi \alpha \frac{\partial v_{\varepsilon}}{\partial x_i} (\xi) + O(\alpha^2).$$
(89)

Finally, (84) follows by (85)-(86) and (89).

Step 4 For i = 1, 2, j = 0, 1, 2, we have

$$\int_{\Omega} e^{U_{\varepsilon}} Z_{j,\varepsilon} \frac{\partial u_{\varepsilon}}{\partial x_i} dx = -\frac{\alpha}{\delta} \left(\frac{16}{3\mu} \pi \delta_{ij} + O(\alpha^2) \right).$$
(90)

For i = 1, 2 and j = 0, 1, 2. Note that we have the identity

$$\frac{\partial}{\partial x_i} e^{U_{\varepsilon}} Z_{j,\varepsilon} = \frac{e^{U_{\varepsilon}}}{\delta \mu} \left(\delta_{ij} (Z_{0,\varepsilon} + 1) - \delta_{j0} Z_{i,\varepsilon} - 3 Z_{i,\varepsilon} Z_{j,\varepsilon} \right).$$

Setting $\Psi_{ij} := \delta_{ij}(Z_0 + 1) - \delta_{j0}Z_i - 3Z_iZ_j$ and applying the divergence theorem, we find

$$\begin{split} \int_{\Omega} e^{U_{\varepsilon}} Z_{j,\varepsilon} \frac{\partial u_{\varepsilon}}{\partial x_{i}} dx &= -\int_{\Omega} u_{\varepsilon} \frac{d}{dx_{i}} \left(e^{U_{\varepsilon}} Z_{j,\varepsilon} \right) dx \\ &= -\frac{1}{\delta \mu} \int_{\Omega} u_{\varepsilon} e^{U_{\varepsilon}} \left(\delta_{ij} (Z_{0,\varepsilon} + 1) - \delta_{j0} Z_{i,\varepsilon} - 3 Z_{i,\varepsilon} Z_{j,\varepsilon} \right) dx \\ &= -\frac{1}{\delta \mu} \int_{\Omega - \frac{\varepsilon}{\delta}} u_{\varepsilon} (\xi + \delta y) e^{\bar{U}} \Psi_{ij} dy \\ &= -\frac{1}{\delta \mu} \int_{B(0,\frac{\sigma}{\delta})} u_{\varepsilon} (\xi + \delta y) e^{\bar{U}} \Psi_{ij} dy + O(\beta \delta^{2}), \end{split}$$

where in the last equality we used that

$$u_{\varepsilon} = O(\beta)$$
 and $e^{\bar{U}}\Psi_{ij} = O(|y|^{-5}),$ (91)

for $|y| \ge \frac{\sigma}{\delta}$. By Lemma 2.6 we have

$$u_{\varepsilon}(\xi + \delta y) = \beta + \alpha \bar{U}(y) + O(\alpha^3) + O(\delta|y|), \qquad (92)$$

for $y \in B(0, \frac{\sigma}{\delta})$. Using again (91), we get that

$$\int_{B(0,\frac{\sigma}{\delta})} e^{\bar{U}} \Psi_{ij} dy = \underbrace{\int_{\mathbb{R}^2} e^{\bar{U}} \Psi_{ij} dy}_{=0} + O(\delta^3).$$

Similarly, we have

$$\begin{split} \int_{B(0,\frac{\sigma}{\delta})} \overline{U} e^{\bar{U}} \Psi_{ij} dy &= \int_{\mathbb{R}^2} \overline{U} e^{\bar{U}} \Psi_{ij} dy + O(\beta^2 \delta^3) \\ &= \frac{16}{3} \pi \delta_{ij} + O(\beta^2 \delta^3), \end{split}$$

and (90) is proved.

Step 5 Proof of (81).

Let us set

$$a_{ij,arepsilon} = a_{ij,arepsilon}(\xi,\mu) := -rac{3\mu}{16\pi}rac{\delta}{lpha}\int_{\Omega}e^{U_{arepsilon}}Z_{j,arepsilon}rac{\partial u_{arepsilon}}{\partial x_{i}}d\sigma.$$

According to Step 4, we have $a_{i0,\varepsilon} = O(\alpha^2)$ if i = 1, 2. Moreover the matrix $A = (a_{ij,\varepsilon})_{i,j\in\{1,2\}}$ is invertible and its inverse $A^{-1} = (a_{\varepsilon}^{ij})_{ij\in\{1,2\}}$ satisfies

$$a_{\varepsilon}^{ij} = \delta_{ij} + O(\alpha^2), \quad i, j = 1, 2.$$

Then (81) follows by (84), just setting

$$a_{i,\varepsilon} := \sum_{j=1}^{2} a_{\varepsilon}^{ij} a_{0j,\varepsilon}$$

It is important to point out that (81) cannot be considered a precise uniform expansion of $\kappa_{i,\varepsilon}$. Indeed, (80) and the rough (but difficult to improve) estimate $a_{i,\varepsilon} = O(\alpha^2)$ yield only $\kappa_{0,\varepsilon}a_{i,\varepsilon} = O(\alpha^5)$. Since $\delta \ll \alpha^5$ it is not possible to identify the leading term in the RHS of (81). However, it is clear that the term involving $\frac{\partial v_{\varepsilon}}{\partial x_i}$ becomes dominant when $\kappa_{0,\varepsilon}$ vanishes. This is enough for our argument.

Proof of Theorem 1.3 completed

Proof. Let us consider the vector field

$$B_{\varepsilon}(\mu,\xi) = \left(\frac{1}{6\pi\alpha^3}\kappa_{0,\varepsilon}, \frac{2}{3\delta\mu}\left(\kappa_{1,\varepsilon} + \kappa_{0,\varepsilon}a_{1,\varepsilon}\right), \frac{2}{3\delta\mu}\left(\kappa_{2,\varepsilon} + \kappa_{0,\varepsilon}a_{2,\varepsilon}\right)\right).$$

By construction, for any $\varepsilon > 0$, B_{ε} depends continuously on μ and ξ . Moreover, thanks to (80), (81) and Lemma 2.2, we have

$$B_{\varepsilon} \to \bar{B}(\mu,\xi) := \left(2 - \log\left(\frac{8}{\mu^2}\right), \nabla u_0(\xi)\right)$$

as $\varepsilon \to 0$, uniformly for $\mu \in \mathcal{U}$ and $\xi \in B(\xi_0, \sigma)$. By assumption (A2), \overline{B} has a C^0 -stable zero at the point (μ_0, ξ_0) , with $\mu_0 = \sqrt{8}e^{-1}$. Then, for ε small enough, there exist $\xi = \xi(\varepsilon) \to \xi_0, \ \mu = \mu(\varepsilon) \to \mu_0$ as $\varepsilon \to 0$ such that $B_{\varepsilon}(\mu(\varepsilon), \xi(\varepsilon)) = 0$. Clearly, this is equivalent to $\kappa_{i,\varepsilon,\mu(\varepsilon),\xi(\varepsilon)} = 0, \ i = 0, 1, 2$. That concludes the proof. \Box

Appendix A. The proof of Lemma 2.4

Proof. The third equation in (28) allows to write δ as a function of $\alpha, \beta, \varepsilon, \mu, \xi$:

$$\log \frac{1}{\delta^2} = \frac{\beta}{2\alpha} + \frac{V_{\varepsilon,\alpha,\xi}(\xi)}{2\alpha} - \frac{c_{\mu,\xi}}{2},$$

and the second equation in (28) gives α as a function of $\beta, \varepsilon, \mu, \xi$:

$$\alpha = (2\beta + \beta^{\varepsilon} + \varepsilon\beta^{\varepsilon})^{-1}$$

Then, (after a simple computation) it is sufficient to prove that there exists $\beta = \beta(\varepsilon, \mu, \xi)$ such that

$$\frac{1}{\beta} \left(\log \lambda + \frac{c_{\mu,\xi}}{2} \right) + 2 \frac{\log \beta}{\beta} + \underbrace{\left(\frac{1}{2} \beta^{\varepsilon} - u_0(\xi) \right)}_{:=\theta_{\varepsilon}(\xi,\mu)} - \left(V_{\varepsilon,\alpha,\xi}(\xi) - u_0(\xi) \right) \\
+ \frac{\log \left(2 + \beta^{\varepsilon-1} + \varepsilon \beta^{\varepsilon-1} \right)}{\beta} - \frac{1}{2} \varepsilon \beta^{\varepsilon} - \frac{1}{2} V_{\varepsilon,\alpha,\xi}(\xi) \left(\beta^{\varepsilon-1} + \varepsilon \beta^{\varepsilon-1} \right) = 0.$$
(93)

Now, we choose $\beta^{\varepsilon} := 2u_0(\xi) + \theta_{\varepsilon}(\xi,\mu)$ with $\|\theta_{\varepsilon}\|_{C^0(\overline{B(\xi_0,\sigma)}\times\overline{\mathcal{U}})}$ so small that

$$2u_0(\xi) + \theta_{\varepsilon}(\xi,\mu) \ge \eta > 1$$
 in $\overline{B(\xi_0,\sigma)} \times \overline{\mathcal{U}}$.

This is possible because of (22). With this choice we have $\frac{1}{\beta} = O\left(\eta^{-\frac{1}{\varepsilon}}\right)$. It is easy to show that (93) has a solution θ_{ε} because of a simple fixed point argument. Indeed (93) rewrites as $\theta_{\varepsilon} = \mathcal{T}(\theta_{\varepsilon})$ where \mathcal{T} is a contraction mapping on the ball

$$\left\{\theta_{\varepsilon} \in C^{0}(\overline{B(\xi_{0},\sigma)} \times \overline{\mathcal{U}}) : \|\theta_{\varepsilon}\|_{C^{0}(\overline{B(\xi_{0},\sigma)} \times \overline{\mathcal{U}})} \leq \rho_{\varepsilon}\right\},\$$

where $\rho_{\varepsilon} := \rho \min\left\{\frac{1}{\varepsilon}\eta^{-\frac{1}{\varepsilon}}, \|v_{\varepsilon} - u_0\|_{C^0(\overline{\Omega})}\right\}$ for some $\rho > 0$ and $\rho_{\varepsilon} \to 0$ as $\varepsilon \to 0$. Here we use the expression of $V_{\varepsilon,\alpha,\xi}(\xi)$ in (11) and (ii) of Lemma 2.2.

Appendix B. A Stampacchia type estimate

In this section we prove domain-independent estimates for solutions of the Poisson equation $-\Delta u = f$, under Dirichlet boundary conditions, with $f \in L^p(\Omega)$ and p approaching 1. Our strategy is based on the Stampacchia method. **Lemma B.1 ([28], Lemma 4.1)** Let $\psi : \mathbb{R}^+ \mapsto \mathbb{R}^+$ be a nonincreasing function. Assume that there exist $M > 0, \gamma > 0, \delta > 1$ such that

$$\psi(h) \leq \frac{M\psi(k)^{\delta}}{(h-k)^{\gamma}} \qquad \forall \ h > k > 0.$$

Then $\psi(d) = 0$, where $d = M^{\frac{1}{\gamma}} \psi(0)^{\frac{\delta-1}{\gamma}} 2^{\frac{\delta}{\delta-1}}$.

Let $\Omega \subseteq \mathbb{R}^2$ be a bounded smooth domain. For any q > 1, let $S_q(\Omega)$ be the Sobolev's constant for the embedding of $H_0^1(\Omega)$ in $L^q(\Omega)$, namely

$$S_q(\Omega) = \inf_{u \in H_0^1(\Omega)} \frac{\|u\|_{H_0^1(\Omega)}}{\|u\|_{L^q(\Omega)}}.$$

It is known that $0 < S_q(\Omega) < +\infty$ and that (see [26] Lemma 2.2)

$$\lim_{q \to +\infty} \sqrt{q} S_q(\Omega) = \sqrt{8\pi e}$$

Theorem B.2 Let Ω be a bounded smooth domain. For p > 1, $f \in L^p(\Omega)$, the unique solution $u \in H^1_0(\Omega)$ of the equation $-\Delta u = f$ satisfies

$$\|u\|_{L^{\infty}(\Omega)} \leq 4S_{\frac{3p+1}{p-1}}(\Omega)^{-2} \|f\|_{L^{p}} |\Omega|^{\frac{p^{2}-1}{3p^{2}+p}}.$$

Proof. We want to apply the previous lemma to the function

$$\psi(k) := |A_k|, \quad A_k := \{ x \in \Omega : |u(x)| > k \}.$$

For any k > 0, let us consider the function

$$v_k(x) := \begin{cases} 0 & |u(x)| \le k, \\ u(x) - k & u(x) > k, \\ -u(x) - k & u(x) < -k. \end{cases}$$

Note that $v_k \in H_0^1(\Omega)$ and $|\nabla v_k| = |\nabla u|\chi_{A_k}$. If we test the equation against v_k we get

$$\int_{\Omega} \nabla u \cdot \nabla v_k \, dx = \int_{\Omega} f v_k dx. \tag{94}$$

For any $q \in (1, p)$ Hölder's inequality gives

$$\int_{\Omega} f v_k \, dx = \int_{A_k} f v_k \, dx \le \|f\|_{L^q(A_k)} \|v_k\|_{L^{\frac{q}{q-1}}(A_k)} \le \|f\|_{L^p} |A_k|^{\frac{p-q}{pq}} \|v_k\|_{L^{\frac{q}{q-1}}(A_k)}.$$
 (95)

By Sobolev's inequality, we have that

$$\int_{\Omega} \nabla u \cdot \nabla v_k \, dx = \int_{A_k} |\nabla v_k|^2 dx \ge S_{\frac{q}{q-1}}(\Omega)^2 \|v_k\|_{L^{\frac{q}{q-1}}}^2.$$
(96)

By (94)-(96), we have

$$\|v_k\|_{L^{\frac{q}{q-1}}} \le S_{\frac{q}{q-1}}(\Omega)^{-2} \|f\|_{L^p} |A_k|^{\frac{p-q}{pq}}.$$

Now, for any h > k, we have that $A_h \subseteq A_k$ and $v_k \ge (h - k)$ in A_h , hence

$$\int_{\Omega} |v_k|^{\frac{q}{q-1}} dx = \int_{A_k} v_k^{\frac{q}{q-1}} dx \ge \int_{A_h} v_k^{\frac{q}{q-1}} dx \ge (h-k)^{\frac{q}{q-1}} |A_h|.$$

In conlcusion, we find

$$(h-k)|A_h|^{\frac{q-1}{q}} \le S_{\frac{q}{q-1}}(\Omega)^{-2} ||f||_{L^p} |A_k|^{\frac{p-q}{pq}},$$

or, equivalently,

$$\psi(h) \le \frac{S_{\frac{q}{q-1}}(\Omega)^{-\frac{2q}{q-1}} \|f\|_{L^p}^{\frac{q}{q-1}} \psi(k)^{\frac{p-q}{p(q-1)}}}{(h-k)^{\frac{q}{q-1}}}$$

Then, we are in position to apply Lemma B.1 to ψ with $M = S_{\frac{q}{q-1}}(\Omega)^{-\frac{2q}{q-1}} \|f\|_{L^p}^{\frac{q}{q-1}}$, $\gamma = \frac{q}{q-1}$, and $\delta = \frac{p-q}{p(q-1)}$. For this, we need to impose that $\delta = \frac{p-q}{p(q-1)}$, that is $q < \frac{2p}{p+1}$. Note that $1 < \frac{2p}{p+1} < p$. According to Stampacchia's Lemma, we have

$$\psi(d) = 0 \qquad \text{where} \qquad d = M^{\frac{1}{\gamma}} \psi(0)^{\frac{\delta-1}{\gamma}} 2^{\frac{\delta}{\delta-1}} = S^{2}_{\frac{q}{q-1}} \|f\|_{L^{p}} |\Omega|^{\frac{2p-q(p+1)}{pq}} 2^{\frac{p-q}{2p-q(p+1)}}.$$

This implies that

$$||u||_{L^{\infty}(\Omega)} \leq S_{\frac{q}{q-1}}(\Omega)^{-2} ||f||_{L^{p}} |\Omega|^{\frac{2p-q(p+1)}{pq}} 2^{\frac{p-q}{2p-q(p+1)}}.$$

This is true for any choice of $q \in (1, \frac{2p}{p+1})$. If we take for example p the midpoint of $(1, \frac{2p}{p+1})$, that is $q = \frac{1}{2} + \frac{p}{p+1} = \frac{3p+1}{2(p+1)}$, then we find that

$$\frac{q}{q-1} = \frac{3p+1}{p-1}, \quad \frac{2p-q(p+1)}{pq} = \frac{p^2-1}{3p^2+p}, \quad \frac{p-q}{2p-q(p+1)} = \frac{2p+1}{p+1} \le 2,$$

e get the conclusion.

and we get the conclusion.

Corollary B.3 Given K > 0 and p > 1, there exists a constant C = C(K, p) such that, for any domain $\Omega \subseteq \mathbb{R}^2$ with $|\Omega| \leq K$ and any $f \in L^p(\Omega)$ the unique solution $u \in H^1_0(\Omega)$ of $-\Delta u = f$ satisfies

$$\|u\|_{L^{\infty}(\Omega)} \le C \|f\|_{L^{p}(\Omega)}.$$

Corollary B.4 Given K > 0, there exist $p_0 = p_0(K)$ and C = C(K) such that, for any $1 , any domain <math>\Omega \subseteq \mathbb{R}^2$ with $|\Omega| \leq K$, and any $f \in L^p(\Omega)$, the unique solution $u \in H^1_0(\Omega)$ of $-\Delta u = f$ satisfies

$$\|u\|_{L^{\infty}(\Omega)} \leq \frac{C}{p-1} \|f\|_{L^{p}(\Omega)}.$$

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