# Convergence of Adaptive Filtered schemes for first order evolutive Hamilton-Jacobi equations * 

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#### Abstract

We consider a class of "filtered" schemes for first order time dependent Hamilton-Jacobi equations and prove a general convergence result for this class of schemes. A typical filtered scheme is obtained mixing a high-order scheme and a monotone scheme according to a filter function $F$ which decides where the scheme has to switch from one scheme to the other. A crucial role for this switch is played by a parameter $\varepsilon=\varepsilon(\Delta t, \Delta x)>0$ which goes to 0 as ( $\Delta t, \Delta x$ ) is going to 0 and does not depend on the time $t_{n}$. The tuning of this parameter in the code is rather delicate and has an influence on the global accuracy of the filtered scheme. Here we introduce an adaptive and automatic choice of $\varepsilon=\varepsilon^{n}(\Delta t, \Delta x)$ at every iteration modifying the classical set up. The adaptivity is controlled by a smoothness indicator which selects the regions where we modify the regularity threshold $\varepsilon^{n}$. A convergence result and some error estimates for the new adaptive filtered scheme are proved, this analysis relies on the properties of the scheme and of the smoothness indicators. Finally, we present some numerical tests to compare the adaptive filtered scheme with other methods.


Keywords: High-order Filtered schemes Hamilton-Jacobi equations Convergence Smoothness indicators

## 1 Introduction

Here we propose and analyze a new adaptive filter scheme and prove its convergence to the viscosity solution of the scalar evolutive Hamilton-Jacobi equation

$$
\begin{cases}v_{t}+H\left(v_{x}\right)=0, & (t, x) \in[0, T] \times \mathbb{R},  \tag{1}\\ v(0, x)=v_{0}(x), & x \in \mathbb{R},\end{cases}
$$

where Hamiltonian $H$ and the initial data $v_{0}$ are Lipschitz continuous functions. A precise result of existence and uniqueness in the framework of weak viscosity solutions can be found in 44 and the precise setting of assumptions will be given in Section 2.

[^0]The accurate numerical solution of Hamilton-Jacobi (HJ) equations is a challenging topic of growing importance in many fields of application, e.g. control theory, KAM theory, image processing and material science. Due to the lack of regularity of viscosity solutions, this issue is delicate and the construction of high-order methods can be rather complicated and the proof of convergence is challenging. It is well known that simple monotone schemes are at most first order accurate as shown in [7] so monotonicity should be abandoned to get high-order convergence. Our goal is to present a rather simple way to construct convergent schemes to the viscosity solution $v$ of (11) with the property to be of high-order in the region of regularity.
In recent years a general approach to the construction of high-order methods using filters has been proposed by Lions and Souganidis in [17] and further developed by Oberman and Salvador [18]. Let us remind that a typical feature of a filtered scheme $S^{F}$ is that at the node $x_{j}$ the scheme is combination of a high-order scheme $S^{A}$ and a monotone scheme $S^{M}$ according to a filter function $F$. The scheme is written as

$$
\begin{equation*}
u_{j}^{n+1} \equiv S^{F}\left(u^{n}\right)_{j}:=S^{M}\left(u^{n}\right)_{j}+\varepsilon \Delta t F\left(\frac{S^{A}\left(u^{n}\right)_{j}-S^{M}\left(u^{n}\right)_{j}}{\epsilon \Delta t}\right), \quad j \in \mathbb{Z}, \tag{2}
\end{equation*}
$$

where $\varepsilon=\varepsilon_{\Delta t, \Delta x}>0$ is a fixed parameter going to 0 as $(\Delta t, \Delta x)$ is going to 0 and does not depend on $n$. Filtered schemes are high-order accurate where the solution is smooth, monotone otherwise, and this feature is crucial to prove a convergence result as in [5. Note that the choice of the parameter $\varepsilon$ is delicate because it plays a crucial role in the switching so its tuning is rather important (see 5 for a detailed discussion of this point). Then it seems natural to adapt its choice to the regularity of the solution in the cell via a smoothness indicator. Here we improve the filtered scheme (2) introducing an adaptive and automatic choice of the parameter $\varepsilon=\varepsilon^{n}$ at every iteration. To this end, we introduce a smoothness indicator to select the regions where we have to update the regularity threshold $\varepsilon^{n}$, this indicator is chosen according to the analysis proposed in [15] although other proposals with similar properties can be applied.

To set this paper into perspective let us remind that the construction of high order methods for hyperbolic equations has been a very active research area started by the seminal paper [14. Several techniques have been proposed to improve the accuracy leading to essentially non oscillatory schemes ENO and weighted ENO (so called WENO) for conservation laws as in [13, 2, 12, 1], for a survey on these high-order techniques we refer to [20, 21]. More recently a centered and more efficient version (called CWENO) has been proposed in [8]. We should also mention that high-order methods have been proposed for Hamilton-Jacobi either extending the ENO approach as in [15, 16, 6] or by semi-lagrangian techniques as extensively discussed in [10]. For a recent survey on the numerical approximation of Hamilton-Jacobi equations we refer the interested reader to [9].

The paper is organized as follows:
In Section 2 we construct the new adaptive filtered scheme and present in detail all its building blocks, the main assumptions are given there. Section 3 is focused on the analysis of the smoothness indicators in one dimension. In Section 4 we state and prove the main convergence result (that was announced in [11), some technical lemmas are proved in Appendix A at the end of this paper. Finally in Sect. 5 we present several tests to show the effectiveness of the adaptive scheme with respect to the basic filtered scheme and to other state-of-the-art methods. Sect. 6 contains the conclusions with final comments.

## 2 A new Adaptive Filtered scheme

Consider the first order evolutive Hamilton-Jacobi equation (11) where the hamiltonian $H$ and the initial data $v_{0}$ are Lipschitz continuous functions. It is well known that with these assumptions we have the existence and uniqueness of the viscosity solution. Notice that to keep the ideas clear we are considering the most simple scalar case with the hamiltonian depending only on the derivative of the solution, with more general situations following directly. Our aim is to present a rather simple way to construct convergent schemes to the viscosity solution $v$ of (1) with the property to be of high-order whenever some regularity is detected.

Starting from the ideas of [5] on filtered schemes, we proceed in this study introducing a procedure to compute the regularity threshold $\varepsilon$ in an automatic way, in order to exploit the local regularity of the solution.

Let us begin defining a uniform grid in space $x_{j}=j \Delta x, j \in \mathbb{Z}$, and in time $t_{n}=t_{0}+n \Delta t, n \in$ $[0, N]$, with $(N-1) \Delta t<T \leq N \Delta t$. Then, we compute the numerical approximation $u_{j}^{n}=u\left(t_{n}, x_{j}\right)$ with the simple formula

$$
\begin{equation*}
u_{j}^{n+1}=S^{A F}\left(u^{n}\right)_{j}:=S^{M}\left(u^{n}\right)_{j}+\phi_{j}^{n} \varepsilon^{n} \Delta t F\left(\frac{S^{A}\left(u^{n}\right)_{j}-S^{M}\left(u^{n}\right)_{j}}{\varepsilon^{n} \Delta t}\right) \tag{3}
\end{equation*}
$$

where $u_{j}^{n+1}:=u\left(t_{n+1}, x_{j}\right), S^{M}$ and $S^{A}$ are respectively the monotone and the high-order scheme, $F$ is the filter function needed to switch between the two schemes, $\varepsilon^{n}$ is the switching parameter at time $t_{n}$ and $\phi_{j}^{n}$ is the smoothness indicator function at the node $x_{j}$ and time $t_{n}$. More details on the components of the schemes will be given in the following sections.

Notice that if $\varepsilon^{n} \equiv \varepsilon \Delta x$, with $\varepsilon>0$ and $\phi_{j}^{n} \equiv 1$, we get the Basic Filtered Scheme (2).

### 2.1 Assumptions on the schemes

In this section we present in detail the basic components of our scheme, which are a monotone finite difference scheme $S^{M}$ and a high-order, possibly unstable, scheme $S^{A}$. Let us begin by giving the assumptions on the monotone scheme.

Assumptions on $S^{M}$.
(M1) The scheme can be written in differenced form

$$
u_{j}^{n+1} \equiv S^{M}\left(u_{j}^{n}\right):=u_{j}^{n}-\Delta t h^{M}\left(D^{-} u_{j}^{n}, D^{+} u_{j}^{n}\right)
$$

for a function $h^{M}\left(p^{-}, p^{+}\right)$, with $D^{ \pm} u_{j}^{n}:= \pm \frac{u_{j \pm 1}^{n}-u_{j}^{n}}{\Delta x} ;$
(M2) $h^{M}$ is a Lipschitz continuous function;
(M3) (Consistency) $\forall v, h^{M}(v, v)=H(v)$;
(M4) (Monotonicity) for any functions $u, v$,

$$
u \leq v \quad \Rightarrow \quad S^{M}(u) \leq S^{M}(v)
$$

Under assumption (M2), the consistency property (M3) is equivalent to say that for all functions $v \in C^{2}([0, T] \times \mathbb{R})$, there exists a constant $C_{M} \geq 0$ independent on $\Delta=(\Delta t, \Delta x)$ such that

$$
\begin{equation*}
\mathcal{E}_{M}(v)(t, x):=\left|\frac{v(t+\Delta t, x)-S^{M}(v(t, \cdot))(x)}{\Delta t}\right| \leq C_{M}\left(\Delta t\left\|v_{t t}\right\|_{\infty}+\Delta x\left\|v_{x x}\right\|_{\infty}\right) \tag{4}
\end{equation*}
$$

where $\mathcal{E}_{M}$ is the consistency error. The last relation clearly shows the bound on the accuracy of the monotone schemes, which are at most first order accurate even for regular solutions.

Remark 2.1. As pointed out in [5], under the Lipschitz assumption (M2) the monotonicity property (M4) can be restated in terms of some quantities that can be easily computed. In fact, it is enough to require, for a.e. $\left(p^{-}, p^{+}\right) \in \mathbb{R}^{2}$,

$$
\begin{equation*}
\frac{\partial h^{M}}{\partial p^{-}}\left(p^{-}, p^{+}\right) \geq 0, \quad \frac{\partial h^{M}}{\partial p^{+}}\left(p^{-}, p^{+}\right) \leq 0 \tag{5}
\end{equation*}
$$

and the CFL condition

$$
\begin{equation*}
\frac{\Delta t}{\Delta x}\left(\frac{\partial h^{M}}{\partial p^{-}}\left(p^{-}, p^{+}\right)-\frac{\partial h^{M}}{\partial p^{+}}\left(p^{-}, p^{+}\right)\right) \leq 1 \tag{6}
\end{equation*}
$$

We call the CFL number, dependent on the hamiltonian of the considered problem, the constant ratio $\lambda:=\frac{\Delta t}{\Delta x}$ such that (6) is satisfied. Notice that working with explicit finite difference schemes this number can always be computed.

Example 2.2. We give some examples of monotone schemes in differenced form which satisfy (M1)-(M4). Other examples may be found in the pioneering work [7] or in [20].

- For the eikonal equation,

$$
v_{t}+\left|v_{x}\right|=0
$$

we can use the simple numerical hamiltonian

$$
\begin{equation*}
h^{M}\left(p^{-}, p^{+}\right):=\max \left\{p^{-},-p^{+}\right\} . \tag{7}
\end{equation*}
$$

- For general equations, instead, we recall the Central Upwind scheme of [16]

$$
\begin{equation*}
h^{M}\left(p^{-}, p^{+}\right):=\frac{1}{a^{+}-a^{-}}\left[a^{-} H\left(p^{+}\right)-a^{+} H\left(p^{-}\right)-a^{+} a^{-}\left(p^{+}-p^{-}\right)\right], \tag{8}
\end{equation*}
$$

with $a^{+}=\max \left\{H_{p}\left(p^{-}\right), H_{p}\left(p^{+}\right), 0\right\}$ and $a^{-}=\min \left\{H_{p}\left(p^{-}\right), H_{p}\left(p^{+}\right), 0\right\}$, using the usual notation $H_{p}$ for the derivative of $H$ with respect to $v_{x}$.

- Another numerical hamiltonian we could use is the Lax-Friedrichs hamiltonian

$$
\begin{equation*}
h^{M}\left(p^{-}, p^{+}\right):=H\left(\frac{p^{-}+p^{-}}{2}\right)-\frac{\theta}{2}\left(p^{+}-p^{-}\right) \tag{9}
\end{equation*}
$$

where $\theta>0$ is a constant. The scheme is monotone under the restrictions $\max _{p}\left|H_{p}(p)\right|<\theta$ and $\theta \lambda \leq 1$.

Next, we define the requirements on the high-order scheme.

Assumptions on $S^{A}$.
(A1) The scheme can be written in differenced form

$$
u_{j}^{n+1}=S^{A}\left(u^{n}\right)_{j}:=u_{j}^{n}-\Delta t h^{A}\left(D^{k,-} u_{j}, \ldots, D^{-} u_{j}^{n}, D^{+} u_{j}^{n}, \ldots, D^{k,+} u_{j}^{n}\right)
$$

for some function $h^{A}\left(p^{-}, p^{+}\right)$(in short), with $D^{k, \pm} u_{j}^{n}:= \pm \frac{u_{j \pm k}^{n}-u_{j}^{n}}{k \Delta x}$;
(A2) $h^{A}$ is a Lipschitz continuous function.
(A3) (High-order consistency) Fix $k \geq 2$ order of the scheme, then for all $l=1, \ldots, k$ and for all functions $v \in C^{l+1}$, there exists a constant $C_{A, l} \geq 0$ such that

$$
\begin{aligned}
\mathcal{E}_{A}(v)(t, x):= & \left|\frac{v(t+\Delta t, x)-S^{A}(v(t, \cdot))(x)}{\Delta t}\right| \\
& \leq C_{A, l}\left(\Delta t^{l}\left\|\partial_{t}^{l+1} v\right\|_{\infty}+\Delta x^{l}\left\|\partial_{x}^{l+1} v\right\|_{\infty}\right)
\end{aligned}
$$

It is interesting to notice that we are not making any assumption on the stability of the high-order scheme, that is because filtered schemes are able to stabilize a possibly unstable scheme.

Before giving some examples of high-order schemes satisfying (A1)-(A3), let us state an interesting property of the solution $v$ of (1) in case of enough regularity. Notice that we are considering the simplest case of $H$ dependent only on the gradient of $v$.

Lemma 2.1. Let $v$ be the solution of (1). Then, if $v \in C^{r}\left(\Omega_{(t, x)}\right), r \geq 2$, where $\Omega_{(t, x)}$ is a neighborhood of a point $(t, x) \in \Omega:=[0, T] \times \mathbb{R}$, it holds

$$
\begin{align*}
\frac{\partial^{k} v(t, x)}{\partial t^{k}} & =(-1)^{k} \frac{\partial^{k-2}}{\partial x^{k-2}}\left(H_{p}^{k}\left(v_{x}(t, x)\right) v_{x x}(t, x)\right)  \tag{10}\\
& =(-1)^{k} \frac{\partial^{k-2}}{\partial x^{k-2}}\left(H_{p}^{k-1}\left(v_{x}(t, x)\right) \frac{\partial}{\partial x} H\left(v_{x}(t, x)\right)\right)
\end{align*}
$$

for $k=2, \ldots, r$.
Proof. Let us proceed by induction on $2 \leq k \leq r$, omitting the dependence on $(t, x)$ to simplify the notation. For $k=2$, we have

$$
v_{t t}=\frac{\partial}{\partial t}\left(-H\left(v_{x}\right)\right)=-H_{p}\left(v_{x}\right) v_{x t}=-H_{p}\left(v_{x}\right) \frac{\partial}{\partial x}\left(-H\left(v_{x}\right)\right)=H_{p}^{2}\left(v_{x}\right) v_{x x}
$$

and the statement holds in this case. Suppose now that 10 holds for $2<k<r-1$, then we can
compute

$$
\begin{aligned}
\frac{\partial^{k+1} v}{\partial t^{k+1}} & =\frac{\partial}{\partial t}\left(\frac{\partial^{k} v}{\partial t^{k}}\right) \\
& \left.=\frac{\partial}{\partial t}\left((-1)^{k} \frac{\partial^{k-2}}{\partial x^{k-2}}\left(H_{p}^{k}\left(v_{x}\right)\right) v_{x x}\right)\right) \quad \text { by inductive hypothesis } \\
& \left.=(-1)^{k} \frac{\partial^{k-2}}{\partial x^{k-2}}\left(\frac{\partial}{\partial t}\left(H_{p}^{k}\left(v_{x}\right)\right) v_{x x}\right)\right) \\
& =(-1)^{k} \frac{\partial^{k-2}}{\partial x^{k-2}}\left(\frac{\partial}{\partial p}\left(H_{p}^{k}\left(v_{x}\right)\right) v_{x t} v_{x x}+H_{p}^{k}\left(v_{x}\right) v_{x x t}\right) \\
& =(-1)^{k} \frac{\partial^{k-2}}{\partial x^{k-2}}\left(\frac{\partial}{\partial x}\left(H_{p}^{k}\left(v_{x}\right)\right) v_{x t}+H_{p}^{k}\left(v_{x}\right) \frac{\partial}{\partial x}\left(v_{x t}\right)\right) \\
& =(-1)^{k} \frac{\partial^{k-1}}{\partial x^{k-1}}\left(H_{p}^{k}\left(v_{x}\right) v_{t x}\right) \\
& =(-1)^{k+1} \frac{\partial^{k-1}}{\partial x^{k-1}}\left(H_{p}^{k+1}\left(v_{x}\right) v_{x x}\right),
\end{aligned}
$$

as we wanted.
Let us now consider the value of the solution at $v(t+\Delta t, x)$, with $\Delta t>0$ and its Taylor expansion of order $r \geq 2$ around the point $(t, x)$. Using Lemma 2.1, we can rewrite

$$
\begin{align*}
v(t+\Delta t, x)= & v(t, x)+\Delta t v_{t}(t, x)+\sum_{k=2}^{r} \frac{\Delta t^{k}}{k!} \frac{\partial^{k} v(t, x)}{\partial t^{k}}+O\left(\Delta t^{r+1}\right) \\
= & v(t, x)-\Delta t H\left(v_{x}(t, x)\right)+ \\
& \sum_{k=2}^{r} \frac{(-\Delta t)^{k}}{k!} \frac{\partial^{k-2}}{\partial x^{k-2}}\left(H_{p}^{k}\left(v_{x}(t, x)\right) v_{x x}(t, x)\right)+O\left(\Delta t^{r+1}\right), \tag{11}
\end{align*}
$$

which for $r=2$ simply reads

$$
\begin{equation*}
v(t+\Delta t, x)=v(t, x)-\Delta t H\left(v_{x}(t, x)\right)+\frac{\Delta t^{2}}{2} H_{p}^{2}\left(v_{x}(t, x)\right) v_{x x}(t, x)+O\left(\Delta t^{3}\right) . \tag{12}
\end{equation*}
$$

Remark 2.3. Using this last relation we could show that, assuming (A1)-(A2), the consistency property is equivalent to require that for $l=2, \ldots, k$, and for all $v \in C^{l+1}$,

$$
\begin{align*}
\mathcal{E}_{A}(v)(t, x):= & \left|h^{A}\left(D^{-} v, D^{+} v\right)-H\left(v_{x}\right)+\frac{\Delta t}{2} H_{p}^{2}\left(v_{x}\right) v_{x x}\right| \\
& \leq C_{A, l}\left(\Delta t^{l}\left\|\partial_{t}^{l+1} v\right\|_{\infty}+\Delta x^{l}\left\|\partial_{x}^{l+1} v\right\|_{\infty}\right) . \tag{13}
\end{align*}
$$

Now, let us give some examples of high-order schemes satisfying (A1)-(A3) with $l=2$.
Example 2.4. As a first example let us consider the class of schemes obtained combining a highorder in space numerical hamiltonian $h_{*}^{A}$ and the second order Runge-Kutta SSP (or Heun scheme). To explain the simple procedure, let us consider the semidiscrete problem

$$
\left.u_{t}=h_{*}^{A}\left(D^{-} u(t, x), D^{+} u(t, x)\right)\right),
$$

where $h_{*}^{A}$, is a high-order in space numerical hamiltonian of second order,

$$
\begin{equation*}
h_{*}^{A}\left(D^{-} v_{j}^{n}, D^{+} v_{j}^{n}\right)=H\left(v_{x}\left(t^{n}, x_{j}\right)\right)+O\left(\Delta x^{2}\right), \tag{14}
\end{equation*}
$$

such as the simple second order central approximation

$$
\begin{equation*}
h_{*}^{A}\left(D^{-} u_{j}^{n}, D^{+} u_{j}^{n}\right)=H\left(\frac{D^{-} u_{j}^{n}+D^{+} u_{j}^{n}}{2}\right), \tag{15}
\end{equation*}
$$

then to obtain the same accuracy in time we discretize using the second order SSP Runge-Kutta scheme,

$$
\left\{\begin{array}{l}
u^{*}=u^{n}-\Delta t h_{*}^{A}\left(D^{-} u^{n}, D^{+} u^{n}\right)  \tag{16}\\
u^{n+1}=\frac{1}{2} u^{n}+\frac{1}{2} u^{*}-\frac{\Delta t}{2} h_{*}^{A}\left(D^{-} u^{*}, D^{+} u^{*}\right) .
\end{array}\right.
$$

The scheme can be written in differenced form in the sense of (A1)-(13) defining

$$
\begin{equation*}
h^{A}\left(D^{-} u^{n}, D^{+} u_{j}^{n}\right)=\frac{1}{2}\left[h_{*}^{A}\left(D^{-} u^{n}, D^{+} u^{n}\right)+h_{*}^{A}\left(D^{-} u^{*}, D^{+} u^{*}\right)\right] . \tag{17}
\end{equation*}
$$

To verify that the scheme is second order we can use the Taylor expansion to see that

$$
\begin{aligned}
h_{*}^{A}\left(D^{-} v_{j}^{*}, D^{+} v_{j}^{*}\right) & =H\left(v_{x}^{n}\left(x_{j}\right)-\Delta t \frac{d}{d x} H\left(v_{x}^{n}\left(x_{j}\right)\right)\right)+O\left(\Delta x^{2}\right) \\
& =H\left(v_{x}^{n}\left(x_{j}\right)\right)-\Delta t\left[H_{p}\left(v_{x}^{n}\left(x_{j}\right)\right) v_{x x}^{n}\left(x_{j}\right)\right] H_{p}\left(v_{x}^{n}\left(x_{j}\right)\right)+O\left(\Delta x^{2}\right),
\end{aligned}
$$

having exploited the relation $v^{*}=v^{n}-\Delta t\left[H\left(v_{x}^{n}\left(x_{j}\right)+O\left(\Delta x^{2}\right)\right]\right.$, the Lipschitz continuity of $H$ and having assumed a CFL condition $\lambda=\frac{\Delta t}{\Delta x}=$ const; whence, again using the consistency property (14)

$$
h^{A}\left(D^{-} v^{n}, D^{+} v^{n}\right)=H\left(v_{x}^{n}\left(x_{j}\right)\right)-\frac{\Delta t}{2} H_{p}^{2}\left(v_{x}^{n}\left(x_{j}\right)\right) v_{x x}^{n}\left(x_{j}\right)+O\left(\Delta x^{2}\right),
$$

as we wanted. Notice that through this procedure the stencil of the scheme (14) becomes doubled for $h^{A}$. Notice also that this procedure can be easily extended to the case of hamiltonian dependent on the space variable $x$.

Example 2.5. Then we propose a couple of numerical hamiltonians $h^{A}$ obtained discretizing directly the formula (12) or, equivalently, obtained from the same Lax-Wendroff schemes for conservation laws by the substitution $u_{j}^{n}=\frac{v_{j+1}^{n}-v_{j}^{n}}{\Delta x}$. The first is the original Lax-Wendroff scheme

$$
\begin{align*}
h^{A}\left(D^{-} u_{j}^{n}, D^{+} u_{j}^{n}\right)=\frac{1}{2}\{ & H\left(D^{+} u_{j}^{n}\right)+H\left(D^{-} u_{j}^{n}\right)+ \\
& \left.-\frac{\Delta t}{\Delta x} H_{p}\left(\frac{D^{-} u_{j}^{n}+D^{+} u_{j}^{n}}{2}\right)\left[H\left(D^{+} u_{j}^{n}\right)-H\left(D^{-} u_{j}^{n}\right)\right]\right\}, \tag{18}
\end{align*}
$$

and the second is its variation proposed by Richtmyer,

$$
\begin{equation*}
h^{A}\left(D^{-} u_{j}^{n}, D^{+} u_{j}^{n}\right)=H\left(\frac{D^{-} u_{j}^{n}+D^{+} u_{j}^{n}}{2}-\frac{\Delta t}{2 \Delta x}\left[H\left(D^{+} u_{j}^{n}\right)-H\left(D^{-} u_{j}^{n}\right)\right]\right) . \tag{19}
\end{equation*}
$$

Example 2.6. Following the approach of the Lax-Wendroff schemes and making use of the expansion (11), we can easily write higher order schemes, in both space and time, using very compact stencils. The idea is simply to discretize directly the above expansion using finite difference approximations of the right order. For example, if we want to write a fourth order Lax-Wendroff scheme using only five points, one of the possibilities is to define

$$
\begin{aligned}
H_{1}= & H\left(\frac{u_{j-2}-8 u_{j-1}+8 u_{j+1}-u_{j+2}}{12 \Delta x}\right) \\
H_{2}= & H_{p}^{2}\left(\frac{u_{j-2}-8 u_{j-1}+8 u_{j+1}-u_{j+2}}{12 \Delta x}\right)\left(\frac{-u_{j-2}+16 u_{j-1}-30 u_{j}+16 u_{j+1}-u_{j+2}}{12 \Delta x^{2}}\right), \\
H_{3}= & \frac{1}{2 \Delta x}\left[H_{p}^{3}\left(\frac{u_{j+2}-u_{j}}{2 \Delta x}\right)\left(\frac{u_{j+2}-2 u_{j+1}+u_{j}}{\Delta x^{2}}\right)-H_{p}^{3}\left(\frac{u_{j}-u_{j-2}}{2 \Delta x}\right)\left(\frac{u_{j}-2 u_{j-1}+u_{j-2}}{\Delta x^{2}}\right)\right], \\
H_{4}= & \frac{1}{\Delta x^{2}}\left[H_{p}^{4}\left(\frac{u_{j+2}-u_{j}}{2 \Delta x}\right)\left(\frac{u_{j+2}-2 u_{j+1}+u_{j}}{\Delta x^{2}}\right)-2 H_{p}^{4}\left(\frac{u_{j+1}-u_{j-1}}{2 \Delta x}\right)\left(\frac{u_{j+1}-2 u_{j}+u_{j-1}}{\Delta x^{2}}\right)\right. \\
& \left.\quad+H_{p}^{4}\left(\frac{u_{j}-u_{j-2}}{2 \Delta x}\right)\left(\frac{u_{j}-2 u_{j-1}+u_{j-2}}{\Delta x^{2}}\right)\right],
\end{aligned}
$$

and then compute

$$
\begin{equation*}
h^{A}\left(D^{-} u_{j}^{n}, D^{+} u_{j}^{n}\right)=H_{1}-\frac{\Delta t}{2}\left[H_{2}-\frac{\Delta t}{3}\left(H_{3}-\frac{\Delta t}{4} H_{4}\right)\right] \tag{20}
\end{equation*}
$$

It is straightforward to verify that, if the solution $v$ is regular enough, using Taylor expansion we have

- $H_{1}=H\left(v_{x}\right)+O\left(\Delta x^{4}\right)$,
- $H_{2}=H_{p}^{2}\left(v_{x}\right) v_{x x}+O\left(\Delta x^{4}\right)$,
- $H_{3}=\frac{\partial}{\partial x}\left(H_{p}^{3}\left(v_{x}\right) v_{x x}\right)+O\left(\Delta x^{2}\right)$,
- $H_{4}=\frac{\partial^{2}}{\partial x^{2}}\left(H_{p}^{4}\left(v_{x}\right) v_{x x}\right)+O\left(\Delta x^{2}\right)$,
and that the resulting scheme satisfies (A1)-(A3) with $l=4$. Notice that to obtain fourth order it would have been enough to have approximations of one order lower for $H_{2}$ and $H_{4}$, but thanks to the symmetry of the discretizations we can get higher orders without increasing the number of points in the stencil.


### 2.2 Filter function

In order to couple the schemes and their properties, we need to define a function $F$, called filter function $F$, such that
(F1) $F(x) \approx x$ for $|x| \leq 1$,
(F2) $F(x)=0$ for $|x|>1$,
which implies that

- If $\left|S^{A}-S^{M}\right| \leq \Delta t \varepsilon^{n}$ and $\phi_{j}^{n}=1 \Rightarrow S^{A F} \approx S^{A}$
- If $\left|S^{A}-S^{M}\right|>\Delta t \varepsilon^{n}$ or $\phi_{j}^{n}=0 \Rightarrow S^{A F}=S^{M}$.

It is clear that, with just these two requirements, we are left with several possible choices for $F$. In the following, we present some examples of filter functions satisfying the previous relations, which differ especially for regularity properties. We number the functions in order to be clearer in Figure 1.

Example 2.7. As a first example let us present the filter function we use in our numerical tests, defined in [5] as

$$
F_{1}(x)= \begin{cases}x & \text { if }|x| \leq 1  \tag{21}\\ 0 & \text { otherwise }\end{cases}
$$

which is clearly discontinuous at $x=-1,1$ and satisfies trivially the properties (F1)-(F2).
Example 2.8. As a second possibility we propose the family of regular filter functions given by the formula

$$
F(x)=x \exp \left(-c(|x|-a)^{b}\right),
$$

for appropriate choices of the parameters $a, b$ and $c$.
Remark 2.9. We give some hints on how to chose the parameters. We notice that

- a controls the amplitude of the transition phase around 1 and -1 ;
- b controls the slope of the transition phase;
- c can be used to make the exponent approach 0 faster when $x \approx 1,-1$.

In particular, in Figure 1 we represent two choices for the parameters,

$$
\begin{equation*}
F_{2}(x)=x \exp \left(-4(|x|-0.25)^{20}\right) \quad(a=0.25, b=20, c=4) \tag{22}
\end{equation*}
$$

and a variant graphically more similar to $F_{1}$,

$$
\begin{equation*}
F_{3}(x)=x \exp \left(-\frac{(|x|-0.01)^{50}}{100}\right) \quad(a=0.01, b=50, c=0.01) . \tag{23}
\end{equation*}
$$

These functions are very regular $\left(F \in C^{\infty}\right)$ and developing with Taylor we can see that they satisfy (F1)-(F2).
Example 2.10. As last examples let us consider some functions which satisfy (F1)-(F2) and are continuous, but are not necessarily derivable. First, let us consider the family of functions

$$
F(x)= \begin{cases}x \exp \left(-\frac{a}{b-|x|}\right) & \text { if }|x| \leq b  \tag{24}\\ 0 & \text { otherwise },\end{cases}
$$

varying the parameters $a$ and $b$. We propose the choice

$$
F_{4}(x)= \begin{cases}x \exp \left(-\frac{0.001}{1.05-|x|}\right) & \text { if }|x| \leq 1.05  \tag{25}\\ 0 & \text { otherwise }\end{cases}
$$

where we chose the value $b=1.05$ in order to make the function approach better the value 1 for $x=-1,1$. Finally, we recall also the filter defined in [18] as

$$
F_{5}(x)= \begin{cases}x & |x| \leq 1  \tag{26}\\ 0 & |x| \geq 2 \\ -x+2 & 1 \leq x \leq 2 \\ -x-2 & -2 \leq x \leq-1\end{cases}
$$



Figure 1: Possible choices for the filter function $F$.

After extensive computations, we noticed that the results obtained with our adaptive filtered (AF) scheme are not sensitive with respect to changes in regularity of the filter function, even with very large transition phases. That is probably because, as we will see in the next section, the parameter $\varepsilon^{n}$ is designed to obtain the property (F1) whenever possible, then in regions of regularity of the solution the argument of $F$ lies most probably in $[-1,1]$, where all the filter functions are practically the same. Some major differences, instead, can be seen in the results obtained with the basic filtered scheme, for which the threshold $\varepsilon$ is fixed at the beginning.

### 2.3 Tuning of the parameter $\varepsilon^{n}$

The last step is to show how to compute the switching parameter $\varepsilon^{n}$, which is the real core of the adaptivity of our scheme. Then, if we want the scheme (3) to switch to the high-order scheme when some regularity is detected, we have to choose $\varepsilon^{n}$ such that

$$
\begin{equation*}
\left|\frac{S^{A}\left(v^{n}\right)_{j}-S^{M}\left(v^{n}\right)_{j}}{\varepsilon^{n} \Delta t}\right|=\left|\frac{h^{A}(\cdot)-h^{M}(\cdot)}{\varepsilon^{n}}\right| \leq 1, \quad \text { for }(\Delta t, \Delta x) \rightarrow 0 \tag{27}
\end{equation*}
$$

in the region of regularity at time $t_{n}$, that is

$$
\begin{equation*}
\mathcal{R}^{n}=\left\{x_{j}: \phi_{j}^{n}=1\right\} . \tag{28}
\end{equation*}
$$

For the moment, to simplify the presentation we assume the existence of a function $\phi$ such that

$$
\phi_{j}^{n}= \begin{cases}1 & \text { if the solution } u^{n} \text { is regular in } I_{j}  \tag{29}\\ 0 & \text { if } I_{j} \text { contains a point of singularity },\end{cases}
$$

referring to the next section for some examples of practical computation of the function $\phi$.
Therefore, assuming $v$ sufficiently smooth, computing directly by Taylor expansions, we have for the monotone scheme

$$
h^{M}\left(D^{-} v_{j}^{n}, D^{+} v_{j}^{n}\right)=H\left(v_{x}^{n}\left(x_{j}\right)\right)+\frac{\Delta x}{2} v_{x x}^{n}\left(x_{j}\right)\left(\partial_{p_{+}} h_{j}^{M}-\partial_{p_{-}} h_{j}^{M}\right)+O\left(\Delta x^{2}\right),
$$

where we used the relation

$$
\left.D^{ \pm} v_{j}^{n}=v_{x}^{n}\left(x_{j}\right)\right) \pm \frac{\Delta x}{2} v_{x x}^{n}\left(x_{j}\right)+O\left(\Delta x^{2}\right)
$$

while for the high-order scheme, by the consistency property,

$$
h^{A}\left(D^{-} v_{j}^{n}, D^{+} v_{j}^{n}\right)=H\left(v_{x}^{n}\left(x_{j}\right)\right)-\frac{\Delta t}{2} H_{p}^{2}\left(v_{x}^{n}\right) v_{x x}+O\left(\Delta t^{2}\right)+O\left(\Delta x^{2}\right) .
$$

Whence, from (27) we obtain

$$
\begin{equation*}
\varepsilon^{n} \geq\left|\frac{\Delta x}{2} v_{x x}^{n}\left(\partial_{p_{+}} h_{j}^{M}-\partial_{p_{-}} h_{j}^{M}+\lambda H_{p}^{2}\left(v_{x}^{n}\right)\right)+O\left(\Delta t^{2}\right)+O\left(\Delta x^{2}\right)\right| . \tag{30}
\end{equation*}
$$

Finally, we use a numerical approximation of the lower bound on the right hand side of the previous inequality to obtain the following formula for $\varepsilon^{n}$,

$$
\begin{align*}
& \varepsilon^{n}=\max _{x_{j} \in \mathcal{R}^{n}} K \mid H\left(D u_{j}^{n}\right)-H\left(D u_{j}^{n}-\lambda\left[H\left(D^{+} u_{j}^{n}\right)-H\left(D^{-} u_{j}^{n}\right)\right]\right) \\
&+ {\left[h^{M}\left(D u_{j}^{n}, D^{+} u_{j}^{n}\right)-h^{M}\left(D u_{j}^{n}, D^{-} u_{j}^{n}\right)\right] } \\
&-\left[h^{M}\left(D^{+} u_{j}^{n}, D u_{j}^{n}\right)-h^{M}\left(D^{-} u_{j}^{n}, D u_{j}^{n}\right)\right] \mid \tag{31}
\end{align*}
$$

with $K>\frac{1}{2}, \lambda:=\frac{\Delta t}{\Delta x}$ and $D u_{j}^{n}:=\frac{u_{j+1}^{n}-u_{j-1}^{n}}{2 \Delta x}$. Notice that if we assume enough regularity on the solution $v$, then (31) gives a second order approximation of the right hand side of (30) multiplied by $2 K$.

## 3 Smoothness indicator function

In the previous section we assumed the existence of a smoothness indicator function $\phi$, in the sense that

$$
\phi_{j}^{n}=\phi\left(\omega_{j}^{n}\right):= \begin{cases}1 & \text { if the solution } u^{n} \text { is regular in } I_{j},  \tag{32}\\ 0 & \text { if } I_{j} \text { contains a point of singularity },\end{cases}
$$

where $I_{j}=\left(x_{j-1}, x_{j+1}\right)$ and $\omega_{j}^{n}$ is the smoothness indicator at the node $x_{j}$ depending on the values of the approximate solution $u^{n}$. The aim of this section is precisely to show a simple construction of a function satisfying (29) which makes use of smoothness indicators widely known in literature. Moreover, in the process we review the theory of the smoothness indicators of [15], defined for the construction of the WENO schemes for (1],

$$
\begin{equation*}
\beta_{k}=\beta_{k}\left(u^{n}\right)_{j}:=\sum_{l=2}^{r} \int_{x_{j-1}}^{x_{j}} \Delta x^{2 l-3}\left(P_{k}^{(l)}(x)\right)^{2} d x, \tag{33}
\end{equation*}
$$

for $k=0, \ldots, r-1$, where $P_{k}$ is the Lagrange polynomial of degree $r$ interpolating the values of $u^{n}$ on the stencil $\mathcal{S}_{j+k}=\left\{x_{j+k-r}, \ldots, x_{j+k}\right\}$.

Then, before proceeding with the construction of $\phi$, let us state a fundamental result on the behavior of the indicators (33).

Proposition 3.1. Assume $f \in C^{r+1}\left(\Omega \backslash\left\{x_{s}\right\}\right)$, with $\Omega$ a neighborhood of $x_{s}$, and $f^{\prime}\left(x_{s}^{-}\right) \neq f^{\prime}\left(x_{s}^{+}\right)$. Then, for $k=0, \ldots, r-1$ and $j \in \mathbb{Z}$, the followings are true:
i) If $x_{s} \in \Omega \backslash \stackrel{\circ}{\mathcal{S}}_{j+k} \Rightarrow \beta(f)_{k}=O\left(\Delta x^{2}\right)$,
ii) If $x_{s} \in \stackrel{\circ}{\mathcal{S}}_{j+k} \Rightarrow \beta(f)_{k}=O(1)$,
where $\mathcal{S}_{j+k}=\left\{x_{j-r+k}, \ldots, x_{j+k}\right\}$ and $\stackrel{\circ}{\mathcal{S}}_{j+k}=\left(x_{j-r+k}, x_{j+k}\right)$.
We skip the proof, which is rather technical, but the interested reader can find it in the Appendix.
Remark 3.2. Notice that we could avoid the restrictions on $f$ in the points of regularity by adding a small quantity $\sigma_{h}=\sigma \Delta x^{2}$, for some constant $\sigma>0$, to the indicators $\beta_{k}$ and consider instead

$$
\begin{equation*}
\widetilde{\beta}_{k}:=\beta_{k}+\sigma_{h}, \tag{34}
\end{equation*}
$$

as it has been done in [3]. We will use this assumption in the sequel, choosing $\sigma=1$.
Our aim is to identify the points (or the intervals) in which the approximate solution $u^{n}$ presents a singularity in the first derivative. To be precise, here with $u^{n}$ we mean any continuous function with nodal values $u_{j}^{n}, j \in \mathbb{Z}$. Let us focus the attention on a point $x_{j}$ of the grid and consider the simplest case of $r=2$, which is enough for our purpose. Let us consider separately the intervals $\left(x_{j-1}, x_{j}\right]$ and $\left[x_{j}, x_{j+1}\right)$ defining

$$
\begin{equation*}
\beta_{k}^{-}=\Delta x \int_{x_{j-1}}^{x_{j}}\left(P_{k}^{\prime \prime}(x)\right)^{2} d x \tag{35}
\end{equation*}
$$

for $k=0,1$, where $P_{0}, P_{1}$ are the polynomials interpolating the solution, respectively, on the stencils $\left\{x_{j-2}, x_{j-1}, x_{j}\right\}$ e $\left\{x_{j-1}, x_{j}, x_{j+1}\right\}$; and symmetrically

$$
\begin{equation*}
\beta_{k}^{+}=\Delta x \int_{x_{j}}^{x_{j+1}}\left(P_{k}^{\prime \prime}(x)\right)^{2} d x \tag{36}
\end{equation*}
$$

for $k=0,1$, where now $P_{0}, P_{1}$ are the interpolating polynomials on $\left\{x_{j-1}, x_{j}, x_{j+1}\right\}$ and $\left\{x_{j}, x_{j+1}, x_{j+2}\right\}$. From the definition it is clear that $\left(\beta^{+}\right)_{j}=\left(\beta^{-}\right)_{j+1}$ so we have to compute the quantities just once. Then, we define as in [15]

$$
\begin{equation*}
\alpha_{k}^{ \pm}=\frac{1}{\left(\beta_{k}^{ \pm}+\sigma_{h}\right)^{2}}, \tag{37}
\end{equation*}
$$

with $\sigma_{h}=\sigma \Delta x^{2}$ the parameter we introduced in Remark 3.2, and focus on the information given by the interpolating polynomial on $\left\{x_{j-1}, x_{j}, x_{j+1}\right\}$ defining

$$
\omega_{+}=\frac{\alpha_{0}^{+}}{\alpha_{0}^{+}+\alpha_{1}^{+}},
$$

to inspect the regularity on $\left[x_{j}, x_{j+1}\right)$ and in the same way for $\left(x_{j-1}, x_{j}\right]$,

$$
\omega_{-}=\frac{\alpha_{1}^{-}}{\alpha_{0}^{-}+\alpha_{1}^{-}} .
$$

By Proposition 3.1 and Remark 3.2 we know that $\widetilde{\beta}_{k}=O\left(\Delta x^{2}\right)$ if there is no singularity in the stencil, and $\widetilde{\beta}_{k}=O(1)$ otherwise, so in presence of a singularity we can only fall in one of the following cases:

- If $x_{j-2}<x_{s} \leq x_{j-1}$, then $\widetilde{\beta}_{0}^{-}=O(1), \widetilde{\beta}_{1}^{-}=\widetilde{\beta}_{0}^{+}=O\left(\Delta x^{2}\right), \widetilde{\beta}_{1}^{+}=O\left(\Delta x^{2}\right)$,
- If $x_{j-1}<x_{s}<x_{j}$, then $\widetilde{\beta}_{0}^{-}=O(1), \widetilde{\beta}_{1}^{-}=\widetilde{\beta}_{0}^{+}=O(1), \widetilde{\beta}_{1}^{+}=O\left(\Delta x^{2}\right)$,
- If $x_{s}=x_{j}$, then $\widetilde{\beta}_{0}^{-}=O\left(\Delta x^{2}\right), \widetilde{\beta}_{1}^{-}=\widetilde{\beta}_{0}^{+}=O(1), \widetilde{\beta}_{1}^{+}=O\left(\Delta x^{2}\right)$,
- If $x_{j}<x_{s}<x_{j+1}$, then $\widetilde{\beta}_{0}^{-}=O\left(\Delta x^{2}\right), \widetilde{\beta}_{1}^{-}=\widetilde{\beta}_{0}^{+}=O(1), \widetilde{\beta}_{1}^{+}=O(1)$,
- If $x_{j+1} \leq x_{s}<x_{j+2}$, then $\widetilde{\beta}_{0}^{-}=O\left(\Delta x^{2}\right), \widetilde{\beta}_{1}^{-}=\widetilde{\beta}_{0}^{+}=O\left(\Delta x^{2}\right), \widetilde{\beta}_{1}^{+}=O(1)$,
with $x_{s}$ point of singularity. Now, we can compute

$$
\begin{equation*}
\frac{\alpha_{1}^{ \pm}-\alpha_{0}^{ \pm}}{\alpha_{0}^{ \pm}}=\frac{\left(\beta_{0}^{ \pm}+\sigma_{h}\right)^{2}-\left(\beta_{1}^{ \pm}+\sigma_{h}\right)^{2}}{\left(\beta_{1}^{ \pm}+\sigma_{h}\right)^{2}}=\left(\frac{\beta_{0}^{ \pm}-\beta_{1}^{ \pm}}{\beta_{1}^{ \pm}+\sigma_{h}}\right)\left(\frac{\beta_{0}^{ \pm}+\beta_{1}^{ \pm}+2 \sigma_{h}}{\beta_{1}^{ \pm}+\sigma_{h}}\right), \tag{38}
\end{equation*}
$$

which, noticing that, if the function is smooth in both stencils of $\beta_{0}^{ \pm}$and $\beta_{1}^{ \pm}$, we have

$$
\begin{align*}
& \frac{\beta_{0}^{ \pm}-\beta_{1}^{ \pm}}{\beta_{1}^{ \pm}+\sigma_{h}}=-2 h \frac{f_{j}^{\prime \prime} f_{j}^{\prime \prime \prime}}{\left(f^{\prime \prime}\right)^{2}+\sigma}+O\left(h^{2}\right)=O(h)  \tag{39}\\
& \frac{\beta_{0}^{ \pm}+\beta_{1}^{ \pm}+2 \sigma_{h}}{\beta_{1}^{ \pm}+\sigma_{h}}=2+O(h)=O(1),
\end{align*}
$$

leads to

$$
\begin{equation*}
\alpha_{1}^{ \pm}=\alpha_{0}^{ \pm}(1+O(\Delta x)) . \tag{40}
\end{equation*}
$$

Whence we can deduce that if the solution is regular enough in both stencils

$$
\begin{equation*}
\omega_{ \pm}=\frac{\alpha_{k}^{ \pm}}{\alpha_{0}^{ \pm}+\alpha_{1}^{ \pm}}=\frac{1}{2}+O(\Delta x) \tag{41}
\end{equation*}
$$

with $k=0$ for the superscript " + " and $k=1$ for " - ". On the other hand, if there is a singularity in at least one of the stencils

$$
\alpha_{k}^{ \pm}= \begin{cases}O(1) & \text { if } f \text { is not smooth in } \mathcal{S}_{j+k}  \tag{42}\\ O\left(\Delta x^{-4}\right) & \text { if } f \text { is smooth in } \mathcal{S}_{j+k}\end{cases}
$$

then it is easy to verify that the behavior of our $\omega_{ \pm}$falls in the following cases:

- If $x_{j-2}<x_{s} \leq x_{j-1}$, then $\omega_{-}=1+O\left(\Delta x^{4}\right), \omega_{+}=1 / 2+O(\Delta x)$
- If $x_{j-1}<x_{s}<x_{j}$, then $\omega_{-}=O(1), \omega_{+}=O\left(\Delta x^{4}\right)$
- If $x_{s}=x_{j}$, then $\omega_{-}=O\left(\Delta x^{4}\right), \omega_{+}=O\left(\Delta x^{4}\right)$
- If $x_{j}<x_{s}<x_{j+1}$, then $\omega_{-}=O\left(\Delta x^{4}\right), \omega_{+}=O(1)$
- If $x_{j+1} \leq x_{s}<x_{j+2}$, then $\omega_{-}=1 / 2+O(\Delta x), \omega_{+}=1+O\left(\Delta x^{4}\right)$,
where with $\omega_{ \pm}=O(1)$ we mean a number dependent on the jump of the derivative. To be precise, here the $O(1)$ comes from the fact that, by Proposition 3.1, we would not have the property (42) in presence of a singularity, so using again the expansion (38) we have to notice that

$$
\frac{\beta_{0}^{ \pm}-\beta_{1}^{ \pm}}{\beta_{1}^{ \pm}+\sigma_{h}}=O(1)
$$

Now, defining $\omega_{j}=\min \left\{\omega_{-}, \omega_{+}\right\}$we can rewrite

$$
\omega_{j}= \begin{cases}O\left(\Delta x^{4}\right) & \text { if } x_{j-1}<x_{s}<x_{j+1} \\ \frac{1}{2}+O(\Delta x) & \text { otherwise } .\end{cases}
$$

Unfortunately, we noticed through numerical tests that the $O(\Delta x)$ term in regular regions may produce heavy oscillations around the optimal value $\bar{\omega}=1 / 2$. To increase the accuracy, we can use higher order smoothness indicator $(r>2)$, but we would need a bigger reconstruction stencil, or we can use the mappings defined in [12],

$$
\begin{equation*}
g(\omega)=\frac{\omega\left(\bar{\omega}+\bar{\omega}^{2}-3 \bar{\omega} \omega+\omega^{2}\right)}{\bar{\omega}^{2}+\omega(1-2 \bar{\omega})}, \quad \bar{\omega} \in(0,1) \tag{43}
\end{equation*}
$$

which have the properties that $g(0)=0, g(1)=1, g(\bar{\omega})=\bar{\omega}, g^{\prime}(\bar{\omega})=0$ and $g^{\prime \prime}(\bar{\omega})=0$. Then, we define

$$
\begin{aligned}
\omega_{ \pm}^{*} & =g\left(\omega_{ \pm}\right) \\
& =g(\bar{\omega})+g^{\prime}(\bar{\omega})\left(\omega_{ \pm}-\bar{\omega}\right)+\frac{g^{\prime \prime}(\bar{\omega})}{2}\left(\omega_{ \pm}-\bar{\omega}\right)^{2}+\frac{g^{\prime \prime \prime}(\bar{\omega})}{6}\left(\omega_{ \pm}-\bar{\omega}\right)^{3}+O\left(\Delta x^{4}\right) \\
& =\bar{\omega}+\frac{\left(\omega_{ \pm}-\bar{\omega}\right)^{3}}{\bar{\omega}-\bar{\omega}^{3}}+O\left(\Delta x^{4}\right) \\
& =\bar{\omega}+O\left(\Delta x^{3}\right) .
\end{aligned}
$$

Remark 3.3. Notice that with respect to the definition in [12] we avoided the second weighting which seems unnecessary in our case. More explicitly, the mapping we use is

$$
\begin{equation*}
g(\omega)=4 \omega\left(\frac{3}{4}-\frac{3}{2} \omega+\omega^{2}\right) . \tag{44}
\end{equation*}
$$

Finally, what is left is to define the function $\phi$ such that $\phi=1$ if $\omega$ is close to $\frac{1}{2}$ and $\phi=0$, otherwise. Notice that in the latter are included both cases in which the function has a singularity in the first derivative $\left(\omega=O\left(\Delta x^{4}\right)\right)$ and when the second derivative is discontinuous ( $\omega=O(1)$ ). The simplest choice is to take

$$
\begin{equation*}
\phi(\omega)=\chi_{\{\omega \geq M\}}, \tag{45}
\end{equation*}
$$

with $M<\frac{1}{2}$, a number possibly dependent on $\Delta x$.
Remark 3.4. Notice that to construct the function $\phi$ using the indicators (33) with $r=2$ we need only five points to inspect the regularity in $I_{j}$.

Next, we show that if we make a particular choice for $M$ we are able to prove the following result, which can be seen as an "inverse" of Proposition 3.1 for numerical solutions and (probably)
gives a useful tool for the analysis of the next section. Unfortunately, at the moment, this result is valid only for indicators $\omega$ using the standard construction for $r=2$, without the possibility to introduce any of the presented modifications, or higher order indicators. Moreover, as it will be briefly discussed in Remark 4.4, it introduces some limitations in the applicability even when using the standard indicators, testifying the necessity of some improvements in the argument used.

Before proceeding, let us remind that we are working with structured grids, then if we consider a one-parameter family of grid values $\left\{u_{j}(\Delta x)\right\}_{j \in J(\Delta x)}$, as $\Delta x$ goes to 0 , the indexed family of sets of indices $J(\Delta x)$ is expanding, in the sense that if $\Delta x_{2}<\Delta x_{1}$, then $J\left(\Delta x_{1}\right) \subset J\left(\Delta x_{2}\right)$, where $J(\Delta x) \subseteq \mathbb{Z}$, for all $\Delta x>0$. Moreover, we define $I_{s}(\Delta x)$ as the set of indices $j$ such that $\phi_{j}=0$ and assume, for simplicity, $\left|I_{s}(\Delta x)\right|<\infty$ and $I_{s}(\Delta x) \equiv I_{s}$, for $\Delta x>0$.

Lemma 3.1. Let $\omega$ be computed using (35)-(36) and $\phi$ be defined by (45) with $M(\Delta x)=\frac{1}{2}-C \Delta x$, for some constant $C$ such that $0<M(\Delta x)<\frac{1}{2}$. Consider a one-parameter family of sequences $\left\{u_{j}(\Delta x)\right\}_{j \in J(\Delta x)}$, and a partition $\left\{R_{i}\right\}_{i=0, \ldots,\left|I_{s}\right|}$ of the regularity set $\mathcal{R}=\left\{j \in \mathbb{Z}: \phi_{j}=1\right\}=\bigcup_{i} R_{i}$, and $\mathcal{R}=\mathbb{Z}$ if $I_{s}=\emptyset$. Then, if for all $i=0, \ldots,\left|I_{s}\right|$, there exists $j_{i} \in R_{i}$, such that $\left|D^{2} u_{j_{i}}(\Delta x)\right|<\infty$, we have that

$$
\begin{equation*}
\left|D^{2} u_{j}(\Delta x)\right|=\frac{\left|u_{j+1}(\Delta x)-2 u_{j}(\Delta x)+u_{j-1}(\Delta x)\right|}{\Delta x^{2}} \leq B, \quad \forall j \in \mathcal{R} \tag{46}
\end{equation*}
$$

for a constant $B$ independent of $\Delta x$.
Proof. It is clear that, since $\left|I_{s}\right|<\infty$ by hypothesis, it is enough to prove the assertion just for one $i \in I_{s}$, or more simply in the regular case $\mathcal{R}=\mathbb{Z}$. Let us assume then that the sequence is "regular" and, without loss of generality, that there exists $j_{\min } \in \mathcal{R}$ such that $D^{2} u_{j_{\min }}(\Delta x)=0$, which happens for example when the sequence has compact support. In the following, we drop the dependence on $\Delta x$ for clarity of presentation, since it should not cause confusion.

Since we are considering the case of $\omega$ computed by non-mapped indicators, we have that, by definition of $\phi$ and $\omega$, if $\phi_{j}=1$ then at least one $\omega_{ \pm}>M$. Without loss of generality, let us assume $\omega_{+}>M$, with the other case being symmetrical. Then, by definition,

$$
\omega_{+}=\frac{\left(\beta_{1}^{+}+\sigma_{h}\right)^{2}}{\left(\beta_{1}^{+}+\sigma_{h}\right)^{2}+\left(\beta_{0}^{+}+\sigma_{h}\right)^{2}}>M
$$

which leads by simple computations to

$$
\left|\beta_{0}^{+}\right|<\sqrt{\frac{1-M}{M}}\left|\beta_{1}^{+}\right|+\left(\sqrt{\frac{1-M}{M}}-1\right) \sigma_{h}
$$

and dividing by $\Delta x^{2}$, we get

$$
\left|D^{2} u_{j}\right|^{2}<\sqrt{\frac{1-M}{M}}\left|D^{2} u_{j+1}\right|^{2}+\left(\sqrt{\frac{1-M}{M}}-1\right)
$$

where we used the definition of $\sigma_{h}=\sigma \Delta x^{2}$, with $\sigma=1$ for simplicity. Whence, iterating till $L_{j} \in \mathbb{N}$
such that $j+L_{j}=j_{\text {min }}$, we have

$$
\begin{aligned}
\left|D^{2} u_{j}\right|^{2}<\ldots & <\left(\frac{1-M}{M}\right)^{\frac{L_{j}}{2}}\left|D^{2} u_{j_{\min }}\right|^{2}+\left(\sqrt{\frac{1-M}{M}}-1\right) \sum_{k=1}^{L_{j}}\left(\frac{1-M}{M}\right)^{\frac{k-1}{2}} \\
& =\left(\sqrt{\frac{1-M}{M}}-1\right) \sum_{k=0}^{L_{j}-1}\left(\frac{1-M}{M}\right)^{\frac{k}{2}} \\
& =\left(\sqrt{\frac{1-M}{M}}-1\right) \frac{1-\left(\frac{1-M}{M}\right)^{\frac{L_{j}}{2}}}{1-\sqrt{\frac{1-M}{M}}}=\left(\frac{1-M}{M}\right)^{\frac{L_{j}}{2}}-1 .
\end{aligned}
$$

Now, if notice that forall $j \in \mathcal{R}$ we can find a constant $L>0$ independent on $j$ such that $L_{j} \leq \frac{L}{\Delta x}$ and recall the hypothesis on $M=\frac{1}{2}-C \Delta x$, we can conclude

$$
\begin{aligned}
\left|D^{2} u_{j}\right|^{2} & \leq\left(\frac{1-M}{M}\right)^{\frac{L_{j}}{2}}-1 \leq\left(\frac{1}{M}-1\right)^{\frac{L}{2 \Delta x}}-1 \\
& =\left(\frac{2}{1-2 C \Delta x}-1\right)^{\frac{L}{2 \Delta x}}-1 \\
& \approx(1+4 C \Delta x)^{\frac{L}{2 \Delta x}}-1 \xrightarrow{\Delta x \rightarrow 0} e^{2 L C}-1,
\end{aligned}
$$

by the well known notable limit. Then, the statement follows taking $B:=\sqrt{e^{2 L C}-1}$.
Remark 3.5. Notice that the previous lemma strongly relies on the fact that $\omega$ is computed using (35)-(36) without introducing the mappings (44). In fact, if we were to use (44), we could develop the algebra until the inequality

$$
\left|D^{2} u_{j}\right|^{2} \leq\left(\frac{1}{g^{-1}(M)}-1\right)^{\frac{L}{2 \Delta x}}-1,
$$

but, by definition, $g^{-1}$ can not be expanded in Taylor series around the point $\frac{1}{2}$, whence we could not use the notable limit to conclude.

On the other hand, through extensive numerical simulations on various critical situations, we could acknowledge that a weaker result seems to hold also for more general indicators. More precisely, we collected numerical evidence that, fixed $j \in \mathbb{Z}$, if the sequence of the second order increments

$$
D^{2} u_{j}(\Delta x)=\frac{u_{j+1}(\Delta x)-2 u_{j}(\Delta x)+u_{j-1}(\Delta x)}{\Delta x^{2}},
$$

presents some kind of discontinuity, then we have

$$
\omega_{j}=O(1) .
$$

Consequently, choosing $M(\Delta x)=\frac{1}{2}-C \Delta x$, for a constant $C$ such that $0<M(\Delta x)<\frac{1}{2}$, or even more simply, $M(\Delta x) \rightarrow 0$ as $\Delta x \rightarrow 0$, we could infer

$$
j \notin \mathcal{R}, \quad \text { for } \Delta x \rightarrow 0
$$

Notice that this property can be proven almost directly in the case the discontinuity of the second order increment is "caused" by that of the first order finite difference, in the sense

$$
D^{ \pm} u_{j}(\Delta x):= \pm \frac{u_{j \pm 1}(\Delta x)-u_{j}(\Delta x)}{\Delta x} \rightarrow u_{x}^{ \pm}
$$

with $u_{x}^{+} \neq u_{x}^{-}$. On the contrary, if the sequence of $D^{ \pm} u_{j}(\Delta x)$ is "regular", the detection of a discontinuity in $D^{2} u_{j}(\Delta x)$ is more involved. It is noteworthy to point out that we are interested mainly in detecting unbounded second order increments. Unfortunately, without any further assumption on the sequences $\left\{u_{j}(\Delta x)\right\}_{j \in J(\Delta x)}$, such a result would not suffice to infer that $\left|D^{2} u_{j}\right|<B$ if $j \in \mathcal{R}$, for some $B>0$, since we could not secure the boundedness of second order increments at points in some neighborhood of a cell (point) at which the sequence is "regular" but has unbounded second order increment.

Therefore, we are forced to add a "technical" assumption in order to justify the proof of Proposition 4.1. More precisely, when using the alternative constructions for $\omega$ (using the mapping (44)), we define the region of regularity $\mathcal{R}$ detected by the function $\widetilde{\phi}$ as the set

$$
\mathcal{R}=\left\{j \in \mathbb{Z}: \widetilde{\phi}\left(\omega_{j}\right)=1\right\}, \quad \text { with } \quad \widetilde{\phi}_{j}= \begin{cases}1 & \text { if } \phi\left(\omega_{j}\right)=1 \text { and }\left|D u_{j}^{2}\right|<B  \tag{47}\\ 0 & \text { otherwise },\end{cases}
$$

for some constant $B \gg 0$. Notice that with this definition, which, we recall, is needed only for theoretical reasons, it is not necessary to require $M(\Delta x) \rightarrow 0$, then we can simply choose a constant $M>0$ small enough (e.g. $M=0.1$ ), as we will do in the numerical tests of Section 5 .

## 4 Convergence result

We are now able to present our main result, but before doing so let us state a useful proposition about the numerical solution and the parameter $\varepsilon^{n}$.

Proposition 4.1. Let $u^{n}$ be the solution obtained by the scheme (3)-(31) and assume that $v_{0}$ and $H$ are Lipschitz continuous functions. Assume also that $\mathcal{R}^{n}$ is defined by (28) or (47), with $\phi$ given by (45), and that $\lambda=\Delta t / \Delta x=$ constant. Then, $\varepsilon^{n}$ is well defined and $u^{n}$ satisfies, for any $i$ and $j$, the discrete Lipschitz estimate

$$
\frac{\left|u_{i}^{n}-u_{j}^{n}\right|}{\Delta x} \leq L
$$

for some constant $L>0$, for $0 \leq n \leq T / \Delta t$. Moreover, there exists a constant $C>0$ such that

$$
\varepsilon^{n} \leq C \Delta x .
$$

Proof. Before proceeding with the proof let us notice that, if $u^{n}$ satisfies 4.1) for a constant $L_{n}>0$, calling for brevity

$$
D^{*} u_{j}:=D u_{j}^{n}-\lambda\left[H\left(D^{+} u_{j}^{n}\right)-H\left(D^{-} u_{j}^{n}\right)\right]
$$

we have that

$$
\begin{align*}
& \varepsilon^{n}=\max _{x_{j} \in \mathcal{R}^{n}} K \mid H\left(D u_{j}^{n}\right)-H\left(D^{*} u_{j}\right)+\left[h^{M}\left(D u_{j}^{n}, D^{+} u_{j}^{n}\right)-h^{M}\left(D u_{j}^{n}, D^{-} u_{j}^{n}\right)\right] \\
& -\left[h^{M}\left(D^{+} u_{j}^{n}, D u_{j}^{n}\right)-h^{M}\left(D^{-} u_{j}^{n}, D u_{j}^{n}\right)\right] \mid \\
& =\max _{x_{j} \in \mathcal{R}^{n}} K \left\lvert\,\left[\Delta t\left(\frac{H\left(D u_{j}^{n}\right)-H\left(D^{*} u_{j}\right)}{D u_{j}^{n}-D^{*} u_{j}}\right)\left(\frac{H\left(D^{+} u_{j}^{n}\right)-H\left(D^{-} u_{j}^{n}\right)}{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}\right)\right.\right. \\
& +\Delta x\left(\frac{h^{M}\left(D u_{j}^{n}, D^{+} u_{j}^{n}\right)-h^{M}\left(D u_{j}^{n}, D^{-} u_{j}^{n}\right)}{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}\right) \\
& \left.-\Delta x\left(\frac{h^{M}\left(D^{+} u_{j}^{n}, D u_{j}^{n}\right)-h^{M}\left(D^{-} u_{j}^{n}, D u_{j}^{n}\right)}{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}\right)\right] \left.\left(\frac{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}{\Delta x}\right) \right\rvert\, \\
& =\max _{x_{j} \in \mathcal{R}^{n}} K \left\lvert\,\left[\Delta t\left(\frac{H\left(D u_{j}^{n}\right)-H\left(D^{*} u_{j}\right)}{D u_{j}^{n}-D^{*} u_{j}}\right)\left(\frac{H\left(D^{+} u_{j}^{n}\right)-H\left(D^{-} u_{j}^{n}\right)}{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}\right)\right.\right. \\
& +\Delta x\left(\frac{h^{M}\left(D u_{j}^{n}, D^{+} u_{j}^{n}\right)-h^{M}\left(D u_{j}^{n}, D^{-} u_{j}^{n}\right)}{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}\right) \\
& \left.-\Delta x\left(\frac{h^{M}\left(D^{+} u_{j}^{n}, D u_{j}^{n}\right)-h^{M}\left(D^{-} u_{j}^{n}, D u_{j}^{n}\right)}{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}\right)\right] \left.\frac{\sqrt{\beta_{0}^{+}\left[u_{j}^{n}\right]}}{\Delta x} \right\rvert\, \\
& \leq K\left|\left(\Delta t L_{H 2} L_{H}+2 \Delta x L_{h^{M}}\right) B\right| \\
& =K B\left(\lambda L_{H 2} L_{H}+2 L_{h^{M}}\right) \Delta x, \tag{48}
\end{align*}
$$

where $L_{H}$ and $L_{H 2}$ are the local Lipschitz constant of $H$ on $\left[-L_{n}, L_{n}\right]$ and $\left[-2 L_{n}-\Delta t L_{H} B, 2 L_{n}+\right.$ $\left.\Delta t L_{H} B\right]$, respectively, and $L_{n}$ is the Lipschitz constant of $u^{n}$. Notice that if the function $H$ is globally Lipschitz continuous we have the same estimate with $L_{H 2}=L_{H}$, where now $L_{H}$ is the global Lipschitz constant of $H$. Notice also that we have used the fact that, by definition,

$$
\frac{\sqrt{\beta_{0}^{+}\left[u_{j}^{n}\right]}}{\Delta x}=\frac{D^{+} u_{j}^{n}-D^{-} u_{j}^{n}}{\Delta x}=D^{2} u_{j}^{n},
$$

and that $x_{j} \in \mathcal{R}^{n} \Rightarrow D^{2} u_{j}^{n}<B$, for some constant $B>0$ independent on $n$, by Lemma 3.1 or by the definition of $\mathcal{R}^{n}(47)$. Then, the last statement would follow with $C=K B\left(\lambda L_{H}^{2}+2 \overline{L_{h^{M}}}\right)$.

Let us now prove the main statement proceeding, as usual, by induction on $n \geq 0$ and noticing that it is sufficient to prove (4.1) for $i$ and $j$ such that $i=j \pm 1$.

For $n=0$, as we take $u_{j}^{\sigma}=v_{0}\left(x_{j}\right)$ for $j \in \mathbb{Z}$, we have that 4.1 is satisfied by the Lipschitz continuity assumption on $v_{0}$, with constant $L_{0}$.

Now, assuming that 4.1 is satisfied for $n-1>0$ so that $\varepsilon^{k}$ for $k=0, \ldots, n-1$ are bounded
by (48), we can compute

$$
\begin{aligned}
\frac{\left|u_{i}^{n}-u_{j}^{n}\right|}{\Delta x} & =\frac{1}{\Delta x}\left|S^{M}\left(u^{n-1}\right)_{i}+\phi_{i} \varepsilon^{n-1} \Delta t F(\cdot)_{i}-S^{M}\left(u^{n-1}\right)_{j}-\phi_{j} \varepsilon^{n-1} \Delta t F(\cdot)_{j}\right| \\
& \leq \frac{1}{\Delta x}\left(\left|S^{M}\left(u^{n-1}\right)_{i}-S^{M}\left(u^{n-1}\right)_{j}\right|+\varepsilon^{n-1} \Delta t\left|\phi_{i} F(\cdot)_{i}-\phi_{j} F(\cdot)_{j}\right|\right) \\
& \leq \frac{\left|u_{i}^{n-1}-u_{j}^{n-1}\right|}{\Delta x}+\frac{2 \Delta t}{\Delta x} \varepsilon^{n-1}
\end{aligned}
$$

then, iterating back and using the same arguments,

$$
\begin{aligned}
\frac{\left|u_{i}^{n}-u_{j}^{n}\right|}{\Delta x} & \leq \frac{\left|u_{i}^{n-1}-u_{j}^{n-1}\right|}{\Delta x}+2 \Delta t C \leq \ldots \\
& \leq \frac{\left|u_{i}^{1}-u_{j}^{1}\right|}{\Delta x}+2(n-1) \Delta t C \\
& \leq \frac{\left|u_{i}^{0}-u_{j}^{0}\right|}{\Delta x}+2 n \Delta t C \\
& \leq L_{0}+2 \frac{T}{\Delta t} \Delta t C=L
\end{aligned}
$$

where $C$ is well defined by 48 . Notice that we have used the monotonicity of $S^{M}$ and the fact that $|F| \leq 1,|\phi| \leq 1$.

Therefore, it is clear that by construction our scheme is $\varepsilon$-monotone, in the sense of the following Definition 4.2 ( $\varepsilon$-monotonicity). A numerical scheme $S$ is $\varepsilon$-monotone if for any functions $u, v$,

$$
u \leq v \Rightarrow S(u) \leq S(v)+C \varepsilon \Delta t
$$

where $C$ is constant and $\varepsilon \rightarrow 0$ as $\Delta=(\Delta t, \Delta x) \rightarrow 0$.
Finally, we conclude this section giving our convergence result for the Adaptive Filtered Schemes.
Theorem 4.3. Let the assumptions on $S^{M}$ and $S^{A}$ be satisfied. Assume that $v_{0}$ and $H$ are Lipschitz continuous functions, $u_{j}^{n+1}$ is computed by (3)-(31), with $K>1 / 2$ and $\lambda=\frac{\Delta t}{\Delta x}$, a constant such that (6) is satisfied. Let us denote by $v_{j}^{n}:=v\left(t^{n}, x_{j}\right)$ the values of the viscosity solution on the nodes of the grid. Then,
i) the AF scheme (3) satisfies Crandall-Lions estimate [7]

$$
\left\|u^{n}-v^{n}\right\|_{\infty} \leq C_{1} \sqrt{\Delta x}, \quad \forall n=0, \ldots, N
$$

for some constant $C_{1}>0$ independent of $\Delta x$.
ii) (First order convergence for regular solutions) Moreover, if $v \in C^{2}([0, T] \times \mathbb{R})$, then

$$
\left\|u^{n}-v^{n}\right\|_{\infty} \leq C_{2} \Delta x, \quad \forall n=0, \ldots, N
$$

for some constant $C_{2}>0$ independent of $\Delta x$.
iii) (High-order local consistency) Let $k \geq 2$ be the order of the scheme $S^{A}$. If $v \in C^{l+1}$ in some neighborhood of a point $(t, x) \in[0, T] \times \mathbb{R}$, then for $1 \leq l \leq k$,

$$
\mathcal{E}_{A F}\left(v^{n}\right)_{j}=\mathcal{E}_{A}\left(v^{n}\right)_{j}=O\left(\Delta x^{l}\right)+O\left(\Delta t^{l}\right)
$$

for $t^{n}-t, x_{j}-x, \Delta t, \Delta x$ sufficiently small.
Proof. i) Let us proceed as has been done in [5] defining $w_{j}^{n+1}=S^{M}\left(w^{n}\right)_{j}$, the solution computed with the monotone scheme alone with $w_{j}^{0}=v_{0}\left(x_{j}\right)$. Then by definition,

$$
\begin{equation*}
u_{j}^{n+1}-w_{j}^{n+1}=S^{M}\left(u^{n}\right)_{j}-S^{M}\left(w^{n}\right)_{j}+\phi_{j} \varepsilon^{n} \Delta t F\left(\frac{S^{A}\left(u^{n}\right)_{j}-S^{M}\left(u^{n}\right)_{j}}{\varepsilon^{n} \Delta t}\right), \tag{49}
\end{equation*}
$$

whence, exploiting the nonexpansivity in $L^{\infty}$ of $S^{M}$, the definition of $\varepsilon^{n}$ and that $|F| \leq 1$,

$$
\begin{equation*}
\max _{j}\left|u_{j}^{n+1}-w_{j}^{n+1}\right| \leq \max _{j}\left|u_{j}^{n}-w_{j}^{n}\right|+\varepsilon^{n} \Delta t . \tag{50}
\end{equation*}
$$

Then, proceeding recursively on $n \leq N$ and recalling that by Proposition 4.1 there exists a constant $C>0$ such that $\varepsilon^{n} \leq C \Delta x:=\varepsilon$ for each $n$,

$$
\begin{equation*}
\max _{j}\left|u_{j}^{n}-w_{j}^{n}\right| \leq \sum_{k=0}^{n-1} \varepsilon^{k} \Delta t \leq n \varepsilon \Delta t \leq T \varepsilon . \tag{51}
\end{equation*}
$$

At this point, by the triangular inequality

$$
\begin{equation*}
\max _{j}\left|u_{j}^{n+1}-v_{j}^{n+1}\right| \leq \max _{j}\left|u_{j}^{n+1}-w_{j}^{n+1}\right|+\max _{j}\left|w_{j}^{n+1}-v_{j}^{n+1}\right|, \tag{52}
\end{equation*}
$$

whence we have that

$$
\begin{equation*}
\max _{j}\left|u_{j}^{n+1}-v_{j}^{n+1}\right| \leq \max _{j}\left|w_{j}^{n}-v_{j}^{n}\right|+\varepsilon T \leq\left(C_{C L}+C T\right) \sqrt{\Delta x}, \tag{53}
\end{equation*}
$$

with $C_{C L}>0$ given by the Crandall-Lions estimate for $S^{M}$.
ii) Let us recall that by (4), in the case of $v \in C^{2}$ the consistency error for the monotone scheme is such that $\mathcal{E}_{M}\left(v^{n}\right)_{j} \leq C_{M}(\Delta t+\Delta x)$. Then we can compute

$$
\begin{aligned}
\left|u_{j}^{n+1}-v_{j}^{n+1}\right| & =\left|S^{M}\left(u^{n}\right)_{j}+\phi_{j} \varepsilon^{n} \Delta t F(\cdot)-v_{j}^{n+1}\right| \\
& \leq\left|S^{M}\left(u^{n}\right)_{j}-S^{M}\left(v^{n}\right)_{j}\right|+\left|S^{M}\left(v^{n}\right)_{j}-v_{j}^{n+1}\right|+\varepsilon^{n} \Delta t \\
& \leq\left\|u^{n}-v^{n}\right\|_{\infty}+\Delta t\left(\mathcal{E}_{M}\left(v^{n}\right)+\varepsilon^{n}\right),
\end{aligned}
$$

whence, by recursion on $n \leq N$ and recalling what we have done in the previous point,

$$
\begin{equation*}
\left\|u^{n}-v^{n}\right\|_{\infty} \leq\left\|u^{0}-v^{0}\right\|_{\infty}+T\left(\max _{k=0, \ldots, n-1}\left\|\mathcal{E}_{M}\left(v^{k}\right)\right\|_{\infty}+\varepsilon\right) . \tag{54}
\end{equation*}
$$

To finish this proof what is left is to use the estimate on $\mathcal{E}_{M}$ and Proposition 4.1.
iii) In order to show that $S^{A F}\left(v^{n}\right)_{j}=S^{A}\left(v^{n}\right)_{j}$ for $\Delta t$ e $\Delta x$ small enough it is sufficient to prove that

$$
\begin{equation*}
\frac{\left|S^{A}\left(v^{n}\right)_{j}-S^{M}\left(v^{n}\right)_{j}\right|}{\varepsilon^{n} \Delta t} \leq 1, \quad \text { for }(\Delta t, \Delta x) \rightarrow 0 \tag{55}
\end{equation*}
$$

which follows directly from the computation we have done in section 2.3 for the tuning of the parameter $\varepsilon^{n}$. In fact, if we plug (31) inside the previous inequality, we can deduce that

$$
\frac{\left|S^{A}\left(v^{n}\right)_{j}-S^{M}\left(v^{n}\right)_{j}\right|}{\varepsilon^{n} \Delta t} \leq \frac{1}{2 K}+O(\Delta x)+O(\Delta t)
$$

which, using that $K>1 / 2$ by assumption, leads to the thesis as $(\Delta t, \Delta x) \rightarrow 0$. Notice that we have used the property $\varepsilon^{n}=O(\Delta x)$ and exploited the CFL condition.

Remark 4.4. Notice that the assumption $M(\Delta x)=\frac{1}{2}-C \Delta x$, for some constant $C>0$ such that $M(\Delta x)>0$, needed to apply Lemma 3.1, may give some problems in the proof of third assertion of the previous theorem. In fact, applying the standard definition (41) to the viscosity solution $v$ at a point $x_{j}$ and recalling the computations that led to (39), we get that

$$
\omega_{j}^{ \pm}=\frac{1}{2} \mp \Delta x \frac{4 v_{j}^{\prime \prime} v_{j}^{\prime \prime \prime}}{\left(v_{j}^{\prime \prime}\right)^{2}+\sigma}+O\left(\Delta x^{2}\right) .
$$

Consequently, in order to be sure that if $v \in C^{3}$, then $j \in \mathcal{R}$, we have to choose the constant $C$ such that

$$
C \geq\left|\frac{4 v_{j}^{\prime \prime} v_{j}^{\prime \prime \prime}}{\left(v_{j}^{\prime \prime}\right)^{2}+\sigma}\right|,
$$

or require additional smoothness assumptions on $v$, for example $v_{j}^{\prime \prime \prime} \ll v_{j}^{\prime \prime}$. This in fact poses a strong limitation on the applicability of Lemma 3.1, at least in the present formulation.

## 5 Numerical Tests

In this section we present some one-dimensional examples designed to show the properties of our scheme, stated by Theorem 4.3. Our goal is also to compare the performances of our Adaptive Filtered Scheme $S^{A F}$ with those of the Basic Filtered Scheme $S^{F}$ of 5 and with the WENO scheme of second/third order of [15]. For each test we specify the monotone and high-order schemes composing the filtered scheme and compute the errors and orders in $L^{\infty}$ and $L^{1}$ norm. At the end of the section, we will also show briefly how to use these schemes in order to approximate simple two dimensional problems. To be precise, in the following examples we will refer to the standard CFL condition

$$
\begin{equation*}
\lambda \max \left|H_{p}(p)\right| \leq 1, \tag{56}
\end{equation*}
$$

to define $\lambda$, which is in fact equivalent to (6) and more easily computed.
Example 1: Transport equation. In order to test the capability of our scheme to handle both regular and singular regions, let us begin with a simple linear example and consider the problem

$$
\left\{\begin{array}{l}
v_{t}(t, x)+v_{x}(t, x)=0 \quad \text { in }(0, T) \times \Omega \\
v(0, x)=v_{0}(x),
\end{array}\right.
$$

with periodic boundary conditions, in two different situations. At first, aiming to test the full accuracy of the schemes, we consider the regular initial data (Case a),

$$
\begin{equation*}
v_{0}(x)=\sin (\pi x), \quad x \in \Omega \tag{57}
\end{equation*}
$$

with $\Omega=[-2,2]$ ans $T=0.9$. Then, as a second test, we take the mixed initial data (Case b),

$$
v_{0}(x)= \begin{cases}\min \left\{(1-x)^{2},(1+x)^{2}\right\} & \text { if }-1 \leq x \leq 1  \tag{58}\\ \sin ^{2}(\pi(x-2)) & \text { if } 2 \leq x \leq 3 \\ 0 & \text { otherwise }\end{cases}
$$

with $\Omega=[-1.5,3.5]$ and $T=2$. The latter problem models the transport of a function composed by two peaks, the first with one point of singularity while the second is in $C^{2}$. For these tests we use the Central Upwind scheme (8) as monotone scheme and the simple Heun-Centered (HC) scheme (15)-(16) as high-order scheme, with $\lambda=0.9$ for Case $a$ and $\lambda=0.4$ for Case $b$. We also compare the results obtained using $S^{A F}$ with the 4th order Lax-Wendroff scheme 20) as high-order scheme. We recall that the latter high-order scheme has a very compact 5 -points stencil, while the WENO scheme of second/third order (coupled with the third order Runge Kutta scheme) has a stencil of nine points.


Figure 2: (Example 1a.) Plots at time $T=0.9$ with the AF-HC scheme on the left and WENO on the right for $\Delta x=0.05$.

Table 1: (Example 1a.) Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | F-HC $(5 \Delta x)$ |  | AF-HC |  | AF-LW4ord |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 40 | 10 | $1.36 \mathrm{e}-02$ |  | $1.36 \mathrm{e}-02$ |  | $1.37 \mathrm{e}-04$ |  | $8.02 \mathrm{e}-02$ |  |
| 80 | 20 | $2.56 \mathrm{e}-03$ | 2.41 | $2.56 \mathrm{e}-03$ | 2.41 | $8.66 \mathrm{e}-06$ | 3.98 | $2.62 \mathrm{e}-02$ | 1.62 |
| 160 | 40 | $5.76 \mathrm{e}-04$ | 2.15 | $5.76 \mathrm{e}-04$ | 2.15 | $5.43 \mathrm{e}-07$ | 4.00 | $4.50 \mathrm{e}-03$ | 2.54 |
| 320 | 80 | $1.40 \mathrm{e}-04$ | 2.04 | $1.40 \mathrm{e}-04$ | 2.04 | $3.40 \mathrm{e}-08$ | 4.00 | $1.95 \mathrm{e}-04$ | 4.52 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 40 | 10 | $3.58 \mathrm{e}-02$ |  | $3.58 \mathrm{e}-02$ |  | $3.62 \mathrm{e}-04$ |  | $2.07 \mathrm{e}-01$ |  |
| 80 | 20 | $6.66 \mathrm{e}-03$ | 2.43 | $6.66 \mathrm{e}-03$ | 2.43 | $2.25 \mathrm{e}-05$ | 4.01 | $4.14 \mathrm{e}-02$ | 2.32 |
| 160 | 40 | $1.48 \mathrm{e}-03$ | 2.17 | $1.48 \mathrm{e}-03$ | 2.17 | $1.40 \mathrm{e}-06$ | 4.01 | $5.09 \mathrm{e}-03$ | 3.02 |
| 320 | 80 | $3.57 \mathrm{e}-04$ | 2.05 | $3.57 \mathrm{e}-04$ | 2.05 | $8.69 \mathrm{e}-08$ | 4.01 | $3.08 \mathrm{e}-04$ | 4.05 |

In the first test all the schemes are very accurate and achieve optimal order in both norms, as shown in Figure 2 and Table 1. In this case both filtered schemes have the same numerical results
and coincide with the HC high-order scheme (we avoided to report the same results), as wanted. Moreover, we can see that our fourth order scheme is much more accurate even than the WENO scheme, despite the smaller stencil required.


Figure 3: (Example 1b.) Plots of the solution at time $T=2$ with $\Delta x=0.025$. Top: basic filtered scheme with HC on the left, adaptive on the right. Bottom: fourth order scheme AF scheme on the left and WENO on the right.

For the second case, looking at Figure 3 we can observe that the adaptive tuning of $\varepsilon^{n}$ is able to contain the oscillations behind the peaks produced by the unstable HC scheme, which are clearly visible instead in the case of $S^{F}$ with $\varepsilon=5 \Delta x$. We can also see that our scheme coupled with the fourth order scheme produces again the best results in terms of errors and orders in both norms (see Table 2) and gives the best resolution of the peaks, preserving better the kink of the singularity and the feet of the regular part, without introducing any oscillation.

Example 2: Eikonal equation. As a first nonlinear problem let us consider the eikonal equation

$$
\left\{\begin{array}{l}
v_{t}(t, x)+\left|v_{x}(t, x)\right|=0  \tag{59}\\
v_{0}(x)=\max \left\{1-x^{2}, 0\right\}^{4}
\end{array} \text { in }(0,0.3) \times(-2,2),\right.
$$

where $v_{0}$ is a Lipschitz continuous initial data with high regularity. Then, we repeat the simulation

Table 2: (Example 1b.) Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | F-HC $(10 \Delta x)$ |  | AF-HC |  | AF-LW4ord |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 50 | 50 | $3.46 \mathrm{e}-01$ |  | $2.98 \mathrm{e}-01$ |  | $2.65 \mathrm{e}-01$ |  | $3.47 \mathrm{e}-01$ |  |
| 100 | 100 | $1.41 \mathrm{e}-01$ | 1.29 | $1.78 \mathrm{e}-01$ | 0.75 | $1.56 \mathrm{e}-01$ | 0.77 | $2.07 \mathrm{e}-01$ | 0.75 |
| 200 | 200 | $9.69 \mathrm{e}-02$ | 0.54 | $1.12 \mathrm{e}-01$ | 0.66 | $9.08 \mathrm{e}-02$ | 0.78 | $1.28 \mathrm{e}-01$ | 0.70 |
| 400 | 400 | $7.29 \mathrm{e}-02$ | 0.41 | $7.05 \mathrm{e}-02$ | 0.67 | $5.06 \mathrm{e}-02$ | 0.84 | $7.66 \mathrm{e}-02$ | 0.74 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 50 | 50 | $4.34 \mathrm{e}-01$ |  | $2.94 \mathrm{e}-01$ |  | $2.21 \mathrm{e}-01$ |  | $3.62 \mathrm{e}-01$ |  |
| 100 | 100 | $1.41 \mathrm{e}-01$ | 1.63 | $9.77 \mathrm{e}-02$ | 1.59 | $4.26 \mathrm{e}-02$ | 2.38 | $1.39 \mathrm{e}-01$ | 1.39 |
| 200 | 200 | $4.24 \mathrm{e}-02$ | 1.73 | $3.06 \mathrm{e}-02$ | 1.67 | $9.22 \mathrm{e}-03$ | 2.21 | $3.83 \mathrm{e}-02$ | 1.86 |
| 400 | 400 | $1.38 \mathrm{e}-02$ | 1.62 | $1.01 \mathrm{e}-02$ | 1.60 | $2.61 \mathrm{e}-03$ | 1.82 | $8.39 \mathrm{e}-03$ | 2.19 |

with the "reversed" initial data

$$
\begin{equation*}
v_{0}(x)=-\max \left\{1-x^{2}, 0\right\}^{4}, \tag{60}
\end{equation*}
$$

which presents also a major problem in the origin because of the saddle point in the hamiltonian, where two directions of propagation occur. Here the aim is mainly to compare the results obtained by the unfiltered high-order schemes with their filtered versions, in order to show the stabilization property of the filtering process. For the monotone scheme we use the numerical hamiltonian (7), while to achieve high-order we use the Lax-Wendroff-Richtmyer (LWR) scheme 19). Moreover, as in the previous example, we present also the results obtained with the AF scheme coupled with the fourth order LW scheme. The CFL number is set to 0.375 for both simulations.


Figure 4: (Example 2a.) Initial data (left) and plots of the solution at time $T=0.3$ with the AF scheme (center) and the LWR scheme (right) for $\Delta x=0.025$.

Let us first point out that, as Figures 45 clearly show, the LWR scheme is unstable in the origin in both situations, while the AF scheme (and the basic filtered scheme) is not. Then, for the first case, looking at Table 3 we can see that the filtered-LWR schemes give almost the same results, are of high-order in both norms and have better errors even than the WENO scheme in almost all simulations. Moreover, we can recognize the typical, as will be shown also by the following examples, improvements and drawbacks of the fourth order LW scheme, which has a slightly wider stencil. In fact, the scheme has bigger errors in the $L^{\infty}$ norm in the first three refinements of the grid, while has way better errors and orders in the $L^{1}$ norm, achieving almost optimal order, which

Table 3: (Example 2a.) Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | F-LWR $(5 \Delta x)$ |  | AF-LWR |  | AF-LW4ord |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 40 | 8 | $1.96 \mathrm{e}-02$ |  | $1.64 \mathrm{e}-02$ |  | $2.18 \mathrm{e}-02$ |  | $6.81 \mathrm{e}-02$ |  |
| 80 | 16 | $4.48 \mathrm{e}-03$ | 2.13 | $4.00 \mathrm{e}-03$ | 2.04 | $9.98 \mathrm{e}-03$ | 1.13 | $3.42 \mathrm{e}-02$ | 1.00 |
| 160 | 32 | $1.06 \mathrm{e}-03$ | 2.08 | $1.11 \mathrm{e}-03$ | 1.85 | $1.35 \mathrm{e}-03$ | 2.89 | $1.62 \mathrm{e}-02$ | 1.08 |
| 320 | 64 | $2.56 \mathrm{e}-04$ | 2.05 | $2.56 \mathrm{e}-04$ | 2.12 | $2.31 \mathrm{e}-04$ | 2.55 | $7.52 \mathrm{e}-03$ | 1.11 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 40 | 8 | $1.52 \mathrm{e}-02$ |  | $1.16 \mathrm{e}-02$ |  | $1.11 \mathrm{e}-02$ |  | $2.05 \mathrm{e}-02$ |  |
| 80 | 16 | $3.78 \mathrm{e}-03$ | 2.01 | $3.71 \mathrm{e}-03$ | 1.65 | $1.05 \mathrm{e}-03$ | 3.40 | $4.68 \mathrm{e}-03$ | 2.13 |
| 160 | 32 | $8.94 \mathrm{e}-04$ | 2.08 | $8.96 \mathrm{e}-04$ | 2.05 | $7.28 \mathrm{e}-05$ | 3.85 | $9.55 \mathrm{e}-04$ | 2.29 |
| 320 | 64 | $2.09 \mathrm{e}-04$ | 2.09 | $2.09 \mathrm{e}-04$ | 2.10 | $7.14 \mathrm{e}-06$ | 3.35 | $1.40 \mathrm{e}-04$ | 2.78 |

testifies the overall improvement.


Figure 5: (Example 2b.) Plots at time $T=0.3$ with the AF and WENO schemes for $\Delta x=0.05$ (left) and LWR scheme for $\Delta x=0.0125$.

For Case b, as testified by Table 4, we can repeat almost the same considerations, but this time the improvements given by the adaptive filtering are evident. The AF-LWR scheme is again of high-order especially in the $L^{1}$ norm, without the need to introduce any limiter as has been done in [5], and the numerical results are always comparable to those obtained by the WENO scheme of second/third order, while the AF-LW4ord scheme has again worse $L^{\infty}$ errors for the first discretizations and better errors and orders in the $L^{1}$ norm.

Example 3: Burgers' equation. Let us consider now the Burgers' equation for HJ with a regular initial data

$$
\left\{\begin{array}{l}
v_{t}(t, x)+\frac{1}{2}\left(v_{x}(t, x)+1\right)^{2}=0 \quad \text { in }(0, T) \times(0,2),  \tag{61}\\
v_{0}(x)=-\cos (\pi x),
\end{array}\right.
$$

which is a test case widely used in literature. In order to test the full accuracy of the schemes even in the nonlinear case, we first run the simulation for $T=\frac{4}{5 \pi^{2}}$, when the solution is still regular, with $\lambda=\frac{2}{\pi^{2}} \approx 0.2<\max \left|H_{p}\right|^{-1}=0.5$. Then, we consider the final time $T=\frac{3}{2 \pi^{2}}$ when a moving

Table 4: (Example 2b.) Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | F-LWR $(5 \Delta x)$ |  | AF-LWR |  | AF-LW4ord |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 40 | 8 | $1.91 \mathrm{e}-02$ |  | $1.40 \mathrm{e}-02$ |  | $1.63 \mathrm{e}-02$ |  | $2.33 \mathrm{e}-02$ |  |
| 80 | 16 | $9.24 \mathrm{e}-03$ | 1.04 | $3.37 \mathrm{e}-03$ | 2.06 | $7.51 \mathrm{e}-03$ | 1.11 | $1.02 \mathrm{e}-02$ | 1.19 |
| 160 | 32 | $5.77 \mathrm{e}-03$ | 0.68 | $1.58 \mathrm{e}-03$ | 1.09 | $2.14 \mathrm{e}-03$ | 1.81 | $4.10 \mathrm{e}-03$ | 1.32 |
| 320 | 64 | $3.46 \mathrm{e}-03$ | 0.74 | $7.09 \mathrm{e}-04$ | 1.16 | $6.92 \mathrm{e}-04$ | 1.63 | $1.22 \mathrm{e}-03$ | 1.75 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 40 | 8 | $2.38 \mathrm{e}-02$ |  | $2.01 \mathrm{e}-02$ |  | $1.29 \mathrm{e}-02$ |  | $2.96 \mathrm{e}-02$ |  |
| 80 | 16 | $8.48 \mathrm{e}-03$ | 1.49 | $5.70 \mathrm{e}-03$ | 1.82 | $2.05 \mathrm{e}-03$ | 2.65 | $7.04 \mathrm{e}-03$ | 2.07 |
| 160 | 32 | $3.41 \mathrm{e}-03$ | 1.32 | $1.82 \mathrm{e}-03$ | 1.65 | $3.20 \mathrm{e}-04$ | 2.68 | $1.43 \mathrm{e}-03$ | 2.30 |
| 320 | 64 | $1.52 \mathrm{e}-03$ | 1.17 | $5.84 \mathrm{e}-04$ | 1.64 | $6.38 \mathrm{e}-05$ | 2.33 | $2.82 \mathrm{e}-04$ | 2.34 |

(to the right) singularity appears, taking $\lambda=\frac{15}{8 \pi^{2}} \approx 0.19$. For both simulations we use the the Central Upwind monotone scheme and the LWR scheme for both the filtered schemes and compare the results as before with the WENO scheme and the fourth order AF scheme.


Figure 6: (Example 3.) From left to right: initial data of problem 61) and plots of the solution at time $T=4 /\left(5 \pi^{2}\right)$ and $T=3 /\left(2 \pi^{2}\right)$ for $\Delta x=0.025$.

This example summarizes all the behaviors already seen in the previous cases. In fact, as displayed by Tables 556, if the solution is still regular the fourth order AF scheme gives the best results and achieves the optimal order in both norms, while when the singularity appears has the usual problems in the $L^{\infty}$ norm and better orders (than the second order scheme) in the $L^{1}$ norm. Here we have to notice that the WENO scheme has better errors and orders in the second simulation with respect to all the filtered schemes. Moreover, we can clearly see that the basic filtered scheme depends heavily on the choice of $\varepsilon$, in fact after extensive computations we noticed that choosing for example $\varepsilon=5 \Delta x$ we get worse results in both cases, while if we increase the constant we get better results in the regular case and worse in the latter. In the tables we presented the results for the choice that gives the best results in the singular case, while it has clearly problems in the first situation. This is the main advantage of the adaptive $\varepsilon^{n}$ which is able to tune itself in the right way depending on the local (in time) regularity of the solution.

Example 4: Evolution in 2D by dimensional splitting. In the following we will show a convenient procedure to solve simple two-dimensional problems by making use of the one-dimensional schemes defined in the previous sections. Let us consider a classical problem similar to the Burgers'

Table 5: (Example 3.) $T=4 /\left(5 \pi^{2}\right)$. Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | F-LWR $(10 \Delta x)$ |  | AF-LWR |  | AF-LW4ord |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 40 | 8 | $1.30 \mathrm{e}-02$ |  | $9.61 \mathrm{e}-03$ |  | $1.89 \mathrm{e}-03$ |  | $1.04 \mathrm{e}-02$ |  |
| 80 | 16 | $8.67 \mathrm{e}-03$ | 0.59 | $2.77 \mathrm{e}-03$ | 1.79 | $2.84 \mathrm{e}-04$ | 2.73 | $2.12 \mathrm{e}-03$ | 2.30 |
| 160 | 32 | $5.07 \mathrm{e}-03$ | 0.77 | $7.24 \mathrm{e}-04$ | 1.94 | $2.68 \mathrm{e}-05$ | 3.41 | $1.82 \mathrm{e}-04$ | 3.54 |
| 320 | 64 | $2.66 \mathrm{e}-03$ | 0.93 | $1.83 \mathrm{e}-04$ | 1.99 | $1.89 \mathrm{e}-06$ | 3.83 | $2.05 \mathrm{e}-05$ | 3.15 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 40 | 8 | $3.76 \mathrm{e}-03$ |  | $3.30 \mathrm{e}-03$ |  | $2.76 \mathrm{e}-04$ |  | $3.67 \mathrm{e}-03$ |  |
| 80 | 16 | $1.29 \mathrm{e}-03$ | 1.54 | $8.20 \mathrm{e}-04$ | 2.01 | $1.97 \mathrm{e}-05$ | 3.81 | $6.57 \mathrm{e}-04$ | 2.48 |
| 160 | 32 | $4.49 \mathrm{e}-04$ | 1.52 | $2.04 \mathrm{e}-04$ | 2.01 | $1.50 \mathrm{e}-06$ | 3.71 | $5.43 \mathrm{e}-05$ | 3.60 |
| 320 | 64 | $1.82 \mathrm{e}-04$ | 1.30 | $5.09 \mathrm{e}-05$ | 2.00 | $1.04 \mathrm{e}-07$ | 3.86 | $2.98 \mathrm{e}-06$ | 4.19 |

Table 6: (Example 3.) $T=3 /\left(2 \pi^{2}\right)$. Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | F-LWR $(10 \Delta x)$ |  | AF-LWR |  | AF-LW4ord |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 40 | 16 | $4.88 \mathrm{e}-02$ |  | $5.30 \mathrm{e}-02$ |  | $6.31 \mathrm{e}-02$ |  | $3.89 \mathrm{e}-02$ |  |
| 80 | 32 | $2.47 \mathrm{e}-02$ | 0.98 | $2.47 \mathrm{e}-02$ | 1.10 | $2.87 \mathrm{e}-02$ | 1.13 | $1.61 \mathrm{e}-02$ | 1.27 |
| 160 | 64 | $9.81 \mathrm{e}-03$ | 1.33 | $9.95 \mathrm{e}-03$ | 1.31 | $1.03 \mathrm{e}-02$ | 1.48 | $5.12 \mathrm{e}-03$ | 1.65 |
| 320 | 128 | $2.57 \mathrm{e}-03$ | 1.93 | $2.59 \mathrm{e}-03$ | 1.94 | $2.69 \mathrm{e}-03$ | 1.94 | $8.40 \mathrm{e}-04$ | 2.61 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 40 | 16 | $5.17 \mathrm{e}-03$ |  | $5.28 \mathrm{e}-03$ |  | $3.83 \mathrm{e}-03$ |  | $3.69 \mathrm{e}-03$ |  |
| 80 | 32 | $1.26 \mathrm{e}-03$ | 2.03 | $1.27 \mathrm{e}-03$ | 2.06 | $8.89 \mathrm{e}-04$ | 2.11 | $6.94 \mathrm{e}-04$ | 2.41 |
| 160 | 64 | $2.86 \mathrm{e}-04$ | 2.14 | $2.87 \mathrm{e}-04$ | 2.14 | $1.43 \mathrm{e}-04$ | 2.64 | $8.67 \mathrm{e}-05$ | 3.00 |
| 320 | 128 | $5.68 \mathrm{e}-05$ | 2.33 | $5.68 \mathrm{e}-05$ | 2.34 | $1.82 \mathrm{e}-05$ | 2.97 | $6.40 \mathrm{e}-06$ | 3.76 |

equation, already studied for example in [10, that is

$$
\left\{\begin{array}{l}
v_{t}+\frac{1}{2}\left(v_{x}^{2}+v_{y}^{2}\right)=0  \tag{62}\\
v(0, x, y)=\max \left\{0,1-\left(x^{2}+y^{2}\right)\right\} .
\end{array} \quad \text { in }(0,0.5] \times[-2,2]^{2}\right.
$$

The exact solution has discontinuous second derivatives along the circle of radius 1 centered in the origin, where for $t \geq 1 / 2$ develops also a singularity in the gradient. In particular, for $t \geq 1 / 2$, it can be computed as

$$
v(t, x, y)= \begin{cases}\frac{\left(\sqrt{x^{2}+y^{2}}-1\right)^{2}}{2 t} & \text { if } x^{2}+y^{2} \leq 1 \\ 0 & \text { otherwise }\end{cases}
$$

In this situation, since the hamiltonian can be expressed as a sum of one-dimensional hamiltonian, depending on the evolution along the $x$ and $y$ direction, respectively, we can use a dimensional splitting to solve the problem. More precisely, if we write $H\left(v_{x}, v_{y}\right)=H_{1}\left(v_{x}\right)+H_{2}\left(v_{y}\right)$, we can
approximate the solution by solving sequentially the problems in one space dimension

$$
v_{t}+H_{1}\left(v_{x}\right)=0 \quad \text { and } \quad v_{t}+H_{2}\left(v_{y}\right)=0,
$$

keeping each time the other space variable constant. Since the hamiltonians trivially commute, we can use the simple Lie-Trotter splitting

$$
\begin{equation*}
u^{n+1}=S_{y}^{\Delta t}\left(S_{x}^{\Delta t}\left(u^{n}\right)\right), \tag{63}
\end{equation*}
$$

where $S_{x}^{\Delta t}$ and $S_{y}^{\Delta t}$ are numerical schemes of time step $\Delta t$ for the problems in the $x$ and $y$ direction, respectively, without introducing errors in the time evolution. For more details about dimensional splitting techniques we refer the reader to [19] and the references therein.

For this problem the low regularity of the solution plays a major role. In fact, as shown in Table 7. all the tested schemes reach at most a first order convergence rate, in both norms. In particular, we can acknowledge the better results given by the second order filtered scheme, with the F scheme performing slightly better in most situations, and the fact that the higher order schemes do not produce any relevant improvement, not even in the $L^{1}$ norm. Moreover, from Figure 7 we can see that our scheme gives nicely sharp results even with a coarse grid, at least comparably to those of the WENO scheme.


Figure 7: (Example 4.) Top: Initial datum (left) and exact solution at $T=0.5$ (right). Bottom: solution at $T=0.5$ computed by the AF-HC scheme (left) and the WENO scheme (right) with $\Delta x=0.05$.

Table 7: (Example 4.). Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | F-HC $(10 \Delta x)$ |  | AF-HC |  | AF-LW4ord |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 40 | 20 | $7.50 \mathrm{e}-02$ |  | $7.64 \mathrm{e}-02$ |  | $1.35 \mathrm{e}-01$ |  | $9.97 \mathrm{e}-02$ |  |
| 80 | 40 | $3.92 \mathrm{e}-02$ | 0.94 | $4.57 \mathrm{e}-02$ | 0.74 | $6.79 \mathrm{e}-02$ | 0.99 | $5.42 \mathrm{e}-02$ | 0.88 |
| 160 | 80 | $2.86 \mathrm{e}-02$ | 0.45 | $2.54 \mathrm{e}-02$ | 0.85 | $3.39 \mathrm{e}-02$ | 1.00 | $2.76 \mathrm{e}-02$ | 0.98 |
| 320 | 160 | $1.92 \mathrm{e}-02$ | 0.57 | $1.33 \mathrm{e}-02$ | 0.93 | $1.71 \mathrm{e}-02$ | 0.99 | $1.35 \mathrm{e}-02$ | 1.03 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 40 | 20 | $7.01 \mathrm{e}-02$ |  | $8.69 \mathrm{e}-02$ |  | $6.75 \mathrm{e}-02$ |  | $7.76 \mathrm{e}-02$ |  |
| 80 | 40 | $3.10 \mathrm{e}-02$ | 1.18 | $4.26 \mathrm{e}-02$ | 1.03 | $3.79 \mathrm{e}-02$ | 0.83 | $3.88 \mathrm{e}-02$ | 1.00 |
| 160 | 80 | $1.68 \mathrm{e}-02$ | 0.88 | $2.12 \mathrm{e}-02$ | 1.01 | $1.99 \mathrm{e}-02$ | 0.93 | $1.90 \mathrm{e}-02$ | 1.03 |
| 320 | 160 | $1.08 \mathrm{e}-02$ | 0.65 | $1.03 \mathrm{e}-02$ | 1.05 | $9.70 \mathrm{e}-03$ | 1.05 | $9.23 \mathrm{e}-03$ | 1.04 |

Example 5: Generating singularities in two dimensions. In this example we consider a problem similar to the Burgers' equation in two dimensions, which is strictly connected to (61),

$$
\left\{\begin{array}{l}
v_{t}+\left(v_{x}+1\right)^{2}+\left(v_{y}+1\right)^{2}=0 \quad \text { in }(0, T) \times \Omega,  \tag{64}\\
v(0, x, y)=-0.5(\cos (\pi x)+\cos (\pi y)),
\end{array}\right.
$$

with $\Omega=[0,2]^{2}$ and periodic boundary conditions. As for problem (61), we consider the final time $T=\frac{4}{5 \pi^{2}}$, when the solution is still smooth, and then $T=\frac{3}{2 \pi^{2}}$, time at which an interesting set of singularities develops. It is clear that we can use the dimensional splitting also in this situation and solve the problem using one-dimensional schemes and the Lie-Trotter splitting, since the hamiltonians $H_{1}$ and $H_{2}$ commute. We use the same schemes as in Example 3 and a slightly more restrictive CFL number with respect to problem (61) in order to use coarser grids, which is set to $\lambda=\frac{4}{5 \pi^{2}} \approx 0.08$ for the first test, and $\lambda=\frac{3}{4 \pi^{2}} \approx 0.076$ for the latter.

The exact solution is computed by the Hopf-Lax formula,

$$
v(t, x, y)=\left(\min _{a \in A} \frac{1}{2} \cos (x-a t)+\frac{1}{4} a^{2}-a+\min _{b \in A} \frac{1}{2} \cos (y-b t)+\frac{1}{4} b^{2}-b\right),
$$

with $A=[-5,5]$, where we used the fact that the evolution can be seen as the sum of separate one-dimensional evolutions.

As we could expect, in this example we have analogous result with respect to Example 3, with the AF scheme performing well in both situations and better than the F scheme in the regular case (see Tables 8-9). Here again the basic filtered scheme has slightly better results after the singularities develop, due to the action of the $\phi$ function in the regions of singularity, but the loss of accuracy is in fact minimal. Moreover, our scheme performs as good as the WENO scheme when the solution is still regular, while the latter performs better in the second case.

## 6 Conclusions

We have presented a rather simple way to construct convergent schemes coupling a monotone and a high-order scheme via a filter function. A typical feature of filtered schemes is their highorder accuracy in the regions of regularity for the solution. In fact, the filter function can stabilize


Figure 8: (Example 5.) Top: Initial datum (left) and exact solution at $T=3 /\left(2 \pi^{2}\right)$ (right). Bottom: solution at $T=4 /\left(5 \pi^{2}\right)$ (left) and $T=3 /\left(2 \pi^{2}\right)$ (right) computed by the AF scheme with $\Delta x=0.1$.
an otherwise unstable (high-order) scheme, still preserving its accuracy. The main novelty here is the adaptive and automatic choice of the parameter $\varepsilon^{n}$ which improves the scheme in [5. The computation of the switching parameter $\varepsilon^{n}$, although more expensive, is still affordable in low dimension. The adaptive scheme is able to reduce the oscillations which may appear choosing a constant $\varepsilon$ and, as shown by the numerical tests, gives always better results. Finally, we note that the accuracy of adaptive filtered schemes is close to WENO schemes of the same order but filter schemes are easier to implement, give a rather flexible way to couple different schemes and, as we proved, converge to the viscosity solution.

## Acknowledgements

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## References

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Table 8: (Example 5.) $T=4 /\left(5 \pi^{2}\right)$. Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | LWR |  | F-LWR (10 $\boldsymbol{c} x)$ |  | AF-LWR |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 20 | 10 | $7.45 \mathrm{e}-02$ |  | $7.75 \mathrm{e}-02$ |  | $8.32 \mathrm{e}-02$ |  | $8.66 \mathrm{e}-02$ |  |
| 40 | 20 | $3.38 \mathrm{e}-02$ | 1.14 | $5.12 \mathrm{e}-02$ | 0.60 | $2.77 \mathrm{e}-02$ | 1.59 | $3.59 \mathrm{e}-02$ | 1.27 |
| 80 | 40 | $1.49 \mathrm{e}-02$ | 1.18 | $3.25 \mathrm{e}-02$ | 0.66 | $6.58 \mathrm{e}-03$ | 2.08 | $1.30 \mathrm{e}-02$ | 1.47 |
| 160 | 80 | $6.42 \mathrm{e}-03$ | 1.22 | $1.94 \mathrm{e}-02$ | 0.75 | $1.78 \mathrm{e}-03$ | 1.89 | $4.87 \mathrm{e}-03$ | 1.41 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 20 | 10 | $3.67 \mathrm{e}-02$ |  | $4.42 \mathrm{e}-02$ |  | $3.72 \mathrm{e}-02$ |  | $3.71 \mathrm{e}-02$ |  |
| 40 | 20 | $9.53 \mathrm{e}-03$ | 1.94 | $1.21 \mathrm{e}-02$ | 1.87 | $8.50 \mathrm{e}-03$ | 2.13 | $1.00 \mathrm{e}-02$ | 1.89 |
| 80 | 40 | $2.28 \mathrm{e}-03$ | 2.06 | $4.29 \mathrm{e}-03$ | 1.49 | $2.05 \mathrm{e}-03$ | 2.05 | $1.95 \mathrm{e}-03$ | 2.36 |
| 160 | 80 | $6.51 \mathrm{e}-04$ | 1.81 | $1.70 \mathrm{e}-03$ | 1.33 | $5.42 \mathrm{e}-04$ | 1.92 | $4.50 \mathrm{e}-04$ | 2.11 |

Table 9: (Example 5.) $T=3 /\left(2 \pi^{2}\right)$. Errors and orders in $L^{\infty}$ and $L^{1}$ norms.

|  |  | LWR |  | F-LWR (10 $\Delta x)$ |  | AF-LWR |  | WENO 2/3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $N_{x}$ | $N_{t}$ | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord | $L^{\infty}$ Err | Ord |
| 20 | 20 | $1.69 \mathrm{e}-01$ |  | $9.49 \mathrm{e}-02$ |  | $1.04 \mathrm{e}-01$ |  | $8.68 \mathrm{e}-02$ |  |
| 40 | 40 | $6.39 \mathrm{e}-02$ | 1.40 | $3.67 \mathrm{e}-02$ | 1.37 | $3.80 \mathrm{e}-02$ | 1.45 | $2.27 \mathrm{e}-02$ | 1.93 |
| 80 | 80 | $3.23 \mathrm{e}-02$ | 0.98 | $1.41 \mathrm{e}-02$ | 1.38 | $1.47 \mathrm{e}-02$ | 1.37 | $9.08 \mathrm{e}-03$ | 1.32 |
| 160 | 160 | $2.64 \mathrm{e}-02$ | 0.29 | $3.73 \mathrm{e}-03$ | 1.91 | $3.90 \mathrm{e}-03$ | 1.92 | $2.22 \mathrm{e}-03$ | 2.03 |
| $N_{x}$ | $N_{t}$ | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord | $L^{1}$ Err | Ord |
| 20 | 20 | $6.18 \mathrm{e}-02$ |  | $4.62 \mathrm{e}-02$ |  | $4.62 \mathrm{e}-02$ |  | $3.60 \mathrm{e}-02$ |  |
| 40 | 40 | $1.74 \mathrm{e}-02$ | 1.83 | $8.19 \mathrm{e}-03$ | 2.50 | $8.24 \mathrm{e}-03$ | 2.49 | $4.68 \mathrm{e}-03$ | 2.94 |
| 80 | 80 | $4.54 \mathrm{e}-03$ | 1.94 | $1.88 \mathrm{e}-03$ | 2.13 | $1.80 \mathrm{e}-03$ | 2.20 | $6.92 \mathrm{e}-04$ | 2.76 |
| 160 | 160 | $1.13 \mathrm{e}-03$ | 2.00 | $3.73 \mathrm{e}-04$ | 2.33 | $3.78 \mathrm{e}-04$ | 2.25 | $7.62 \mathrm{e}-05$ | 3.18 |

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## A APPENDIX: Technical results

For completeness and reader's convenience we give the proofs of Proposition 3.1 and of the properties of the undivided differences and the binomial coefficients involved. This analysis follows the ideas in [1] where a similar analysis is developed for conservation laws.
of Proposition 3.1. Let us take $r>1$ and without loss of generality, let $x_{s}=0$ (to simplify the notation). Moreover, we will use the convention $h:=\Delta x$.

Let us start by reminding that, using the Newton form of the interpolating polynomial, for $k=0, \ldots, r-1$ and $j \in \mathbb{Z}$, we get

$$
\begin{equation*}
P_{k}(x)=f\left(x_{j-r+k}\right)+\sum_{i=1}^{r} f\left[x_{j-r+k}, \ldots, x_{j-r+k+i}\right] \omega_{i-1}(x), \tag{65}
\end{equation*}
$$

where $\omega_{i}(x)=\left(x-x_{j-r+k}\right) \cdots\left(x-x_{j-r+k+i}\right)$ and $f[\cdot]$ denotes the divided difference of $f$.
Let us proceed with the proof of $i$ ). If we define the function $f_{h}(y):=f\left(x_{j}+h y\right), y \in \mathbb{Z}$, we can write

$$
\begin{aligned}
f\left[x_{j-r+k}, \ldots, x_{j-r+k+i}\right] & =f\left[x_{j}+(k-r) h, \ldots, x_{j}+(k-r+i) h\right] \\
& =\frac{f_{h}[k-r, \ldots, k-r+i]}{i!h^{i}},
\end{aligned}
$$

where $f_{h}[\cdot]$ denotes the undivided difference of $f_{h}$. At this point it is useful to notice that, for $l=0, \ldots, i$ (see Lemma A. 1 for the proof),

$$
\begin{equation*}
f_{h}[k-r, \ldots, k-r+i]=\sum_{n=0}^{i-l}\binom{i-l}{n}(-1)^{i-l-n} f_{h}[(k-r+n), \ldots,(k-r+n+l)] . \tag{66}
\end{equation*}
$$

Moreover, if $0 \notin \stackrel{\circ}{\mathcal{S}}_{j+k}$, expanding with Taylor, it can be shown that for $n=0, \ldots, i-l$ e $h \rightarrow 0$, $f_{h}[(k-r+n), \ldots,(k-r+n+l)] \approx h^{l} f^{(l)}\left(x_{j}\right)$ (see Corollary A.3). Then, we can infer

$$
f_{h}[k-r, \ldots, k-r+i] \rightarrow \begin{cases}o\left(h^{l}\right) & \text { for } l<i \leq r  \tag{67}\\ h^{l} f^{(l)}\left(x_{j}\right) & \text { for } i=l\end{cases}
$$

having exploited the relation $\sum_{j=0}^{i}\binom{i}{j}(-1)^{i-j}=0$, for $i \geq 1$ (see Lemma A.1).
Now, if we define the polynomial

$$
\begin{equation*}
Q_{k}(y):=P_{k}\left(x_{j}+h y\right)=f_{h}(k-r)+\sum_{i=1}^{r} f_{h}[(k-r), \ldots,(k-r+i)] \frac{q_{i-1}(y)}{i!}, \tag{68}
\end{equation*}
$$

where $q_{i}(y)=(y-(k-r)) \cdots(y-(k-r-i)), f_{h}(y)=f\left(x_{j}+h y\right)$, then we can rewrite

$$
\begin{equation*}
P_{k}^{(l)}(x)=\frac{d^{l}}{d x^{l}}\left(Q_{k}\left(\frac{x}{h}\right)\right)=\frac{1}{h^{l}} Q_{k}^{(l)}(y), \quad l=1, \ldots, r . \tag{69}
\end{equation*}
$$

Then, applying the change of variable $y=\left(x-x_{j}\right) / h$ in the integral in (33), we have

$$
\begin{equation*}
h^{2 l-3} \int_{x_{j-1}}^{x_{j}}\left(P_{k}^{(l)}(x)\right)^{2} d x=h^{-2} \int_{-1}^{0}\left(Q_{k}^{(l)}(y)\right)^{2} d y \tag{70}
\end{equation*}
$$

where

$$
\begin{equation*}
Q_{k}^{(l)}(y)=\sum_{i=l}^{r} f_{h}[(k-r), \ldots,(k-r+i)] \frac{q_{i-1}^{(l)}(y)}{i!} \tag{71}
\end{equation*}
$$

Furthermore, keeping in mind (67), we can write

$$
\begin{aligned}
Q_{k}^{(l)}(y)= & f_{h}[(k-r), \ldots,(k-r+l)] \frac{q_{l-1}^{(l)}(y)}{l!}+ \\
& +\sum_{i=l+1}^{r} f_{h}[(k-r), \ldots,(k-r+i)] \frac{q_{i-1}^{(l)}(y)}{i!} \\
= & h^{l} f^{(l)}\left(x_{j}\right)+o\left(h^{l}\right),
\end{aligned}
$$

having noticed that $q_{l-1}^{(l)}(y)=l$ !. From we have just seen, it follows

$$
\begin{aligned}
h^{2 l-3} \int_{x_{j-1}}^{x_{j}}\left(P_{k}^{(l)}(x)\right)^{2} d x & =h^{-2} \int_{-1}^{0}\left(Q_{k}^{(l)}(y)\right)^{2} d y \\
& =h^{2 l-2}\left(f^{(l)}\left(x_{j}\right)\right)^{2}+o\left(h^{2 l-2}\right)
\end{aligned}
$$

as we wanted. In fact, having in mind that by hypothesis $f^{\prime}, f^{\prime \prime} \neq 0$,

$$
\begin{aligned}
\beta_{k} & =\sum_{l=2}^{r} \int_{x_{j-1}}^{x_{j}} h^{2 l-3}\left(P_{k}^{(l)}\right)^{2} d x \\
& =\sum_{l=2}^{r} h^{2 l-2}\left(f^{(l)}\left(x_{j}\right)\right)^{2}+o\left(h^{2 l-2}\right)=O\left(h^{2}\right) .
\end{aligned}
$$

Remark A.1. To prove this point, it could be sufficient to observe that using the regularity of $f$ in $\Omega \backslash\left\{x_{s}\right\}$ and the properties of the interpolating polynomial (which we have just proved) we get

$$
P_{k}^{(l)}(x)=f^{(l)}(x)+O\left(h^{r+1-l}\right), \quad \text { for } x_{j-1} \leq x \leq x_{j}, k=0, \ldots, r-1 .
$$

Moreover, expanding with Taylor, it holds

$$
\begin{equation*}
f^{(l)}(x)=f^{(m)}\left(x_{j}\right) O(h)^{m-l}+o\left(h^{m-l}\right), \tag{72}
\end{equation*}
$$

where $m=\max \{s+1, l\}$ and $s=\max \left\{k: f^{(i)}\left(x_{j}\right)=0, \forall i \leq k\right\} \quad(s \leq r)$. Then, integrating (remembering that by hypothesis $s=0 \Rightarrow m=l$ ),

$$
h^{2 l-3} \int_{x_{j-1}}^{x_{j}}\left(P_{k}^{(l)}(x)\right)^{2} d x=h^{2 l-2}\left(f^{(l)}\left(x_{j}\right)\right)^{2}+o\left(h^{2 l-2}\right)
$$

as before.
Let us continue with the proof of $i \boldsymbol{i}$ ). In this case the proof is a little more complicated and it is better to treat separately the following two cases:
(a) 0 is a point of the grid $\left\{x_{i}\right\}, i \in \mathbb{Z}$;
(b) $0 \in I_{i}=\left(x_{i-1}, x_{i}\right)$ for some $i \in \mathbb{Z}$.

Case $a$. By hypothesis $0 \in \mathcal{S}_{j+k}$ for at least one $k=0, \ldots, r-1$, then, for each fixed $k$, there exists an integer $j_{s} \in\{k-r+1, \ldots, k-1\}$ such that $x_{j}=-j_{s} h$ (for $j_{s}=k-r$ and $j_{s}=k$ we fall in the case treated previously). Substituting in (65),

$$
P_{k}(x)=f\left(\left(-j_{s}+k-r\right) h\right)+\sum_{i=1}^{r} f\left[\left(-j_{s}+k-r\right) h, \ldots,\left(-j_{s}+k-r+i\right) h\right] \omega_{i-1}(x),
$$

with $\omega_{i}(x)=\left(x+\left(j_{s}-k+r\right) h\right) \cdots\left(x+\left(j_{s}-k+r-i\right) h\right)$. As we have done in the proof of $\left.\boldsymbol{i}\right)$, if we define the function $f_{h}(y):=f\left(x_{j}+h y\right)=f\left(h\left(y-j_{s}\right)\right)$, we can obtain the relations (69)-70) with $Q_{k}^{(l)}(y)$ defined as in 71 .

At this point, it is useful to notice that (66) for $l=1$ reads, for $i=1, \ldots, r$,

$$
f_{h}[k-r, \ldots,(k-r+i)]=\sum_{j=0}^{i-1}\binom{i-1}{j}(-1)^{i-j-1} f_{h}[(k-r+j),(k-r+j+1)] .
$$

In order to simplify the notation let us call $i_{s}:=j_{s}-k+r$, that is to say the index $i_{s} \in\{1, \ldots, r-1\}$ such that $x_{j}+\left(k-r+i_{s}\right) h=0$. Then, by hypothesis, we can write for all $i>t:=\max \left\{i_{s}, l-1\right\}$,

$$
\begin{aligned}
f_{h}[k-r, \ldots,(k-r+i)]= & \sum_{j=0}^{i_{s}-1}\binom{i-1}{j}(-1)^{i-j-1} f_{h}[(k-r+j),(k-r+j+1)] \\
& +\sum_{j=i_{s}}^{i-1}\binom{i-1}{j}(-1)^{i-j-1} f_{h}[(k-r+j),(k-r+j+1)]
\end{aligned}
$$

and, noticing that for $h \rightarrow 0$

$$
f_{h}\left[z+j_{s}, z+j_{s}+1\right]=h \frac{f((z+1) h)-f(z h)}{h} \rightarrow \begin{cases}h f^{\prime}\left(0^{+}\right) & \text {if } z \geq 0 \\ h f^{\prime}\left(0^{-}\right) & \text {otherwise }\end{cases}
$$

we can conclude that

$$
\begin{aligned}
f_{h}[k-r, \ldots,(k-r+i)] & \approx h\left[\sum_{j=0}^{i_{s}-1}\binom{i-1}{j}(-1)^{i-j-1} f^{\prime}\left(0^{-}\right)+\sum_{j=i_{s}}^{i-1}\binom{i-1}{j}(-1)^{i-j-1} f^{\prime}\left(0^{+}\right)\right] \\
& =h\left[f^{\prime}\left(0^{+}\right)-f^{\prime}\left(0^{-}\right)\right] \sum_{j=i_{s}}^{i-1}\binom{i-1}{j}(-1)^{i-j-1} \\
& =h\left[f^{\prime}\left(0^{+}\right)-f^{\prime}\left(0^{-}\right)\right]\binom{i-2}{i_{s}-1}(-1)^{i-i_{s}+1} \neq 0,
\end{aligned}
$$

having exploited the relations $\sum_{j=0}^{i}\binom{i}{j}(-1)^{i-j}=0$ and $\sum_{j=0}^{n}\binom{i}{j}(-1)^{i-j}=\binom{i-1}{n}(-1)^{i-n}$, for $0 \leq$ $n<i$ by Lemma A.1.

Furthermore for $l \leq i \leq i_{s}$, as we have seen in the first point $i$ ) of the proof,

$$
f_{h}[k-r, \ldots,(k-r+i)] \approx h \sum_{j=0}^{i-1}\binom{i-1}{j}(-1)^{i-j-1} f^{\prime}\left(0^{-}\right)=o(h),
$$

which also follows directly from the first relation we have just reminded.
From what we have done so far we can deduce, recalling that $t:=\max \left\{i_{s}, l-1\right\}$,

$$
\begin{aligned}
h^{2 l-3} \int_{x_{j-1}}^{x_{j}}\left(P_{k}^{(l)}(x)\right)^{2} d x & =h^{-2} \int_{-1}^{0}\left(\sum_{i=t+1}^{r} f_{h}[k-r, \ldots,(k-r+i)] \frac{q_{i-1}^{(l)}}{i!}\right)^{2} \\
& \approx C_{r k}\left[f^{\prime}\left(0^{+}\right)-f^{\prime}\left(0^{-}\right)\right]^{2}
\end{aligned}
$$

where $C_{r k}=\int_{-1}^{0}\left(\sum_{i=t+1}^{r}\binom{i-2}{i_{s}-1}(-1)^{i-i_{s}+1} \frac{q_{i-1}^{(l)}}{i!}\right)^{2}$, which is the thesis for case $a$.
Case b. By hypothesis there exists an integer $j_{s} \in\{k-r+1, \ldots, k\}$ and a number $0<a_{s}<1$ such that $x_{j}=\left(-j_{s}+a_{s}\right) h$. It is clear now that we can repeat the same constructions of the previous case defining the function $f_{a_{s}, h}(y):=f\left(h\left(y-j_{s}+a_{s}\right)\right)$ and using it in place of $f_{h}$; so, to obtain 70) it will suffice to apply the change of variables $y=\frac{x}{h}+j_{s}-a_{s}$. Then, naming $i_{s}=j_{s}-k+r$, for $i \geq t:=\max \left\{i_{s}, l\right\}$,

$$
\begin{align*}
f_{a_{s}, h}[k-r, \ldots,(k-r+i)]= & \sum_{j=0}^{i_{s}-2}\binom{i-1}{j}(-1)^{i-j-1} f_{a_{s}, h}[(k-r+j),(k-r+j+1)] \\
& +\binom{i-1}{i_{s}-1}(-1)^{i-i_{s}} f_{a_{s}, h}\left[j_{s}-1, j_{s}\right] \\
& +\sum_{j=i_{s}}^{i-1}\binom{i-1}{j}(-1)^{i-j-1} f_{a_{s}, h}[(k-r+j),(k-r+j+1)], \tag{73}
\end{align*}
$$

whence, noticing that

$$
\begin{aligned}
f_{a_{s}, h}\left[j_{s}-1, j_{s}\right] & =f\left(a_{s} h\right)-f\left(\left(a_{s}-1\right) h\right) \\
& =a_{s} h\left(\frac{f\left(a_{s} h\right)-f(0)}{a_{s} h}\right)+\left(1-a_{s}\right) h\left(\frac{f(0)-f\left(\left(a_{s}-1\right) h\right)}{\left(1-a_{s}\right) h}\right) \\
& \approx a_{s} h f^{\prime}\left(0^{+}\right)+\left(1-a_{s}\right) h f^{\prime}\left(0^{-}\right) \\
& =a_{s} h\left[f^{\prime}\left(0^{+}\right)-f^{\prime}\left(0^{-}\right)\right]+h f^{\prime}\left(0^{-}\right),
\end{aligned}
$$

and that

$$
f_{a_{s}, h}\left[z+j_{s}-1, z+j_{s}\right] \rightarrow \begin{cases}h f^{\prime}\left(0^{+}\right) & \text {if } z \geq 1 \\ h f^{\prime}\left(0^{-}\right) & \text {if } z \leq-1,\end{cases}
$$

we can infer that if $i=i_{s}$ (in this case in (73) on the right side of the equation we have only the second term), then $f_{a_{s}, h}[k-r, \ldots,(k-r+i)] \approx a_{s} h\left[f^{\prime}\left(0^{+}\right)-f^{\prime}\left(0^{-}\right)\right] \neq 0$, while if $i>i_{s}$,

$$
f_{a_{s}, h}[k-r, \ldots,(k-r+i)] \approx h\left[f^{\prime}\left(0^{+}\right)-f^{\prime}\left(0^{-}\right)\right]\left[\binom{i-2}{i_{s}-1}(-1)^{i-i_{s}+1}+a_{s}\binom{i-1}{i_{s}-1}(-1)^{i-i_{s}}\right] .
$$

The last quantity, as it can be easily shown, it is null if and only if $a_{s}=\frac{i-i_{s}}{i-1}$; more precisely, for $k$ fixed there exists an integer $i \geq t$ such that $f_{a_{s}, h}[k-r, \ldots,(k-r+i)] \approx C h\left[f^{\prime}\left(0^{+}\right)-f^{\prime}\left(0^{-}\right)\right]$with $C \neq 0$, whence the thesis even in the last case.

Lemma A.1. Let us assume $i \geq 1$ and write $f[\cdot]$ for the undivided difference of a function $f$. Then, it holds

$$
\begin{equation*}
f[0, \ldots, i]=\sum_{j=0}^{i-l}\binom{i-l}{j}(-1)^{i-l-j} f[j, \ldots, j+l], \quad \text { for } l=0, \ldots, i . \tag{74}
\end{equation*}
$$

Moreover, we have that

$$
\sum_{j=0}^{n}\binom{i}{j}(-1)^{i-j}= \begin{cases}\binom{i-1}{n}(-1)^{i-n} & \text { for } n<i  \tag{75}\\ 0 & \text { for } n=i\end{cases}
$$

Proof. Let us start from the proof of (74) and let us proceed by induction on $i$.
Firstly, let us notice that for $l=i$ the identity is trivially satisfied, whence the case $i=0$ follows directly. Then, for any $l=0, \ldots, i-1$, suppose that the statement holds for $i-1$ and for $i>0$ let us compute,

$$
\begin{array}{rlr}
f[0, \ldots, i]= & f[1, \ldots, i]-f[0, \ldots, i-1] & \text { by definition of } f[\cdot] \\
= & \sum_{j=0}^{i-l-1}\binom{i-l-1}{j}(-1)^{i-l-1-j} f[j+1, \ldots, j+1+l] \\
& -\sum_{j=0}^{i-l-1}\binom{i-l-1}{j}(-1)^{i-l-1-j} f[j, \ldots, j+l] \quad \text { by inductive hyp. } \\
= & f[i-l, \ldots, i]+(-1)^{i-l} f[0, \ldots, l] \\
& +\sum_{j=1}^{i-l-1}\binom{i-l-1}{j-1}(-1)^{i-l-j} f[j, \ldots, j+l] \\
& +\sum_{j=1}^{i-l-1}\binom{i-l-1}{j}(-1)^{i-l-j} f[j, \ldots, j+l] \\
= & f[i-l, \ldots, i]+(-1)^{i-l} f[0, \ldots, l] \\
& +\sum_{j=1}^{i-l-1}\binom{i-l}{j}(-1)^{i-l-j} f[j, \ldots, j+l] \\
= & \sum_{j=0}^{i-l}\binom{i-l}{j}(-1)^{i-l-j} f[j, \ldots, j+l], & \binom{n-1}{k-1}+\binom{n-1}{k}=\binom{n}{k}
\end{array}
$$

as we wanted.
Remark A.2. To simplify the notation we have stated the result for $f[0, \ldots, i]$ but the proof clearly holds for $f[k, \ldots, k+i], \forall k$. In the second identity of the previous chain we have assumed this fact applying the inductive hypothesis on both terms.

Let us focus now on the second relation of the lemma (75) and proceed again by induction, but this time on $n: 0 \leq n<i$. For $n=0$ we have $(-1)^{i}=(-1)^{i}$, than the identity holds. Suppose that
(75) holds for $n-1<i-1$ and compute

$$
\begin{aligned}
\sum_{j=0}^{n}\binom{i}{j}(-1)^{i-j} & =\sum_{j=0}^{n-1}\binom{i}{j}(-1)^{i-j}+\binom{i}{n}(-1)^{1-n} \\
& =\binom{i-1}{n-1}(-1)^{i+1-n}+\binom{i}{n}(-1)^{i-n} \quad \text { by inductive hyp. } \\
& =\binom{i-1}{n-1}(-1)^{i+1-n}-\left[\binom{i-1}{n}+\binom{i-1}{n-1}\right](-1)^{i+1-n} \\
& =\binom{i-1}{n}(-1)^{i-n} .
\end{aligned}
$$

For $n=i$ instead, from what we have just seen we can easily compute

$$
\begin{aligned}
\sum_{j=0}^{i}\binom{i}{j}(-1)^{i-j} & =\sum_{j=0}^{i-1}\binom{i}{j}(-1)^{i-j}+(-1)^{i-i} \\
& =\binom{i-1}{i-1}(-1)^{i-i+1}+1 \\
& =-1+1=0 .
\end{aligned}
$$

For the last result is better to prove first the following technical lemma.
Lemma A.2. Let $x \in \mathbb{R}$ and $n \in \mathbb{N}, n \geq 1$. Then, for $t \in \mathbb{N}, t \geq 1$

$$
\begin{equation*}
(x+n)^{t}=\sum_{j=0}^{t-1}(x+j)(x+n)^{t-j-1}(n)_{j}+(n)_{t}, \tag{76}
\end{equation*}
$$

where $(x)_{k}=x(x-1) \cdots(x-k+1)$ is the falling factorial $\left((x)_{0}=1\right)$. Whence it follows directly that, for $i, t \in \mathbb{N}: i \geq 1$,

$$
\sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n}(x+n)^{t}= \begin{cases}0 & \text { for } 0 \leq t<i  \tag{77}\\ i! & \text { for } t=i\end{cases}
$$

Moreover, for $t>i$,

$$
\sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n}(x+n)^{t}=\sum_{j=0}^{\min \{t-i, i\}}(x+j)(i)_{j} \sum_{n=0}^{i-j}\binom{i-j}{n}(-1)^{i-n-j}((x+j)+n)^{t-1-j}
$$

Proof. Let us proceed again by induction, this time on $t$. For $t=1$ we have $x+n=x+n$ so the
identity is verified. Then suppose the (76) true for $t>1$ and compute

$$
\begin{aligned}
(x+n)^{t+1} & =(x+n)^{t}(x+n) \\
& =\left[\sum_{j=0}^{t-1}(x+j)(x+n)^{t-1-j}(n)_{j}+(n)_{t}\right](x+n) \quad \text { by inductive hyp. } \\
& =\sum_{j=0}^{t-1}(x+j)(x+n)^{t-j}(n)_{j}+(n)_{t}(x+t+n-t) \\
& =\sum_{j=0}^{t}(x+j)(x+n)^{t-j}(n)_{j}+(n)_{t+1}
\end{aligned}
$$

as we wanted.
Let us pass to the relation (77) and proceed by induction on $t<i$. For $t=0$ the identity holds by Lemma A.1. Suppose that (77) holds for $t-1<i-1$ and compute

$$
\begin{aligned}
\sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n}(x+n)^{t}= & \sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n}\left[\sum_{j=0}^{t-1}(x+j)(x+n)^{t-1-j}(n)_{j}+(n)_{t}\right] \\
= & \sum_{j=0}^{t-1}(x+j)(i)_{j} \sum_{n=j}^{i}\binom{i-j}{n-j}(-1)^{i-n}(x+n)^{t-1-j} \\
& +(i)_{t} \sum_{n=t}^{i}\binom{i-t}{n-t}(-1)^{i-n} \\
= & \sum_{j=0}^{t-1}(x+j)(i)_{j} \underbrace{\sum_{n=0}^{i-j}\binom{i-j}{n}(-1)^{i-n-j}((x+j)+n)^{t-1-j}}_{=0 \text { by inductive hypothesis }} \\
& +(i)_{t} \underbrace{\sum_{n=0}^{i-t}\binom{i-t}{n}(-1)^{i-n-t}}_{=0}=0 .
\end{aligned}
$$

Finally, for $t=i$ and repeating the same computations we notice that the last term it is not null but

$$
(i)_{i} \sum_{n=i}^{i}\binom{i-i}{n-i}=(i)_{i}=i(i-1) \cdots(i-i+1)=i!
$$

as we wanted.
For the last relation it suffices to repeat again the same expansions we have just done and notice that if $j>i$, then $(n)_{j}=0$ for $n=0, \ldots, i$, and moreover if $t-i<i$,

$$
\sum_{n=0}^{i-j}\binom{i-j}{n}(-1)^{i-n-j}((x+j)+n)^{t-1-j}=0 \quad \text { per } t-i<j \leq i
$$

for what we have seen in the previous case.

Corollary A.3. Let $f \in C^{r+1}\left(\left[x_{j+k-r}, x_{j+k}\right]\right)$ for $k \in\{0, \ldots, r-1\}, r>1$. Then, for $i=1, \ldots, r$ e $s=0, \ldots, i$,

$$
\begin{equation*}
f_{h}[(k-r+s), \ldots,(k-r+s+i)]=h^{i} f^{(i)}\left(x_{j}\right)+o\left(h^{i}\right) . \tag{78}
\end{equation*}
$$

Proof. Defining $x=k-r+s$, if we write the (74) for $l=0$ we have

$$
\begin{aligned}
f_{h}[x, \ldots, x+i] & =\sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n} f_{h}[x+n] \\
& =\sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n} f\left(x_{j}+(x+n) h\right) .
\end{aligned}
$$

Then, developing with Taylor up to order $r$ we can write

$$
f\left(x_{j}+(x+n) h\right)=\sum_{t=0}^{r} f^{(t)}\left(x_{j}\right) \frac{[h(x+n)]^{t}}{t!}+o\left(h^{r}\right)
$$

whence

$$
\begin{aligned}
f_{h}[x, \ldots, x+i] & =\sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n} \sum_{t=0}^{r} f^{(t)}\left(x_{j}\right) \frac{[h(x+n)]^{t}}{t!}+o\left(h^{r}\right) \\
& =\sum_{t=0}^{r} f^{(t)}\left(x_{j}\right) \frac{h^{t}}{t!} \sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n}(x+n)^{t}+o\left(h^{r}\right) \\
& =h^{i} f^{(i)}\left(x_{j}\right)+\sum_{t=i+1}^{r} f^{(t)}\left(x_{j}\right) \frac{h^{t}}{t!} \sum_{n=0}^{i}\binom{i}{n}(-1)^{i-n}(x+n)^{t}+o\left(h^{r}\right) \\
& =h^{i} f^{(i)}\left(x_{j}\right)+o\left(h^{i}\right)
\end{aligned}
$$

having exploited (77) in the third identity of the chain.


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