WEIGHTED ESTIMATES FOR POWERS AND SMOOTHING ESTIMATES OF SCHRÖDINGER OPERATORS WITH INVERSE-SQUARE POTENTIALS

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ABSTRACT. Let \mathcal{L}_a be a Schrödinger operator with inverse square potential $a|x|^{-2}$ on \mathbb{R}^d , $d \geq 3$. The main aim of this paper is to prove weighted estimates for fractional powers of \mathcal{L}_a . The proof is based on weighted Hardy inequalities and weighted inequalities for square functions associated to \mathcal{L}_a . As an application, we obtain smoothing estimates regarding the propagator $e^{it\mathcal{L}_a}$.

1. INTRODUCTION

In this paper, we consider the following Schrödinger operators with inverse-square potentials on $\mathbb{R}^d, d \geq 3$,

(1)
$$\mathcal{L}_a = -\Delta + \frac{a}{|x|^2} \quad \text{with} \quad a \ge -\left(\frac{d-2}{2}\right)^2.$$

 Set

$$\sigma := \frac{d-2}{2} - \frac{1}{2}\sqrt{(d-2)^2 + 4a}.$$

The Schrödinger operator \mathcal{L}_a is understood as the Friedrichs extension of $-\Delta + \frac{a}{|x|^2}$ defined initially on $C_c^{\infty}(\mathbb{R}^d \setminus \{0\})$. The condition $a \ge -\left(\frac{d-2}{2}\right)^2$ guarantees that \mathcal{L}_a is nonnegative. It is wellknown that \mathcal{L}_a is self-adjoint and the extension may not be unique as $-\left(\frac{d-2}{2}\right)^2 < a < 1 - \left(\frac{d-2}{2}\right)^2$. For further details, we refer the readers to [23, 24, 21, 30, 32, 26].

It is well known that Schrödinger operators with inverse-square potentials \mathcal{L}_a have a wide range of applications in physics and mathematics spanning areas such as combustion theory, the Dirac equation with Coulomb potential, quantum mechanics and the study of perturbations of classic space-time metrics. See for example [4, 5, 33, 21] and their references.

Recently there has been a spate of activity dedicated to the operator \mathcal{L}_a . Strichartz estimates, which are an effective tool for studying the behavior of solutions to nonlinear Schrödinger equations and wave equations related to \mathcal{L}_a , were investigated in [4, 5]. In [17] the authors developed the study of Strichartz estimates for the propagators $e^{it(\Delta+V)}$ with $V(x) \sim |x|^{-2}$. The well-posedness and behaviour of the solutions to the heat equation related to \mathcal{L}_a was studied in [33]. In [34], using Morawetz-type inequalities and Sobolev norm properties related to \mathcal{L}_a , the long-time behavior of solutions to nonlinear Schrödinger equations associated to \mathcal{L}_a was considered. More recently, the authors in [23] established the equivalence between L^p -based Sobolev norms defined in terms of $\mathcal{L}_a^{s/2}$ and in terms of $(-\Delta)^{s/2}$ for all regularities 0 < s < 2.

In this paper, our first objective is to extend the estimates in [23] to weighted estimates. More precisely, we will prove the following result.

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Theorem 1.1. Suppose that $d \ge 3$, $a \ge -\left(\frac{d-2}{2}\right)^2$ and 0 < s < 2. If $r_1 := 1 \lor \frac{d}{d-\sigma} (where <math>a \lor b = \max\{a, b\}$) with convention $\frac{d}{0} = \infty$ then for $w \in A_{p/r_1} \cap RH_{(r_2/p)'}$ we have

(2)
$$\|(-\Delta)^{s/2}f\|_{L^p_w} \lesssim \|\mathcal{L}^{s/2}_a f\|_{L^p_w}.$$

If $1 with <math>p_1 := 1 \lor \frac{d}{d-\sigma} then for <math>w \in A_{p/p_1} \cap RH_{(p_2/p)'}$ we have

(3)
$$\|\mathcal{L}_{a}^{s/2}f\|_{L_{w}^{p}} \lesssim \|(-\Delta)^{s/2}f\|_{L_{w}^{p}}$$

Let us describe the motivation for the results in Theorem 1.1.

- (i) When s = 1, (2) and (3) are known as the boundedness of the Riesz transforms and the reverse Riesz transforms, respectively. Note that the boundedness of the Riesz transforms related to \mathcal{L}_a was obtained in [18]. Hence, Theorem 1.1 can be considered as a natural outgrowth of this direction of research.
- (ii) The second motivation of our present work is the need of the following estimate of the form:

$$|x|^{\beta} \mathcal{L}_a^{\theta} f \|_{L^2} \lesssim ||x|^{\beta} (-\Delta)^{\theta} f \|_{L^2}$$

for certain β and θ . This type of estimate was studied in [6] for certain Schrödinger operators instead of \mathcal{L}_a and played a key role in studying dispersive properties of Schrödinger equations on non-flat waveguides. See also [9] for related weighted estimates in L^p spaces with mixed radial-angular integrability.

(iii) Another motivation of Theorem 1.1 is its utility in obtaining smoothing estimates related to the propagators $e^{it\mathcal{L}_a}$. We give such estimates in Theorem 1.2 below. Note that smoothing estimates related to Schrödinger operators are a topic of interest in PDEs and have a close relationship to Strichartz estimates. For further details, the reader can consult [31, 11, 19, 22, 27 and the references therein.

As an application of Theorem 1.1, we obtain the following smoothing estimates.

Theorem 1.2. Suppose that $d \ge 3$, $a > -\left(\frac{d-2}{2}\right)^2 + \frac{1}{4}$, and consider the Schrödinger flow $e^{it\mathcal{L}_a}f$. Then for all $0 < \epsilon < 1$ we have the following smoothing estimates, with C independent of ϵ :

(4)
$$\int \int \left[\frac{|x|^{\epsilon-1}}{(1+|x|^{\epsilon})^2} |\nabla e^{it\mathcal{L}_a} f|^2 + \frac{|x|^{\epsilon-3}}{1+|x|^{\epsilon}} |e^{it\mathcal{L}_a} f|^2 \right] dx dt \le C\epsilon^{-1} \|(-\Delta)^{1/4} f\|_{L^2}^2$$

and also

(5)
$$\|w(x)^{1/2}(-\Delta)^{1/4}e^{it\mathcal{L}_a}f\|_{L^2(\mathbb{R}^{d+1})} \lesssim C\epsilon^{-1/2}\|f\|_{L^2}, \qquad w(x) = \frac{|x|^{\epsilon-1}}{(1+|x|^{\epsilon})^2}.$$

On the other hand, for the wave flow $e^{it\mathcal{L}_a^{1/2}}f$ we have the estimate

(6)
$$\|w(x)^{1/2}e^{it\mathcal{L}_a^{1/2}}f\|_{L^2(\mathbb{R}^{d+1})} \lesssim C\epsilon^{-1/2}\|f\|_{L^2}, \qquad w(x) = \frac{|x|^{\epsilon-1}}{(1+|x|^{\epsilon})^2}.$$

We have a few comments on the theorem. The condition $a > -\left(\frac{d-2}{2}\right)^2 + \frac{1}{4}$ guarantees that the weight w satisfies the conditions in Theorem 1.1. It is an open question whether the results in Theorem 1.2 holds true for $a > -\left(\frac{d-2}{2}\right)^2$. In the particular case of d = 3, the potential is repulsive which is considered in [15, 2]; however, our results are new for d > 3. PLEASE CHECK!

To prove Theorem 1.1 although we follow the approach in [23], some significant modifications and improvements are required due to the following reasons. The first reason is that we work on the weighted Lebesgue estimates instead of unweighted estimates. The second one we need to point out is that in our present paper we employ the vertical square functions in place of the (discrete) Littlewood–Paley square functions. This allows us to bypass the use of spectral multipliers as in [23]. Moreover, due to the lack of regularity condition of the heat kernels of $e^{-t\mathcal{L}_a}$, certain singular integrals considered in the paper may be beyond the Calderón-Zygmund theory. This causes some challenging matters, and we overcome these problems by using the criteria established in [1, 3] for a sigular integrals to be bounded on weighted Lebesgue spaces.

The organization of the paper is as follows. In Section 2 we recall some preliminaries on the Muckenhoupt weights and two criteria for a sigular integrals to be bounded on weighted and unweighted Lebesgue spaces. Some kernels estimates will be derived in Section 3. In Section 4, we first prove the weighted Hardy inequality and weighted estimates for square functions related to \mathcal{L}_a which are of interest in their own right. We conclude Section 4 by using these results to prove Theorems 1.1 and 1.2.

Throughout the paper, we always use C and c to denote positive constants that are independent of the main parameters involved but whose values may differ from line to line. We will write $A \leq B$ if there is a universal constant C so that $A \leq CB$ and $A \sim B$ if $A \leq B$ and $B \leq A$. For $a, b \in \mathbb{R}$, we denote $a \lor b = \max\{a, b\}$ and $a \land b = \min\{a, b\}$. For $p \in [1, \infty]$, we denote by $p' = \frac{p}{p-1}$ the conjugate exponent of p.

2. Preliminaries

2.1. Muckenhoupt weights. We start with some notations which will be used frequently. For a measurable subset $E \subset \mathbb{R}^d$ and a measurable function f we denote

$$\int_E f(x)dx = \frac{1}{|E|} \int_E f(x)dx.$$

Given a ball B, we denote $S_j(B) = 2^j B \setminus 2^{j-1} B$ for $j = 1, 2, 3, \ldots$, and we set $S_0(B) = B$.

Let $1 \leq q < \infty$. A nonnegative locally integrable function w belongs to the Muckenhoupt class A_q , say $w \in A_q$, if there exists a positive constant C so that

$$\left(\int_B w(x)dx\right)\left(\int_B w^{-1/(q-1)}(x)dx\right)^{q-1} \le C, \quad \text{if } 1 < q < \infty,$$

and

$$\label{eq:product} \int_B w(x) dx \leq C \operatorname{ess-inf}_{x \in B} w(x), \quad \text{if } q = 1,$$

for all balls B in \mathbb{R}^d . We say that $w \in A_\infty$ if $w \in A_q$ for any $q \in [1,\infty)$. We shall denote $w(E) := \int_E w(x) dx$ for any measurable set $E \subset \mathbb{R}^d$.

The reverse Hölder classes are defined in the following way: $w \in RH_r, 1 < r < \infty$, if there is a constant C such that for any ball $B \subset \mathbb{R}^d$,

$$\left(\int_B w^r(x)dx\right)^{1/r} \le C \int_B w(x)dx.$$

The endpoint $r = \infty$ is given by the condition: $w \in RH_{\infty}$ whenever, there is a constant C such that for any ball $B \subset \mathbb{R}^d$,

$$w(x) \le C \oint_B w(y) dy$$
 for a.e. $x \in B$.

Let $w \in A_{\infty}$. For $0 , the weighted space <math>L^p_w(\mathbb{R}^d)$ is defined as the space of w(x)dx-measurable functions f such that

$$||f||_{L^p_w(\mathbb{R}^d)} := \left(\int_{\mathbb{R}^d} |f(x)|^p w(x) dx\right)^{1/p} < \infty.$$

It is well-known that the power weight $w(x) = |x|^{\alpha} \in A_p$ if and only if $-d < \alpha < d(p-1)$. Moreover, $w(x) = |x|^{\alpha} \in RH_q$ if and only if $\alpha q > -d$.

We sum up some of the properties of Muckenhoupt classes and reverse Hölder classes in the following results. See [12, 20].

Lemma 2.1. The following properties hold:

- (i) $w \in A_p, 1 if and only if <math>w^{1-p'} \in A_{p'}$.
- (ii) $A_1 \subset A_p \subset A_q$ for $1 \le p \le q \le \infty$.
- (iii) $RH_{\infty} \subset RH_q \subset RH_p$ for 1 .
- (iv) If $w \in A_p$, 1 , then there exists <math>1 < q < p such that $w \in A_q$.
- (v) If $w \in RH_q$, $1 < q < \infty$, then there exists $q such that <math>w \in RH_p$.
- (vi) $A_{\infty} = \bigcup_{1 \le p < \infty} A_p = \bigcup_{1 < p \le \infty} RH_p.$
- (vii) Let $1 < p_0 < p < q_0 < \infty$. Then we have

$$w \in A_{\frac{p}{p_0}} \cap RH_{(\frac{q_0}{p})'} \Longleftrightarrow w^{1-p'} \in A_{\frac{p'}{q'_0}} \cap RH_{(\frac{p'_0}{p'})'}.$$

2.2. Hardy-Littlewood maximal functions. For r > 0, the Hardy-Littlewood maximal function \mathcal{M}_r is defined by

$$\mathcal{M}_r f(x) = \sup_{B \ni x} \left(\frac{1}{|B|} \int_B |f(y)|^r \, dy \right)^{1/r}, \ x \in \mathbb{R}^d,$$

where the supremum is taken over all balls B containing x. When r = 1, we write \mathcal{M} instead of \mathcal{M}_1 .

We now record the following results concerning the weak type estimates and the weighted estimates of the maximal functions.

Lemma 2.2. Let $0 < r < \infty$. Then we have for p > r and $w \in A_{p/r}$,

$$\|\mathcal{M}_r f\|_{L^p_w} \lesssim \|f\|_{L^p_w}$$

2.3. Two theorems on the boundedness of singular integrals. We recall the definition of linearizable operators in [16]. An operator T defined on $L^2(\mathbb{R}^d)$ is said to be a *linearizable operator* if there exists a Banach space \mathbb{B} and a linear operator U from $L^2(\mathbb{R}^n)$ into $L^2(\mathbb{R}^d, \mathbb{B})$ so that

$$|Tf(x)| = ||Uf(x)||_{\mathbb{B}}$$

for all $f \in L^2(\mathbb{R}^d)$ and a.e. $x \in \mathbb{R}^d$.

It can be verified that a linearizable operator is a sublinear operator. The class of linearizable operator includes linear operators, maximal operators and square functions.

We first recall a theorem which is taken from [3, Theorem 6.6] on a criterion for the singular integrals to be bounded on the weighted Lebesgue spaces.

Theorem 2.3. Let $1 \le p_0 < q_0 \le \infty$ and let T be L^{q_0} bounded linearizable operator. Assume that there exists a family of operators $\{A_t\}_{t>0}$ satisfying that for $j \ge 2$ and every ball B

(7)
$$\left(\int_{S_j(B)} |T(I - \mathcal{A}_{r_B})f|^{q_0} d\mu \right)^{1/q_0} \le \alpha(j) \left(\int_B |f|^{q_0} d\mu \right)^{1/q_0}$$

and

(8)
$$\left(\int_{S_j(B)} |\mathcal{A}_{r_B}f|^{q_0} d\mu\right)^{1/q_0} \le \alpha(j) \left(\int_B |f|^{p_0} d\mu\right)^{1/p_0}$$

for all f supported in B. If $\sum_{j} \alpha(j) 2^{jd} < \infty$, then T is bounded on $L^p_w(\mathbb{R}^d)$ for all $p \in (p_0, q_0)$ and $w \in A_{\frac{p}{2n}} \cap RH_{(\frac{q_0}{2n})'}$.

Note that [3, Theorem 6.6] proves Theorem 2.3 for $q_0 = 2$, but their arguments also work well for any value of q_0 .

The following theorem is a direct consequence of [1, Theorem 3.7] which give a sufficient conditions for a singular integral to be bounded on Lebesgue spaces which plays an important role in the sequel. **Theorem 2.4.** Let $1 \le p_0 < q_0 \le \infty$. Let T be a bounded sublinear operator on $L^{p_0}(\mathbb{R}^d)$. Assume that there exists a family of operators $\{\mathcal{A}_t\}_{t>0}$ satisfying that

(9)
$$\left(\int_{B} \left|T(I-\mathcal{A}_{r_{B}})f\right|^{p_{0}}dx\right)^{1/p_{0}} \leq C\mathcal{M}_{p_{0}}(f)(x),$$

and

(10)
$$\left(\int_{B} \left|T\mathcal{A}_{r_{B}}f\right|^{q_{0}}dx\right)^{1/q_{0}} \leq C\mathcal{M}_{p_{0}}(|Tf|)(x),$$

for all balls B with radius r_B , all $f \in C_c^{\infty}(\mathbb{R}^n)$ and all $x \in B$. Then T is bounded on $L^p(\mathbb{R}^d)$ for all $p_0 .$

3. Some kernel estimates

For a constant $\alpha \in \mathbb{R}$. We denote

$$d_{\alpha} = \begin{cases} \frac{d}{\alpha}, & \alpha > 0\\ \infty, & \alpha \le 0. \end{cases}$$

Theorem 3.1. Let $\{T_t\}_{t>0}$ be a family of linear operators on $L^2(\mathbb{R}^d)$ with their associated kernels $T_t(x, y)$. Assume that there exist C, c > 0 and $\alpha, \beta \in \mathbb{R}$ with $d'_{\beta} < d_{\alpha}$ such that for all t > 0 and $x, y \in \mathbb{R}^d \setminus \{0\}$,

(11)
$$|T_t(x,y)| \le C \left(1 + \frac{\sqrt{t}}{|x|}\right)^{\alpha} \left(1 + \frac{\sqrt{t}}{|y|}\right)^{\beta} t^{-d/2} e^{-\frac{|x-y|^2}{ct}}.$$

Assume that $d'_{\beta} . Then for every <math>t > 0$, any measurable subsets $E, F \subset \mathbb{R}^d$, and all $f \in L^p(E)$, we have:

(12)
$$\|T_t f\|_{L^q(F)} \le Ct^{-\frac{d}{2}(\frac{1}{p} - \frac{1}{q})} e^{-\frac{d(E,F)^2}{ct}} \|f\|_{L^p(E)}$$

To prove this theorem, we need the following elementary results.

Lemma 3.2. (a) Let $\kappa \in (-\infty, d)$. Then there exists C > 0 so that for all r > 0

$$\int_{B(0,r)} \frac{1}{|x|^{\kappa}} dx \le Cr^{d-\kappa}.$$

(b) For $p \in [1, \infty)$ we have

$$\Big(\int_{\mathbb{R}^d} \Big[\frac{1}{t^{d/2}}e^{-\frac{|x-y|^2}{ct}}\Big]^p dy\Big)^{1/p} \leq \frac{C}{t^{\frac{d}{2p'}}}$$

uniformly in $x \in \mathbb{R}^d$.

Proof. The proof of this lemma is simple and we omit details.

We now turn to prove Theorem 3.1.

Proof of Theorem 3.1: From (11), for $x \in F$ we have

$$\begin{split} \|T_t f\|_{L^q(F)} &\leq \left\{ \int_F \Big[\int_E \Big(1 + \frac{\sqrt{t}}{|x|} \Big)^{\alpha} \Big(1 + \frac{\sqrt{t}}{|y|} \Big)^{\beta} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy \Big]^q dx \right\}^{1/q} \\ &\leq \left\{ \int_{F \cap B(0,\sqrt{t})} \Big[\int_{E \cap B(0,\sqrt{t})} \Big(1 + \frac{\sqrt{t}}{|x|} \Big)^{\alpha} \Big(1 + \frac{\sqrt{t}}{|y|} \Big)^{\beta} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy \Big]^q dx \right\}^{1/q} \\ &\quad + \left\{ \int_{F \cap B(0,\sqrt{t})} \Big[\int_{E \setminus B(0,\sqrt{t})} \Big(1 + \frac{\sqrt{t}}{|x|} \Big)^{\alpha} \Big(1 + \frac{\sqrt{t}}{|y|} \Big)^{\beta} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy \Big]^q dx \right\}^{1/q} \\ &\quad + \left\{ \int_{F \setminus B(0,\sqrt{t})} \Big[\int_{E \cap B(0,\sqrt{t})} \Big(1 + \frac{\sqrt{t}}{|x|} \Big)^{\alpha} \Big(1 + \frac{\sqrt{t}}{|y|} \Big)^{\beta} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy \Big]^q dx \right\}^{1/q} \\ &\quad + \left\{ \int_{F \setminus B(0,\sqrt{t})} \Big[\int_{E \setminus B(0,\sqrt{t})} \Big(1 + \frac{\sqrt{t}}{|x|} \Big)^{\alpha} \Big(1 + \frac{\sqrt{t}}{|y|} \Big)^{\beta} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy \Big]^q dx \right\}^{1/q} \\ &\quad + \left\{ \int_{F \setminus B(0,\sqrt{t})} \Big[\int_{E \setminus B(0,\sqrt{t})} \Big(1 + \frac{\sqrt{t}}{|x|} \Big)^{\alpha} \Big(1 + \frac{\sqrt{t}}{|y|} \Big)^{\beta} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy \Big]^q dx \right\}^{1/q} \\ &\leq E_1 + E_2 + E_3 + E_4. \end{split}$$

By Hölder's inequality, Lemma 3.2 and the fact that $\beta p' < d$ we have

$$\begin{split} E_{1} &\leq C e^{-\frac{d(E,F)^{2}}{ct}} \Big[\int_{F \cap B(0,\sqrt{t})} t^{-qd/2} \Big(1 + \frac{\sqrt{t}}{|x|} \Big)^{q\alpha} dx \Big]^{1/q} \Big(\int_{E} |f|^{p} \Big)^{1/p} \Big[\int_{B(0,\sqrt{t})} \Big(1 + \frac{\sqrt{t}}{|y|} \Big)^{\beta p'} dy \Big]^{1/p'} \\ &\lesssim e^{-\frac{d(E,F)^{2}}{ct}} t^{-\frac{d}{2}(\frac{1}{p} - \frac{1}{q})} \Big(\int_{E} |f|^{p} \Big)^{1/p}. \end{split}$$

For the second term, using Hölder's inequality, Lemma 3.2 again with the fact that $\alpha q < d$ we arrive at

$$\begin{split} E_{2} &\lesssim \left\{ \int_{F \cap B(0,\sqrt{t})} \left[\int_{E \setminus B(0,\sqrt{t})} \left(1 + \frac{\sqrt{t}}{|x|} \right)^{\alpha} t^{-d/2} e^{-\frac{|x-y|^{2}}{ct}} |f(y)| dy \right]^{q} dx \right\}^{1/q} \\ &\lesssim e^{-\frac{d(E,F)^{2}}{2ct}} \left\{ \int_{F \cap B(0,\sqrt{t})} \left| \left(1 + \frac{\sqrt{t}}{|x|} \right)^{\alpha} \left(\int_{E} |f|^{p} \right)^{1/p} \left(\int_{\mathbb{R}^{d}} \left[t^{-d/2} e^{-\frac{|x-y|^{2}}{2ct}} \right]^{p'} dy \right)^{1/p'} \right|^{q} dx \right\}^{1/q} \\ &\lesssim e^{-\frac{d(E,F)^{2}}{2ct}} t^{-\frac{d}{2p}} \left(\int_{E} |f|^{p} \right)^{1/p} \left[\int_{F \cap B(0,\sqrt{t})} \left(1 + \frac{\sqrt{t}}{|x|} \right)^{\alpha q} dx \right]^{1/q} \\ &\lesssim e^{-\frac{d(E,F)^{2}}{2ct}} t^{-\frac{d}{2}(\frac{1}{p} - \frac{1}{q})} \left(\int_{E} |f|^{p} \right)^{1/p}. \end{split}$$

By a similar argument we can also dominate E_3 by

$$t^{-\frac{d}{2}(\frac{1}{p}-\frac{1}{q})}e^{-\frac{d(E,F)^2}{2ct}}\Big(\int_E |f|^p\Big)^{1/p}.$$

It remains to estimate the last term E_4 . We observe that

$$E_4 \lesssim \left\{ \int_{F \setminus B(0,\sqrt{t})} \left[\int_{E \setminus B(0,\sqrt{t})} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy \right]^q dx \right\}^{1/q}$$

At this stage, by using the standard argument we can prove that

$$\left\{\int_{F\setminus B(0,\sqrt{t})} \left[\int_{E\setminus B(0,\sqrt{t})} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} |f(y)| dy\right]^q dx\right\}^{1/q} \lesssim t^{-\frac{d}{2}(\frac{1}{p}-\frac{1}{q})} e^{-\frac{d(E,F)^2}{2ct}} \left(\int_E |f|^p\right)^{1/p}.$$

Hence,

$$E_4 \lesssim t^{-\frac{d}{2}(\frac{1}{p}-\frac{1}{q})} e^{-\frac{d(E,F)^2}{2ct}} \left(\int_E |f|^p\right)^{1/p}.$$

This completes the proof of (12).

Theorem 3.3 ([29, 25]). Assume $d \ge 3$ and $a \ge -\left(\frac{d-2}{2}\right)^2$. Let $p_t(x,y)$ be the kernel associated to the semigroups $e^{-t\mathcal{L}_a}$. Then there exist two positive constants C and c such that for all t > 0and $x, y \in \mathbb{R}^d \setminus \{0\}$,

$$p_t(x,y) \le C \left(1 + \frac{\sqrt{t}}{|x|}\right)^\sigma \left(1 + \frac{\sqrt{t}}{|y|}\right)^\sigma t^{-d/2} e^{-\frac{|x-y|^2}{ct}}$$

The following results gives some estimates of the heat kernels $p_z(x, y)$ for $z \in \mathbb{C}_{\pi/4} := \{z \in \mathbb{C} :$ $|\arg z| < \pi/4\}.$

Proposition 3.4. Let $p_z(x,y)$ be the kernels associated to the semigroups $e^{-z\mathcal{L}_a}$ with $z \in \mathbb{C}_{\pi/4}$:= $\{z \in \mathbb{C} : |\arg z| < \pi/4\}$. Then there exists constants C and c such that

(13)
$$|p_z(x,y)| \le C \left(1 + \frac{\sqrt{|z|}}{|x|}\right)^{\sigma} \left(1 + \frac{\sqrt{|z|}}{|y|}\right)^{\sigma} |z|^{-d/2} e^{-\frac{|x-y|^2}{c|z|}}$$

Proof. We adapt the standard argument in [11] to our present situation. It suffices to claim that

(14)
$$|w(x)p_z(x,y)w(y)| \le \frac{C}{|z|^{d/2}}$$

where $w(x) = \left(1 + \frac{\sqrt{|z|}}{|y|}\right)^{-\sigma}$ Now for $f : \mathbb{R}^d \to \mathbb{R}$, we define the norm

$$|f|_{wL^{\infty}} = \sup_{x \to \infty} |f(x)w(x)|.$$

Hence (14) is equivalent to that

$$\|e^{-z\mathcal{L}_a}\|_{L^1_{w^{-1}}\to wL^\infty} \le \frac{C}{|z|^{d/2}}$$

Assume that z = 2t + is where $t \ge 0$ and $s \in \mathbb{R}$. Then we have $t \sim |z|$ and

$$\|e^{-z\mathcal{L}_{a}}\|_{L^{1}_{w^{-1}}\to wL^{\infty}} \leq \|e^{-t\mathcal{L}_{a}}\|_{L^{2}\to wL^{\infty}}\|e^{-is\mathcal{L}_{a}}\|_{L^{2}\to L^{2}}\|e^{-t\mathcal{L}_{a}}\|_{L^{1}_{w^{-1}}\to L^{2}}.$$

Since \mathcal{L}_a is nonnegative and self-adjoint, $\|e^{-is\mathcal{L}_a}\|_{L^2\to L^2} \leq 1$. We now claim that

$$\|e^{-t\mathcal{L}_a}\|_{L^1_{w^{-1}}\to L^2} \lesssim t^{-d/4}$$
 and $\|e^{-t\mathcal{L}_a}\|_{L^2\to wL^{\infty}} \lesssim t^{-d/4}.$

We now show $\|e^{-t\mathcal{L}_a}\|_{L^1_{w^{-1}}\to L^2} \lesssim t^{-d/4}$. The inequality $\|e^{-t\mathcal{L}_a}\|_{L^2\to wL^{\infty}} \lesssim t^{-d/4}$ can be done in the same manner. Indeed, for $f \in L^1_{w^{-1}}$ we have

$$\begin{split} \|e^{-t\mathcal{L}_{a}}f\|_{L^{2}} &\lesssim \left(\int_{\mathbb{R}^{d}} \left|\int_{\mathbb{R}^{d}} \left(1 + \frac{\sqrt{t}}{|x|}\right)^{\sigma} \left(1 + \frac{\sqrt{t}}{|y|}\right)^{\sigma} t^{-d/2} e^{-\frac{|x-y|^{2}}{ct}} |f(y)| dy \Big|^{2} dx\right)^{1/2} \\ &\lesssim \int_{\mathbb{R}^{d}} \left(\int_{\mathbb{R}} \left|\left(1 + \frac{\sqrt{t}}{|x|}\right)^{\sigma} t^{-d/2} e^{-\frac{|x-y|^{2}}{ct}}\right|^{2} dx\right)^{1/2} \left(1 + \frac{\sqrt{t}}{|y|}\right)^{\sigma} |f(y)| dy \end{split}$$

Arguing similarly to the proof of Theorem 3.1 we get that

$$\left(\int_{\mathbb{R}^d} \left| \left(1 + \frac{\sqrt{t}}{|x|}\right)^{\sigma} t^{-d/2} e^{-\frac{|x-y|^2}{ct}} \right|^2 dx \right)^{1/2} \lesssim t^{-d/4}.$$

Hence,

$$\|e^{-t\mathcal{L}_a}f\|_{L^2} \lesssim t^{-d/4} \int_{\mathbb{R}^d} \left(1 + \frac{\sqrt{t}}{|y|}\right)^{\sigma} |f(y)| dy := t^{-d/4} \|f\|_{L^1_{w^{-1}}}$$

which implies

$$||e^{-t\mathcal{L}_a}||_{L^1_{w^{-1}}\to L^2} \lesssim t^{-d/4}.$$

This completes our proof.

As a direct consequence of Proposition 3.4 and Cauchy formula, we obtain the following result.

Proposition 3.5. Assume $d \ge 3$ and $a \ge -\left(\frac{d-2}{2}\right)^2$. For any $k \in \mathbb{N}$, there exist two positive constants C_k and c_k such that for all t > 0 and $x, y \in \mathbb{R}^d \setminus \{0\}$,

$$|p_{t,k}(x,y)| \le C_k \left(1 + \frac{\sqrt{t}}{|x|}\right)^{\sigma} \left(1 + \frac{\sqrt{t}}{|y|}\right)^{\sigma} t^{-(k+d/2)} e^{-\frac{|x-y|^2}{c_k t}}.$$

where $p_{t,k}(x,y)$ is an associated kernel to $\mathcal{L}_a^k e^{-t\mathcal{L}_a}$.

4. RIESZ TRANSFORMS AND SMOOTHING ESTIMATES

4.1. Weighted Hardy inequalities. The Hardy inequality for Laplacian $-\Delta$ was studied in [34] and then was generalized for Schrödinger operators \mathcal{L}_a in [23]. In this section, we extend to the weighted Hardy inequalities for \mathcal{L}_a .

Theorem 4.1. Suppose $0 < s < d, d - s - 2\sigma > 0$, and $d'_{\sigma} . Then for <math>w \in A_{\frac{p}{d'_{\sigma}}} \cap RH_{\left(\frac{d_{s+\sigma}}{p}\right)'}$ we have

$$|||x|^{-s}f||_{L^p_w(\mathbb{R}^d)} \lesssim ||\mathcal{L}^{s/2}_a f||_{L^p_w(\mathbb{R}^d)}$$

Proof. It suffices to prove that

$$|||x|^{-s}\mathcal{L}_{a}^{-s/2}g||_{L_{w}^{p}(\mathbb{R}^{d})} \lesssim ||g||_{L_{w}^{p}(\mathbb{R}^{d})}$$

for $d'_{\sigma} and <math>w \in A_{\frac{p}{d'_{\sigma}}} \cap RH_{\left(\frac{d_{s+\sigma}}{p}\right)'}$. To do so we shall apply Theorem 2.3.

We define a linear operator

$$T_{\mathcal{L}_a,s}f(x) = |x|^{-s}\mathcal{L}_a^{-s/2}f(x).$$

Fix $p \in (d'_{\sigma}, d_{s+\sigma})$ and $w \in A_{\frac{p}{d'_{\sigma}}} \cap RH_{\left(\frac{d_{s+\sigma}}{p}\right)'}$. Then we can find $d'_{\sigma} < p_1 < p < p_2 < d_{s+\sigma}$ so that $w \in A_{\frac{p}{p_1}} \cap RH_{\left(\frac{p_2}{\sigma}\right)'}$.

We now fix a ball $B \subset \mathbb{R}^d$ and $m > d + d/p_2$. For any function f supported in B we claim that

(15)
$$\left(\int_{S_j(B)} |T_{\mathcal{L}_a,s}(I-e^{-r_B^2\mathcal{L}_a})^m f(x)|^{p_2} dx\right)^{1/p_2} \lesssim 2^{-2mj} \left(\int_B |f(x)|^{p_2} dx\right)^{1/p_2}.$$

Indeed, using the formula

$$\mathcal{L}_a^{-s/2} = \frac{1}{\Gamma(s/2)} \int_0^\infty t^{s/2} e^{-t\mathcal{L}_a} \frac{dt}{t}$$

to obtain that

(16)
$$T_{\mathcal{L}_a,s}(I - e^{-r_B^2 \mathcal{L}_a})^m f(x) = \frac{1}{\Gamma(s/2)} \int_0^\infty t^{s/2} |x|^{-s} e^{-t\mathcal{L}_a} (I - e^{-r_B^2 \mathcal{L}_a})^m f(x) \frac{dt}{t}.$$

This along with Minkowski's inequality implies that

$$\begin{split} \|T_{\mathcal{L}_{a},s}(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}f\|_{L^{p_{2}}(S_{j}(B))} &\lesssim \frac{1}{\Gamma(s/2)} \int_{0}^{\infty} t^{s/2} \||x|^{-s} e^{-t\mathcal{L}_{a}}(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}f(x)\|_{L^{p_{2}}(S_{j}(B))} \frac{dt}{t} \\ &\lesssim \frac{1}{\Gamma(s/2)} \int_{0}^{r_{B}^{2}} t^{s/2} \||x|^{-s} e^{-t\mathcal{L}_{a}}(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}f(x)\|_{L^{p_{2}}(S_{j}(B))} \frac{dt}{t} \\ &+ \frac{1}{\Gamma(s/2)} \int_{r_{B}^{2}}^{\infty} t^{s/2} \||x|^{-s} e^{-t\mathcal{L}_{a}}(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}f(x)\|_{L^{p_{2}}(S_{j}(B))} \frac{dt}{t} \\ &\lesssim E_{1} + E_{2}. \end{split}$$

We first take care of E_1 . Observe that

(17)
$$E_1 \le \sum_{k=0}^m C_k^m \int_0^{r_B^2} t^{s/2} || |x|^{-s} e^{-(t+kr_B^2)\mathcal{L}_a} f(x) ||_{L^{p_2}(S_j(B))} \frac{dt}{t}.$$

Note that the associated kernel of the linear operator $f(x) \mapsto |x|^{-s} e^{-(t+kr_B^2)\mathcal{L}_a} f(x)$ is given by $|x|^{-s} p_{t+kr_B^2}(x,y)$ and by Theorem 3.3 it is dominated by

$$\begin{aligned} |x|^{-s} \Big(1 + \frac{\sqrt{t + kr_B^2}}{|x|}\Big)^{\sigma} \Big(1 + \frac{\sqrt{t + kr_B^2}}{|y|}\Big)^{\sigma} (t + kr_B^2)^{-d/2} e^{-\frac{|x - y|^2}{c(t + kr_B^2)}} \\ &\leq (t + kr_B^2)^{-s/2} \Big(1 + \frac{\sqrt{t + kr_B^2}}{|x|}\Big)^{\sigma + s} \Big(1 + \frac{\sqrt{t + kr_B^2}}{|y|}\Big)^{\sigma} (t + kr_B^2)^{-d/2} e^{-\frac{|x - y|^2}{c(t + kr_B^2)}}. \end{aligned}$$

Therefore, applying Theorem 3.1 we get that

$$\| |x|^{-s} e^{-(t+kr_B^2)\mathcal{L}_a} f(x) \|_{L^{p_2}(S_j(B))} \lesssim (t+kr_B^2)^{-s/2} e^{-\frac{4^j r_B^2}{c(t+kr_B^2)}} \left(\int_B |f|^{p_2} \right)^{1/p_2}.$$

Inserting this into (17) to obtain that

$$\begin{split} E_1 &\lesssim \sum_{k=0}^m \int_0^{r_B^2} t^{s/2} (t+kr_B^2)^{-s/2} e^{-\frac{4^j r_B^2}{c(t+kr_B^2)}} \frac{dt}{t} \times \left(\int_B |f|^{p_2}\right)^{1/p_2} \\ &\lesssim \sum_{k=0}^m \int_0^{r_B^2} t^{s/2} (t+kr_B^2)^{-s/2} \left(\frac{t+kr_B^2}{4^j r_B^2}\right)^m \frac{dt}{t} \times \left(\int_B |f|^{p_2}\right)^{1/p_2} \\ &\lesssim 2^{-2mj} \left(\int_B |f|^{p_2}\right)^{1/p_2}. \end{split}$$

To estimate E_2 , we note that

$$(I - e^{-r_B^2 \mathcal{L}_a})^m = \int_0^{r_B^2} \dots \int_0^{r_B^2} \mathcal{L}_a^m e^{-(s_1 + \dots + s_m)\mathcal{L}_a} d\vec{s},$$

where $d\vec{s} = ds_1 \dots ds_m$.

Hence, the associated kernel to the linear operator $f(x) \mapsto |x|^{-s} e^{-t\mathcal{L}_a} (I - e^{-r_B^2 \mathcal{L}_a})^m f(x)$ is given by

$$\int_0^{r_B^2} \dots \int_0^{r_B^2} |x|^{-s} p_{t+s_1+\dots+s_m,m}(x,y) d\vec{s},$$

and hence by using Theorem 3.5 we can dominate it by

$$\begin{split} |x|^{-s}t^{-m}\Big(1 + \frac{\sqrt{t+s_1+\dots+s_m}}{|x|}\Big)^{\sigma}\Big(1 + \frac{\sqrt{t+s_1+\dots+s_m}}{|y|}\Big)^{\sigma}(t+s_1+\dots+s_m)^{-d/2}e^{-\frac{|x-y|^2}{c(t+s_1+\dots+s_m)}}\\ &\lesssim |x|^{-s}t^{-m}\Big(1 + \frac{\sqrt{t}}{|x|}\Big)^{\sigma}\Big(1 + \frac{\sqrt{t}}{|y|}\Big)^{\sigma}t^{-d/2}e^{-\frac{|x-y|^2}{c(t)}}\\ &\lesssim t^{-(s/2+m)}\Big(1 + \frac{\sqrt{t}}{|x|}\Big)^{\sigma+s}\Big(1 + \frac{\sqrt{t}}{|y|}\Big)^{\sigma}t^{-d/2}e^{-\frac{|x-y|^2}{c(t)}}, \end{split}$$

where in the first inequality we used the fact that $t + s_1 + \cdots + s_m \sim t$ for $t \geq r_B^2$ and $s_i \in (0, r_B^2], i = 1, \ldots, m$.

This in combination with Theorem 3.1 implies that

$$\| |x|^{-s} e^{-t\mathcal{L}_a} (I - e^{-r_B^2 \mathcal{L}_a})^m f(x) \|_{L^{p_2}(S_j(B))} \lesssim t^{-(s/2+m)} e^{-\frac{4^j r_B^2}{ct}} \Big(\int_B |f|^{p_2} \Big)^{1/p_2}$$

Inserting this into the expression of E_2 to get that

$$\begin{split} E_2 &\lesssim \int_{r_B^2}^{\infty} t^{s/2} t^{-(s/2+m)} e^{-\frac{4^j r_B^2}{ct}} \frac{dt}{t} \times \left(\int_B |f|^{p_2} \right)^{1/p_2} \\ &\lesssim \int_{r_B^2}^{\infty} t^{s/2} t^{-(s/2+m)} \left(\frac{t}{4^j r_B^2} \right)^m \frac{dt}{t} \times \left(\int_B |f|^{p_2} \right)^{1/p_2} \\ &\lesssim 2^{-2mj} \left(\int_B |f|^{p_2} \right)^{1/p_2}. \end{split}$$

From the estimates of E_1 and E_2 we conclude (15).

With estimate (15) in hand, we can now complete the proof of the theorem. First note that it was proved in [23] that $T_{\mathcal{L}_{a},s}$ is bounded on L^{p} for $p \in (d'_{\sigma}, d_{s+\sigma})$. For each ball B, we now set

$$\mathcal{A}_{r_B} = I - (I - e^{-r_B^2 \mathcal{L}_a})^m.$$

Then from (15), we conclude that '

(18)
$$\left(\oint_{S_j(B)} |T_{\mathcal{L}_a,s}(I - \mathcal{A}_{r_B})f(x)|^{p_2} dx \right)^{1/p_2} \lesssim 2^{-j(2m-d/p_2)} \left(\oint_B |f(x)|^{p_2} dx \right)^{1/p_2}$$

On the other hand from Theorem 3.1 and Proposition 3.5 we imply that '

(19)
$$\left(\int_{S_j(B)} |\mathcal{A}_{r_B} f(x)|^{p_2} dx\right)^{1/p_2} \lesssim e^{-c2^{2j}} |B|^{\frac{1}{p_1} - \frac{1}{p_2}} \left(\int_B |f(x)|^{p_1} dx\right)^{1/p_1}$$

which implies that '

(20)
$$\left(\int_{S_j(B)} |\mathcal{A}_{r_B} f(x)|^{p_2} dx\right)^{1/p_2} \lesssim 2^{-j(2m-d/p_2)} \left(\int_B |f(x)|^{p_1} dx\right)^{1/p_1}.$$

From (18), (20) and Theorem 2.3 we obtained the desired result.

4.2. Weighted estimates for square functions. Let $\alpha \in (0,1)$ we consider the following square function

$$S_{\mathcal{L}_a,\alpha}f(x) = \left(\int_0^\infty |(t\mathcal{L}_a)^{1-\alpha}e^{-t\mathcal{L}_a}f|^2\frac{dt}{t}\right)^{1/2}.$$

Note that by functional calculus theory in [28], the square function $S_{\mathcal{L}_a,\alpha}$ is bounded on L^2 . In the following theorem, we prove the weighted L^p estimates for $S_{\mathcal{L}_a,\alpha}$.

Theorem 4.2. Suppose that $d \ge 3$, $a \ge -\left(\frac{d-2}{2}\right)^2$ and $\alpha \in (0,1)$. Then for all $d'_{\sigma} and <math>w \in A_{\frac{p}{d'_{\sigma}}} \cap RH_{\left(\frac{d_{\sigma}}{p}\right)'}$ we have

$$\|S_{\mathcal{L}_a,\alpha}f\|_{L^p_w} \sim \|f\|_{L^p_w}$$

As a consequence, for 0 < s < 2, $d'_{\sigma} and <math>w \in A_{\frac{p}{d'_{\sigma}}} \cap RH_{(\frac{d_{\sigma}}{p})'}$ we have

$$\left\| \left(\int_0^\infty t^{-s} |t\mathcal{L}_a e^{-t\mathcal{L}_a} f|^2 \frac{dt}{t} \right)^{1/2} \right\|_{L^p_w} \sim \|\mathcal{L}_a^{s/2} f\|_{L^p_w}.$$

Proof. We shall apply again Theorem 2.3. To do so we assume for now that $S_{\mathcal{L}_a,\alpha}$ is bounded on L^p for $p \in [2, d_{\sigma})$. This assumption will be justified later.

Fix $d'_{\sigma} and <math>w \in A_{\frac{p}{d'_{\sigma}}} \cap RH_{(\frac{d_{\sigma}}{p})'}$. Then we can pick $d'_{\sigma} < p_1 < p < p \lor 2 < p_2 < d_{\sigma}$ so that $w \in A_{\frac{p}{p_1}} \cap RH_{(\frac{p_2}{p})'}$. Fix $m > 1 + d + d/p_2$, we will claim that

(21)
$$\left(\int_{S_j(B)} |S_{\mathcal{L}_a,\alpha}(I - e^{-r_B^2 \mathcal{L}_a})^m f|^{p_2} dx \right)^{1/p_2} \lesssim 2^{-2(m-1)j} \left(\int_B |f|^{p_2} dx \right)^{1/p_2}, \qquad j \ge 2$$

for all balls B and all $f \in C^{\infty}$ supported in B.

Indeed, by Minkowski's inequality we have

$$\begin{split} \left(\int_{S_{j}(B)} |S_{\mathcal{L}_{a},\alpha}(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}f|^{p_{2}}dx\right)^{1/p_{2}} &\leq \left(\int_{0}^{r_{B}^{2}} \left\|(t\mathcal{L}_{a})^{1-\alpha}e^{-t\mathcal{L}_{a}}(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}f\right\|_{L^{p_{2}}(S_{j}(B))}^{2}\frac{dt}{t}\right)^{1/2} \\ &+ \left(\int_{r_{B}^{2}}^{\infty} \left\|(t\mathcal{L}_{a})^{1-\alpha}e^{-t\mathcal{L}_{a}}(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}f\right\|_{L^{p_{2}}(S_{j}(B))}^{2}\frac{dt}{t}\right)^{1/2} \\ &:= I_{1} + I_{2}. \end{split}$$

We now take care of I_1 first. Note that

$$\mathcal{L}_{a}^{-\alpha} = \frac{1}{\Gamma(\alpha)} \int_{0}^{\infty} u^{\alpha} e^{-u\mathcal{L}_{a}} \frac{du}{u}$$

and

(22)
$$(I - e^{-r_B^2 \mathcal{L}_a})^m = \int_{[0, r_B^2]^m} \mathcal{L}_a^m e^{-(s_1 + \dots + s_m)\mathcal{L}_a} d\vec{s},$$

where $d\vec{s} = ds_1 \dots ds_m$.

Hence,

$$I_{1} \lesssim \left(\int_{0}^{r_{B}^{2}} \left[\int_{0}^{r_{B}^{2}} \left(\frac{u}{t} \right)^{\alpha} \left\| t\mathcal{L}_{a}e^{-(t+u)\mathcal{L}_{a}} (I - e^{-r_{B}^{2}\mathcal{L}_{a}})^{m} f \right\|_{L^{p_{2}}(S_{j}(B))} \frac{du}{u} \right]^{2} \frac{dt}{t} \right)^{1/2} \\ + \left(\int_{0}^{r_{B}^{2}} \left[\int_{r_{B}^{2}}^{\infty} \left(\frac{u}{t} \right)^{\alpha} \left\| t\mathcal{L}_{a}e^{-(t+u)\mathcal{L}_{a}} (I - e^{-r_{B}^{2}\mathcal{L}_{a}})^{m} f \right\|_{L^{p_{2}}(S_{j}(B))} \frac{du}{u} \right]^{2} \frac{dt}{t} \right)^{1/2} \\ := I_{11} + I_{12}.$$

We now have

$$\begin{split} I_{11} &\lesssim \sum_{k=0}^{m} \left(\int_{0}^{r_{B}^{2}} \left[\int_{0}^{r_{B}^{2}} \left(\frac{u}{t} \right)^{\alpha} \frac{t}{t+u+kr_{B}^{2}} \left\| (t+u+kr_{B}^{2})\mathcal{L}_{a}e^{-(t+u+kr_{B}^{2})\mathcal{L}_{a}}f \right\|_{L^{p_{2}}(S_{j}(B))} \frac{du}{u} \right]^{2} \frac{dt}{t} \right)^{1/2} \\ &\lesssim \sum_{k=0}^{m} \left(\int_{0}^{r_{B}^{2}} \left[\int_{0}^{r_{B}^{2}} \left(\frac{u}{t} \right)^{\alpha} \frac{t}{t+u+kr_{B}^{2}} \exp\left(-\frac{4^{j}r_{B}^{2}}{c(t+u+kr_{B}^{2})} \right) \|f\|_{L^{p_{2}}(B)} \frac{du}{u} \right]^{2} \frac{dt}{t} \right)^{1/2} \\ &\lesssim \sum_{k=0}^{m} \left(\int_{0}^{r_{B}^{2}} \left[\int_{0}^{r_{B}^{2}} \left(\frac{u}{t} \right)^{\alpha} \frac{t}{4^{j}r_{B}^{2}} \left(\frac{t+u+kr_{B}^{2}}{4^{j}r_{B}^{2}} \right)^{m} \|f\|_{L^{p_{2}}(B)} \frac{du}{u} \right]^{2} \frac{dt}{t} \right)^{1/2} \\ &\lesssim 2^{-2mj} \|f\|_{L^{p_{2}}(B)} \,. \end{split}$$

Now we use (22) to obtain that

$$I_{12} \lesssim \left(\int_{0}^{r_{B}^{2}} \left[\int_{[0,r_{B}^{2}]^{m}} \int_{r_{B}^{2}}^{\infty} \left(\frac{u}{t}\right)^{\alpha} \left\| t\mathcal{L}_{a}^{m+1}e^{-(t+u+s_{1}+\ldots+s_{m})\mathcal{L}_{a}}f \right\|_{L^{p_{2}}(S_{j}(B))} \frac{du}{u}d\vec{s} \right]^{2} \frac{dt}{t} \right)^{1/2}$$

which along with Theorem 3.1 and the fact that $u \sim t + u + s_1 + \ldots + s_m$ implies that

$$I_{12} \lesssim \left(\int_{0}^{r_B^2} \left[\int_{[0,r_B^2]^m} \int_{r_B^2}^{\infty} \left(\frac{u}{t}\right)^{\alpha} \frac{t}{u^{m+1}} e^{-\frac{4^j r_B^2}{cu}} \|f\|_{L^{p_2}(B)} \frac{du}{u} d\vec{s} \right]^2 \frac{dt}{t} \right)^{1/2} \\ \lesssim 2^{-2mj} \|f\|_{L^{p_2}(B)}.$$

As a consequence,

$$I_1 \lesssim 2^{-2mj} \|f\|_{L^p(B)}.$$

Similarly, we split I_2 as follows

$$I_{2} \lesssim \left(\int_{r_{B}^{2}}^{\infty} \left[\int_{0}^{r_{B}^{2}} \left(\frac{u}{t} \right)^{\alpha} \left\| t \mathcal{L}_{a} e^{-(t+u)\mathcal{L}_{a}} (I - e^{-r_{B}^{2}\mathcal{L}_{a}})^{m} f \right\|_{L^{p_{2}}(S_{j}(B))} \frac{du}{u} \right]^{2} \frac{dt}{t} \right)^{1/2} \\ + \left(\int_{r_{B}^{2}}^{\infty} \left[\int_{r_{B}^{2}}^{\infty} \left(\frac{u}{t} \right)^{\alpha} \left\| t \mathcal{L}_{a} e^{-(t+u)\mathcal{L}_{a}} (I - e^{-r_{B}^{2}\mathcal{L}_{a}})^{m} f \right\|_{L^{p_{2}}(S_{j}(B))} \frac{du}{u} \right]^{2} \frac{dt}{t} \right)^{1/2} \\ := I_{21} + I_{22}.$$

The argument used to estimate I_{12} can be applied again to show that

$$I_{21} \lesssim 2^{-2mj} \|f\|_{L^p(B)}$$

On the other hand, using this argument, we also dominate I_{22} as follows

$$\begin{split} I_{22} &\lesssim \left(\int_{r_B^2}^{\infty} \left[\int_{[0,r_B^2]^m} \int_{r_B^2}^{\infty} \left(\frac{u}{t} \right)^{\alpha} \frac{t}{(u+t)^{m+1}} e^{-\frac{4^j r_B^2}{c(u+t)}} \, \|f\|_{L^{p_2}(B)} \frac{du}{u} d\vec{s} \right]^2 \frac{dt}{t} \right)^{1/2} \\ &\lesssim \left(\int_{r_B^2}^{\infty} \left[\int_{[0,r_B^2]^m} \int_{r_B^2}^{\infty} \left(\frac{u}{t} \right)^{\alpha} \frac{1}{u(u+t)^{m-1}} e^{-\frac{4^j r_B^2}{c(u+t)}} \, \|f\|_{L^{p_2}(B)} \frac{du}{u} d\vec{s} \right]^2 \frac{dt}{t} \right)^{1/2} \\ &\lesssim 2^{-2(m-1)j} \, \|f\|_{L^{p_2}(B)} \, . \end{split}$$

Therefore,

$$I_{22} \lesssim 2^{-2(m-1)j} \|f\|_{L^{p_2}(B)}.$$

Hence, this completes the proof of (21). At this stage, arguing similarly to Theorem 4.1, we obtain that

$$\|S_{\mathcal{L}_a,\alpha}f\|_{L^p_w} \lesssim \|f\|_{L^p_w}.$$

To prove the reverse inequality, by functional calculus theory for $g \in L_v^{p'}$ with $v = w^{1-p'}$ we have

$$\int_{\mathbb{R}^d} f(x)g(x)dx = c(\alpha) \int_{\mathbb{R}^d} \int_0^\infty (t\mathcal{L}_a)^{2(1-\alpha)} e^{-2t\mathcal{L}_a} f(x)g(x)\frac{dt}{t}dx,$$

where $c(\alpha) = \int_0^\infty t^{2(1-\alpha)} e^{-2t} \frac{dt}{t}$. (Actually, this identity holds true in L^2 first. However, due to the weighted L^p boundeness of $S_{\mathcal{L}_a,\alpha}$ we can extend the convergence to L_w^p .)

By Hölder's inequality, we can write

$$\int_{\mathbb{R}^d} f(x)g(x)dx = c(\alpha) \int_{\mathbb{R}^d} \int_0^\infty (t\mathcal{L}_a)^{1-\alpha} e^{-t\mathcal{L}_a} f(x)(t\mathcal{L}_a)^{1-\alpha} e^{-t\mathcal{L}_a} g(x) \frac{dt}{t} dx$$
$$\lesssim \int_{\mathbb{R}^d} S_{\mathcal{L}_a,\alpha} f(x) S_{\mathcal{L}_a,\alpha} g(x) dx$$
$$\lesssim \|S_{\mathcal{L}_a,\alpha} f\|_{L^p_w} \|S_{\mathcal{L}_a,\alpha} g\|_{L^{p'}_v}.$$

Note that from (vii) Lemma 2.1 we obtain $v \in A_{p'/d_{\sigma} \cap RH_{(d'_{\sigma}/p')'}}$. Hence, from the weighted L^p estimates of $S_{\mathcal{L}_a,\alpha}$ we have proved we get that $\|S_{\mathcal{L}_a,\alpha}g\|_{L_p^{p'}} \lesssim \|g\|_{L_p^{p'}}$, we obtain

$$\int_{\mathbb{R}^d} f(x)g(x)dx \lesssim \|S_{\mathcal{L}_a,\alpha}f\|_{L^p_w} \|g\|_{L^{p'}_v}.$$

As a consequence,

$$\|f\|_{L^p_w} \lesssim \|S_{\mathcal{L}_a,\alpha}f\|_{L^p_w}.$$

To complete the proof, we need to prove the original assertion that $S_{\mathcal{L}_a,\alpha}$ is bounded on L^r for all $r \in (2, d_{\sigma})$. According to Theorem 2.4, for any $q_0 \in (2, d_{\sigma})$ it suffices to prove that

(23)
$$\left(\int_{B} \left|S_{\mathcal{L}_{a},\alpha}(I-\mathcal{A}_{r_{B}})f\right|^{2} dx\right)^{1/2} \leq C\mathcal{M}_{2}(f)(x),$$

and

(24)
$$\left(\int_{B} \left|S_{\mathcal{L}_{a},\alpha}\mathcal{A}_{r_{B}}f\right|^{q_{0}}dx\right)^{1/q_{0}} \leq C\mathcal{M}_{2}(\left|S_{\mathcal{L}_{a},\alpha}f\right|)(x)$$

all balls B with radius r_B , all $f \in C_c^{\infty}(\mathbb{R}^n)$ and all $x \in B$ with $\mathcal{A}_{r_B} = I - (I - e^{-r_B^2 \mathcal{L}_a})^m$, $m > 1 + d + d/p_2$.

To prove (23), we write

$$\left(\int_{B} \left|S_{\mathcal{L}_{a},\alpha}(I-\mathcal{A}_{r_{B}})f\right|^{2} dx\right)^{1/2} \leq \sum_{j=0}^{\infty} \left(\int_{B} \left|S_{\mathcal{L}_{a},\alpha}(I-\mathcal{A}_{r_{B}})f_{j}\right|^{2} dx\right)^{1/2}$$
$$:= \sum_{j=0}^{\infty} I_{j},$$

where $f_j = f \chi_{S_j(B)}$.

For j = 0, 1, using the L^2 -boundedness of $S_{\mathcal{L}_a,\alpha}$ and \mathcal{A}_{r_B} we have

$$I_j \lesssim \mathcal{M}_2(f)(x)$$

For $j \ge 2$, the argument in the proof of (21) shows that

$$I_j \lesssim 2^{-j(2m-d/2)} \left(\int_{S_j(B)} |f|^2 \right)^{1/2}$$

Therefore,

$$\sum_{j=0}^{\infty} I_j \lesssim \mathcal{M}_2(f)(x)$$

which proves (23).

It remains to prove (24). Indeed, we have

$$\begin{split} \left(\int_{B} |S_{\mathcal{L}_{a},\alpha}[I-(I-e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}]f(x)|^{q_{0}}dx\right)^{1/q_{0}} \\ &\lesssim \sum_{k=1}^{m} \left(\int_{B} |S_{\mathcal{L}_{a},\alpha}e^{-kr_{B}^{2}\mathcal{L}_{a}}f(x)|^{q_{0}}dx\right)^{1/q_{0}} \\ &\lesssim \sup_{1\leq k\leq m} \left[\int_{B} \left(\int_{0}^{\infty} |e^{-kr_{B}^{2}\mathcal{L}_{a}}(t\mathcal{L}_{a})^{1-\alpha}e^{-t\mathcal{L}_{a}}f(x)|^{2}\frac{dt}{t}\right)^{q_{0}/2}dx\right]^{1/q_{0}} \\ &\lesssim \sum_{j\geq 0} \sup_{1\leq k\leq m} \left[\int_{B} \left(\int_{0}^{\infty} \left|e^{-kr_{B}^{2}\mathcal{L}_{a}}[(t\mathcal{L}_{a})^{1-\alpha}e^{-t\mathcal{L}_{a}}f\chi_{S_{j}(B)}](x)\right|^{2}\frac{dt}{t}\right)^{q_{0}/2}dx\right]^{1/q_{0}} \end{split}$$

which along with Minkowski's inequality, Theorem 3.5 and Theorem 3.1 gives

$$\left(\int_{B} |S_{\mathcal{L}_{a},\alpha}[I - (I - e^{-r_{B}^{2}\mathcal{L}_{a}})^{m}]f(x)|^{q_{0}}dx \right)^{1/q_{0}}$$

$$\lesssim \sum_{j\geq 0} \sup_{1\leq k\leq m} \left(\int_{0}^{\infty} \left\| e^{-kr_{B}^{2}\mathcal{L}_{a}}[(t\mathcal{L}_{a})^{1-\alpha}e^{-t\mathcal{L}_{a}}f\chi_{S_{j}(B)}] \right\|_{L^{q_{0}}(B)}^{2} \frac{dt}{t} \right)^{1/2}$$

$$\lesssim \sum_{j\geq 0} e^{-c4^{j}}|B|^{-(\frac{1}{2}-\frac{1}{q_{0}})} \left(\int_{0}^{\infty} \left\| (t\mathcal{L}_{a})^{1-\alpha}e^{-t\mathcal{L}_{a}}f \right\|_{L^{2}(S_{j}(B))}^{2} \frac{dt}{t} \right)^{1/2}$$

$$\lesssim \sum_{j\geq 0} e^{-c4^{j}}|B|^{-(\frac{1}{2}-\frac{1}{q_{0}})} \left(\int_{2^{j}B} |S_{\mathcal{L}_{a},\alpha}f(x)|^{2}dx \right)^{1/2}.$$

This implies (24). Hence the proof is complete.

The following result regarding weighted estimates for the difference of square functions will play an essential role in the proofs of the main results.

Theorem 4.3. We have the following estimate

$$\left\| \left(\int_0^\infty t^{-s} \left| \left(t\mathcal{L}_a e^{-t\mathcal{L}_a} + t\Delta e^{t\Delta} \right) f \right|^2 \frac{dt}{t} \right)^{1/2} \right\|_{L^p_w} \lesssim \left\| \frac{f}{|x|^s} \right\|_{L^p_w}$$

 $provided\ that$

(a) $a \ge 0, 1 and <math>w \in A_p$; or (b) $-\left(\frac{d-2}{2}\right)^2 \le a < 0, 1 \lor \frac{d}{d+s-\sigma} < p < d_\sigma$ and $w \in A_{\frac{p}{1\lor \frac{d}{d+s-\sigma}}} \cap RH_{(d_\sigma/p)'}$.

Before proceeding with the proof of the theorem, we need the following technical results on kernel estimates.

Let $D_t(x,y)$ be a kernel of $t\mathcal{L}_a e^{-t\mathcal{L}_a} + t\Delta e^{t\Delta}$. We have the following estimates:

Proposition 4.4. (a) If $a \ge 0$ then

(25)
$$|D_t(x,y)| \lesssim t^{-d/2} \left(1 + \frac{|x| + |y|}{\sqrt{t}}\right)^{-2} e^{-\frac{|x-y|^2}{ct}}$$

for all $x, y \in \mathbb{R}^d$ and t > 0. (b) If $-\left(\frac{d-2}{2}\right)^2 \le a < 0$ then

(26)
$$|D_t(x,y)| \lesssim t^{-d/2} \left(1 + \frac{|x| + |y|}{\sqrt{t}}\right)^{-2} e^{-\frac{|x-y|^2}{ct}}$$

for all t > 0 and $|x|, |y| \ge \sqrt{t}/2$.

We remark that in [23] the authors gave upper bounds for the kernels of $e^{-t\mathcal{L}_a} - e^{t\Delta}$. However, this estimate is not sufficient for us.

Proof. We first give the proof for the case $a \ge 0$. Note that in this case since both kernels of $t\mathcal{L}_a e^{-t\mathcal{L}_a}$ and $t\Delta e^{t\Delta}$ satisfy Gaussian upper bounds, there exists C, c > 0 so that

$$|D_t(x,y)| \le Ct^{-d/2}e^{-\frac{|x-y|^2}{ct}}$$

for all $x, y \in \mathbb{R}^d$ and t > 0.

Hence, it suffices to prove (25) for $x \sim y$ and $|x|, |y| \geq \sqrt{t}/2$. From Duhamel's formula, we obtain that

(27)
$$D_{t}(x,y) = at \int_{\mathbb{R}^{d}} \tilde{p}_{t/2}(x,z)|z|^{-2}p_{t/2}(z,y)dz + at \int_{0}^{t/2} \int_{\mathbb{R}^{d}} \tilde{p}_{t-s,1}(x,z)|z|^{-2}p_{s}(z,y)dz \frac{ds}{t-s} + at \int_{t/2}^{t} \int_{\mathbb{R}^{d}} \tilde{p}_{t-s}(x,z)|z|^{-2}p_{s,1}(z,y)dz \frac{ds}{s} = I_{1} + I_{2} + I_{3},$$

where $\tilde{p}_{t,k}(x,y)$ denotes the kernel of $(-1)^k (t\Delta)^k e^{t\Delta}$.

Using the fact that $0 \le p_t(x,y) \le \tilde{p}_t(x,y)$ to get that

$$I_{1} \lesssim t \int_{\mathbb{R}^{d}} \frac{1}{t^{d/2}} e^{-\frac{|x-z|^{2}}{8t}} |z|^{-2} e^{-\frac{|z-y|^{2}}{8t}} dz$$
$$\lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^{2}}{16t}} t \int_{\mathbb{R}^{d}} \frac{1}{t^{d/2}} e^{-\frac{|x-z|^{2}}{16t}} |z|^{-2} dz$$

which along with the fact that

(28)
$$\int_{\mathbb{R}^d} \frac{1}{t^{d/2}} e^{-\frac{|x-z|^2}{ct}} |z|^{-2} dz \lesssim \frac{1}{|x|^2}$$

implies that

$$I_1 \lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{16t}} \frac{t}{|x|^2}.$$

Similarly, by using the Gaussian upper bounds of $\tilde{p}_{t-s,1}(x,z)$ and $p_s(z,y)$ and the fact that

$$\frac{|x-z|^2}{t-s} + \frac{|z-y|^2}{s} \ge \frac{|x-y|^2}{2t} \text{ for all } s \in (0,t)$$

we also obtain that

$$\begin{split} I_2 &\lesssim e^{-\frac{|x-y|^2}{ct}} t \int_0^{t/2} \int_{\mathbb{R}^d} \frac{1}{(t-s)^{d/2}} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^2}{c's}} |z|^{-2} dz \frac{ds}{t-s} \\ &\lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{ct}} t \int_0^{t/2} \int_{\mathbb{R}^d} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^2}{c's}} |z|^{-2} dz \frac{ds}{t} \\ &\lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{16t}} \frac{t}{|y|^2} \end{split}$$

where in the last inequality we used (28).

Similarly, by a change of variable and arguing as in I_2 ,

$$I_3 \lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{16t}} \frac{t}{|x|^2}.$$

This completes the proof for the case $a \ge 0$.

We now consider the case when $-\left(\frac{d-2}{2}\right)^2 \le a < 0$. In this situation $\sigma > 0$ and thus it is easy to observe that

$$|D_t(x,y)| \lesssim t^{-d/2} e^{-\frac{|x-y|^2}{ct}}$$

whenever $|x|, |y| \ge \sqrt{t}/2$. Hence, it suffices to prove (26) for $|x|, |y| \ge \sqrt{t}$ and $|x| \sim |y|$. By expressing $D_t(x, y)$ as in (27), we will need to estimate I_1, I_2, I_3 for $-\left(\frac{d-2}{2}\right)^2 < a < 0$ and $|x|, |y| \ge \sqrt{t}$ and $|x| \sim |y|$.

Arguing similarly to the case $a \ge 0$, we have

$$I_1 \lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{16t}} \frac{t}{|x|^2}.$$

For the second term I_2 , from the kernel bounds estimates of $\tilde{p}_{t,1}(x,y)$ and $p_t(x,y)$, and arguing similarly to the case 1, we obtain

$$\begin{split} I_{2} &\lesssim e^{-\frac{|x-y|^{2}}{ct}} t \int_{0}^{t/2} \int_{\mathbb{R}^{d}} \frac{1}{(t-s)^{d/2}} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^{2}}{c's}} |z|^{-2} \left(1 + \frac{\sqrt{s}}{|z|}\right)^{\sigma} dz ds \\ &\lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^{2}}{ct}} t \int_{0}^{t/2} \int_{\mathbb{R}^{d}} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^{2}}{c's}} |z|^{-2} \left(1 + \frac{\sqrt{s}}{|z|}\right)^{\sigma} dz ds \\ &\lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^{2}}{ct}} t \int_{0}^{t/2} \int_{|z| \ge \sqrt{s}} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^{2}}{c's}} |z|^{-2} \left(1 + \frac{\sqrt{s}}{|z|}\right)^{\sigma} dz ds \\ &+ \frac{1}{t^{d/2}} e^{-\frac{|x-y|^{2}}{ct}} t \int_{0}^{t/2} \int_{|z| < \sqrt{s}} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^{2}}{c's}} |z|^{-2} \left(1 + \frac{\sqrt{s}}{|z|}\right)^{\sigma} dz ds \end{split}$$

Similarly to the case $a \ge 0$, we have

$$\begin{split} \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{ct}} t \int_0^{t/2} \int_{|z| \ge \sqrt{s}} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^2}{c's}} |z|^{-2} \left(1 + \frac{\sqrt{s}}{|z|}\right)^\sigma dz ds \\ &\lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{ct}} t \int_0^{t/2} \int_{|z| \ge \sqrt{s}} \frac{1}{s^{d/2}} e^{-\frac{|z-y|^2}{c's}} |z|^{-2} dz ds \\ &\lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{16t}} \frac{t}{|x|^2}. \end{split}$$

From the fact that $\sigma + 2 < d$ we have

$$\int_{\mathbb{R}^d} \frac{1}{t^{d/2}} e^{-\frac{|x-z|^2}{ct}} |z|^{-2} \left(1 + \frac{\sqrt{s}}{|z|}\right)^{\sigma} dz \lesssim \frac{1}{|x|^2}.$$

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This implies that

$$\frac{1}{t^{d/2}}e^{-\frac{|x-y|^2}{ct}}t\int_0^{t/2}\int_{|z|<\sqrt{s}}\frac{1}{s^{d/2}}e^{-\frac{|z-y|^2}{c's}}|z|^{-2}\left(1+\frac{\sqrt{s}}{|z|}\right)^{\sigma}dzds \lesssim \frac{1}{t^{d/2}}e^{-\frac{|x-y|^2}{16t}}\frac{t}{|x|^2}$$

Likewise, we get that

$$I_3 \lesssim \frac{1}{t^{d/2}} e^{-\frac{|x-y|^2}{16t}} \frac{t}{|x|^2}.$$

This completes our proof.

Proof of Theorem 4.3: We consider two cases. Case 1: $a \ge 0$

Fix $1 and <math>w \in A_p$. Observe that by Proposition 4.4

$$\begin{split} \left(\int_{0}^{\infty} t^{-s} \left| \left(t\mathcal{L}_{a}e^{-t\mathcal{L}_{a}} + t\Delta e^{t\Delta} \right) f(x) \right|^{2} \frac{dt}{t} \right)^{1/2} \\ &\leq \left(\sum_{j \in \mathbb{Z}} \int_{2^{2j}}^{2^{2(j+1)}} t^{-s} \left| \left(t\mathcal{L}_{a}e^{-t\mathcal{L}_{a}} + t\Delta e^{t\Delta} \right) f(x) \right|^{2} \frac{dt}{t} \right)^{1/2} \\ &\leq \left[\sum_{j \in \mathbb{Z}} \int_{2^{2j}}^{2^{2(j+1)}} t^{-s} \left(\int_{\mathbb{R}^{d}} |D_{t}(x,y)| \left| f(y) \right| dy \right)^{2} \frac{dt}{t} \right]^{1/2} \\ &\leq \left[\sum_{j \in \mathbb{Z}} \int_{2^{2j}}^{2^{2(j+1)}} 2^{-2js} \left(\int_{\mathbb{R}^{d}} 2^{-jd} \left(1 + \frac{|x| + |y|}{2^{j}} \right)^{-2} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \right)^{2} \frac{dt}{t} \right]^{1/2} \\ &\leq \sum_{j \in \mathbb{Z}} 2^{-js} \int_{\mathbb{R}^{d}} 2^{-jd} \left(1 + \frac{|x| + |y|}{2^{j}} \right)^{-2} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \end{split}$$

where in the last inequality we used the fact that $\ell_1 \hookrightarrow \ell_2$.

As in [23] we split the right hand side term above into two terms with respect to low-energy and high-energy cases.

$$\left(\int_{0}^{\infty} t^{-s} \left| \left(t\mathcal{L}_{a}e^{-t\mathcal{L}_{a}} + t\Delta e^{t\Delta} \right) f(x) \right|^{2} \frac{dt}{t} \right)^{1/2} \\ \leq \int_{\mathbb{R}^{d}} \sum_{j \in \mathbb{Z}: 2^{j} > |x| + |y|} 2^{-j(d+s)} \left(1 + \frac{|x| + |y|}{2^{j}} \right)^{-2} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \\ + \int_{\mathbb{R}^{d}} \sum_{j \in \mathbb{Z}: 2^{j} \le |x| + |y|} 2^{-j(d+s)} \left(1 + \frac{|x| + |y|}{2^{j}} \right)^{-2} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \\ + L(x)$$

 $:= I_1(x) + I_2(x).$

For the first term, we have

$$\begin{split} I_1(x) &\leq \int_{\mathbb{R}^d} \sum_{j \in \mathbb{Z}: 2^j > |x| + |y|} 2^{-j(d+s)} |f(y)| dy \lesssim \int_{\mathbb{R}^d} \frac{1}{(|x| + |y|)^{d+s}} |f(y)| dy \\ &\lesssim \int_{|y| \leq |x|} \frac{1}{(|x| + |y|)^{d+s}} |f(y)| dy + \int_{|y| > |x|} \frac{1}{(|x| + |y|)^{d+s}} |f(y)| dy \\ &\lesssim I_{11}(x) + I_{12}(x), \end{split}$$

which implies

$$\|I_1(\cdot)\|_{L^p_w} \lesssim \|I_{11}(\cdot)\|_{L^p_w} + \|I_{12}(\cdot)\|_{L^p_w}.$$

It is easy to see that

$$I_{11}(x) \le \int_{|y| \le |x|} \frac{|y|^s}{|x|^{d+s}} \frac{|f(y)|}{|y|^s} dy \le \mathcal{M}\left(\frac{f}{|\cdot|^s}\right)(x)$$

which yields that

$$\|I_{11}(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

Taking $g \in L^{p'}(v), v = w^{1-p'} \in A_{p'}$ we then have

$$\begin{split} \langle I_{12}(\cdot),g \rangle &\leq \int_{\mathbb{R}^d} \int_{|y| > |x|} \frac{1}{(|x| + |y|)^{d+s}} |f(y)| \, |g(x)| dy dx \\ &\leq \int_{\mathbb{R}^d} \int_{|y| > |x|} \frac{1}{|y|^d} \frac{|f(y)|}{|y|^s} \, |g(x)| dy dx \\ &\leq \int_{\mathbb{R}^d} \int_{|x| < |y|} \frac{1}{|y|^d} \frac{|f(y)|}{|y|^s} \, |g(x)| dx dy \lesssim \int_{\mathbb{R}^d} |\mathcal{M}g(y)| \left| \frac{f(y)}{|y|^s} \right| dy \\ &\lesssim \left\| \frac{f}{|\cdot|^s} \right\|_{L^p_w} \|\mathcal{M}g\|_{L^{p'}_v} \lesssim \left\| \frac{f}{|x|^s} \right\|_{L^p_w} \|g\|_{L^{p'}_v}. \end{split}$$

As a consequence,

$$\|I_{12}(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

We turn to the second term $I_2(x)$. For $0 < \epsilon < \frac{2-s}{2}$, we have

$$\begin{split} I_{2}(x) \lesssim & \int_{\mathbb{R}^{d}} \sum_{j \in \mathbb{Z}: 2^{j} \le |x| + |y|} \frac{1}{|x - y|^{d - \epsilon}} \frac{1}{(|x| + |y|)^{s + \epsilon}} \left(\frac{|x| + |y|}{2^{j}}\right)^{-2 + s + \epsilon} |f(y)| dy \\ \lesssim & \int_{\mathbb{R}^{d}} \frac{1}{|x - y|^{d - \epsilon}} \frac{1}{(|x| + |y|)^{s + \epsilon}} |f(y)| dy \lesssim \int_{\mathbb{R}^{d}} \frac{1}{|x - y|^{d - \epsilon}} \frac{1}{(|x| + |y|)^{\epsilon}} \frac{|f(y)|}{|y|^{s}} dy \\ \coloneqq & \int_{\Gamma_{1}(x)} \dots + \int_{\Gamma_{2}(x)} \dots + \int_{\Gamma_{3}(x)} \dots + \int_{\Gamma_{4}(x)} \dots \\ \coloneqq & I_{21}(x) + I_{22}(x) + I_{23}(x) + I_{24}(x), \end{split}$$

where $\Gamma_1(x) = \{y : |y| < |x/2|\}, \ \Gamma_2(x) = \{y : |y| \ge 2|x|\}, \ \Gamma_3(x) = \{y : |x|/2 \le |y| < 2|x|\} \cap B(x, |x|/2) \text{ and } \Gamma_4(x) = \{y : |x|/2 \le |y| < 2|x|\} \cap B(x, |x|/2)^c.$ It is even to dominate $L_1(x)$ as follows

It is easy to dominate $I_{21}(x)$ as follows

$$I_{21}(x) \lesssim \int_{|y| < |x|/2} \frac{1}{|x|^d} \frac{|f(y)|}{|y|^s} dy \lesssim \mathcal{M}\left(\frac{f}{|\cdot|^s}\right)(x),$$

which implies that

$$\|I_{21}(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

For the term $I_{22}(x)$, we have

$$I_{22}(x) \lesssim \int_{|y| \ge 2|x|} \frac{1}{|y|^d} \frac{|f(y)|}{|y|^s} dy.$$

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At this stage, by using the argument in the estimate $I_{12}(\cdot)$ we also get that

$$||I_{22}(\cdot)||_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

In addition, we have

$$I_{23}(x) \lesssim \int_{B(x,2|x|)} \frac{1}{|x|^d} \frac{|f(y)|}{|y|^s} dy \lesssim \mathcal{M}\left(\frac{f}{|\cdot|^s}\right)(x),$$

which implies that

$$\|I_{23}(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

The last term I_{24} can be dealt with as follows.

$$I_{24}(x) \lesssim \int_{B(x,|x|/2)} \frac{1}{|x-y|^{d-\epsilon}|x|^{\epsilon}} \frac{|f(y)|}{|y|^{s}} dy$$

$$\lesssim \sum_{j=0}^{\infty} \int_{2^{-j-1}|x| \le |x-y| < 2^{-j}|x|} \frac{1}{|x-y|^{d-\epsilon}|x|^{\epsilon}} \frac{|f(y)|}{|y|^{s}} dy$$

$$\lesssim \sum_{j=0}^{\infty} 2^{-j\epsilon} \mathcal{M}\left(\frac{f}{|\cdot|^{s}}\right)(x) \lesssim \mathcal{M}\left(\frac{f}{|\cdot|^{s}}\right)(x)$$

which implies that

$$||I_{24}(\cdot)||_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

Case 2: $-\left(\frac{d-2}{2}\right)^2 \leq a < 0$ Fix $p \in \left(1 \lor \frac{d}{d+s-\sigma}, d_{\sigma}\right)$ and $w \in A_{\frac{p}{1\lor \frac{d}{d+s-\sigma}}} \cap RH_{\left(\frac{d\sigma}{p}\right)'}$. Hence, there exist p_1, q_1 so that $1 \lor \frac{d}{d+s-\sigma} < p_1 < p < q_1 < d_{\sigma}$, and $w \in A_{\frac{p}{p_1}} \cap RH_{\left(\frac{q_1}{p}\right)'}$. Similarly to Case 1, by Proposition 4.4 we obtain that

$$\begin{split} \left(\int_{0}^{\infty} t^{-s} \left| \left(t\mathcal{L}_{a} e^{-t\mathcal{L}_{a}} + t\Delta e^{t\Delta} \right) f(x) \right|^{2} \frac{dt}{t} \right)^{1/2} \\ & \leq \int_{\mathbb{R}^{d}} \sum_{j \in \mathbb{Z}: 2^{j} \geq (|x| \wedge |y|)/2} 2^{-j(d+s)} \left(1 + \frac{2^{j}}{|x|} \right)^{\sigma} \left(1 + \frac{2^{j}}{|y|} \right)^{\sigma} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \\ & + \int_{\mathbb{R}^{d}} \sum_{j \in \mathbb{Z}: 2^{j} < (|x| \wedge |y|)/2} 2^{-j(d+s)} \left(1 + \frac{|x| + |y|}{2^{j}} \right)^{-2} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \\ & := J_{1}(x) + J_{2}(x). \end{split}$$

The argument used to estimate $I_2(x)$ in Case 1 also shows that

$$\|J_2(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

It remains to show that

$$\|J_1(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

Indeed, we have

$$\begin{split} J_1(x) &\leq \int_{\mathbb{R}^n} \sum_{j \in \mathbb{Z}: 2^j \geq (|x| \vee |y|)/2} 2^{-j(d+s)} \left(1 + \frac{2^j}{|x|}\right)^{\sigma} \left(1 + \frac{2^j}{|y|}\right)^{\sigma} e^{-\frac{|x-y|^2}{c2^{2j}}} |f(y)| dy \\ &+ \int_{\mathbb{R}^n} \sum_{j \in \mathbb{Z}: |y|/2 > 2^j \geq |x|/2} 2^{-j(d+s)} \left(1 + \frac{2^j}{|x|}\right)^{\sigma} \left(1 + \frac{2^j}{|y|}\right)^{\sigma} e^{-\frac{|x-y|^2}{c2^{2j}}} |f(y)| dy \\ &+ \int_{\mathbb{R}^n} \sum_{j \in \mathbb{Z}: |x|/2 > 2^j \geq |y|/2} 2^{-j(d+s)} \left(1 + \frac{2^j}{|x|}\right)^{\sigma} \left(1 + \frac{2^j}{|y|}\right)^{\sigma} e^{-\frac{|x-y|^2}{c2^{2j}}} |f(y)| dy \\ &:= J_{11}(x) + J_{12}(x) + J_{13}(x). \end{split}$$

For the term J_{11} , one has

$$J_{11}(x) \leq \int_{\mathbb{R}^n} \frac{1}{(|x| + |y|)^{d+s-2\sigma} |x|^{\sigma} |y|^{\sigma}} |f(y)| dy$$

$$\lesssim \int_{|y| \leq |x|} \dots + \int_{|y| \geq |x|} \dots$$

$$:= J_{11}^1(x) + J_{11}^2(x).$$

We now consider the contribution of $J_{11}^1(x)$. In this case, we have

$$J_{11}^1(x) \le \int_{|y| \le |x|} \frac{|y|^{s-\sigma}}{|x|^{d+s-\sigma}} \frac{|f(y)|}{|y|^s} dy$$

If $s - \sigma \ge 0$ then

$$J_{11}^{1}(x) \leq \int_{|y| \leq |x|} \frac{1}{|x|^{d}} \frac{|f(y)|}{|y|^{s}} dy \lesssim \mathcal{M}\left(\frac{f}{|\cdot|}\right)(x),$$

and hence,

$$\|J_{11}^{1}(\cdot)\|_{L^{p}_{w}} \lesssim \left\|\frac{f}{|x|^{s}}\right\|_{L^{p}_{w}}.$$

Otherwise, if $s-\sigma<0$ then by Hölder's inequality

$$J_{11}^{1}(x) \lesssim \left(\int_{|y| \le |x|} \frac{|y|^{(s-\sigma)p_{1}'}}{|x|^{(d+s-\sigma)p_{1}'}} dy \right)^{1/p_{1}'} \left(\int_{|y| \le |x|} \frac{|f(y)|^{p_{1}}}{|y|^{sp_{1}}} dy \right)^{1/p_{1}}.$$

Since $0 < (\sigma - s)p'_1 < d$, by Lemma 3.2, we have

$$J_{11}^{1}(x) \lesssim \left(\frac{1}{|x|^{d}} \int_{|y| \le |x|} \frac{|f(y)|^{p_{1}}}{|y|^{sp_{1}}} dy\right)^{1/p_{1}} \lesssim \mathcal{M}_{p_{1}}\left(\frac{f}{|\cdot|^{s}}\right)(x)$$

which implies that

$$\|J_{11}^1(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

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To estimate the term J_{11}^2 we employ the duality argument. Set $v = w^{1-p'}$, and hence by (vii) Lemma 2.1, $v \in A_{p/d'_{\sigma}}$. For $g \in L_v^{p'}$ we have

$$\begin{split} \langle J_{11}^2, g \rangle \lesssim & \int_{\mathbb{R}^n} \int_{|y| \ge |x|} \frac{1}{(|y|)^{d-\sigma} |x|^{\sigma}} \frac{|f(y)|}{|y|^s} dy |g(x)| dx \\ \lesssim & \int_{\mathbb{R}^d} \int_{|x| \le |y|} \frac{|g(x)|}{|x|^{\sigma}} dx \frac{|f(y)|}{|y|^{d+s-\sigma}} dy \\ \lesssim & \int_{\mathbb{R}^d} \left(\int_{|x| \le |y|} |g(x)|^{p_1'} dx \right)^{1/p_1'} \left(\int_{|x| \le |y|} \frac{1}{|x|^{\sigma p_1}} dx \right)^{1/p_1} \frac{|f(y)|}{|y|^{d+s-\sigma}} dy \end{split}$$

which together with Lemma 3.2 gives

$$\begin{aligned} \langle J_{11}^2, g \rangle &\lesssim \int_{\mathbb{R}^d} \left(\frac{1}{|y|^d} \int_{|x| \le |y|} |g(x)|^{p_1'} dx \right)^{1/p_1'} \frac{|f(y)|}{|y|^s} dy \\ &\lesssim \left\langle \mathcal{M}_{p_1'}g, \frac{f}{|\cdot|^s} \right\rangle \\ &\lesssim \left\| \mathcal{M}_{p_1'}g \right\|_{L_v^{p_1'}} \left\| \frac{f}{|\cdot|^s} \right\|_{L_w^p} \lesssim \left\| g \right\|_{L_v^{p_1'}} \left\| \frac{f}{|\cdot|^s} \right\|_{L_w^p}.\end{aligned}$$

Hence,

$$\|J_{11}^2(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

Let us move on the term J_{12} . We split this term as follows.

$$\begin{split} J_{12}(x) &\lesssim \int_{|y| \ge 2|x|} \sum_{j \in \mathbb{Z}: |y|/2 > 2^{j} \ge |x|/2} 2^{-j(d+s)} \left(1 + \frac{2^{j}}{|x|}\right)^{\sigma} \left(1 + \frac{2^{j}}{|y|}\right)^{\sigma} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \\ &+ \int_{|y| < 2|x|} \sum_{j \in \mathbb{Z}: |y|/2 > 2^{j} \ge |x|/2} 2^{-j(d+s)} \left(1 + \frac{2^{j}}{|x|}\right)^{\sigma} \left(1 + \frac{2^{j}}{|y|}\right)^{\sigma} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \\ &= J_{12}^{1}(x) + J_{12}^{2}(x). \end{split}$$

It is easy to see that

$$\begin{split} J_{12}^2(x) \lesssim & \int_{|y|<2|x|} \sum_{j \in \mathbb{Z}: |y|/2>2^j \ge |x|/2} 2^{-j(d+s)} \left(1 + \frac{2^j}{|x|}\right)^{\sigma} \left(1 + \frac{2^j}{|y|}\right)^{\sigma} e^{-\frac{|x-y|^2}{c2^{2j}}} |f(y)| dy \\ \lesssim & \int_{|y|<2|x|} \frac{1}{|x|^{d+s}} |f(y)| dy \lesssim M\left(\frac{f}{|\cdot|^s}\right)(x) \end{split}$$

which yields that

$$\|J_{12}^2(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}$$

To consider the contribution of J_{12}^1 , we write

$$\begin{split} J_{12}^{1}(x) &\lesssim \int_{|y| \ge 2|x|} \sum_{j \in \mathbb{Z}: |y|/2 > 2^{j} \ge |x|/2} 2^{-j(d+s)} \left(1 + \frac{2^{j}}{|x|}\right)^{\sigma} \left(1 + \frac{2^{j}}{|y|}\right)^{\sigma} e^{-\frac{|x-y|^{2}}{c2^{2j}}} |f(y)| dy \\ &\lesssim \int_{|y| \ge 2|x|} \frac{1}{|x-y|^{d+s-\sigma} |y|^{\sigma}} |f(y)| dy \\ &\lesssim \int_{|y| \ge 2|x|} \frac{1}{|y|^{d+s}} |f(y)| dy \end{split}$$

Arguing similarly to I_{12} in the case 1, we also obtain that

$$\|J_{12}^1(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

Hence,

$$\|J_{12}(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}$$

It remains to show that

$$\|J_{13}(\cdot)\|_{L^p_w} \lesssim \left\|\frac{f}{|x|^s}\right\|_{L^p_w}.$$

The proof of this estimate can be done in the same manner as that of J_{12} . We leave it to the interested reader.

This completes our proof.

4.3. Proof of Theorems 1.1 and Theorem 1.2.

Proof of Theorem 1.1: Fix 0 < s < 2, $d'_{\sigma} and <math>w \in A_{p/d'_{\sigma}} \cap RH_{(d_{s+\sigma}/p)'}$. Then by Theorem 4.2, Theorem 4.3 and Theorem 4.1 we have

$$\begin{split} \|(-\Delta)^{s/2}f\|_{L_{w}^{p}} &\lesssim \left\| \left(\int_{0}^{\infty} t^{-s} |t(-\Delta)e^{t\Delta}f|^{2} \frac{dt}{t} \right)^{1/2} \right\|_{L_{w}^{p}} \\ &\lesssim \left\| \left(\int_{0}^{\infty} t^{-s} \left| (t\mathcal{L}_{a}e^{-t\mathcal{L}_{a}} + t\Delta e^{t\Delta})f \right|^{2} \frac{dt}{t} \right)^{1/2} \right\|_{L_{w}^{p}} + \left\| \left(\int_{0}^{\infty} t^{-s} |t\mathcal{L}_{a}e^{-t\mathcal{L}_{a}}f|^{2} \frac{dt}{t} \right)^{1/2} \right\|_{L_{w}^{p}} \\ &\lesssim \left\| \frac{f}{|x|^{s}} \right\|_{L_{w}^{p}} + \|\mathcal{L}_{a}^{s/2}f\|_{L_{w}^{p}} \\ &\lesssim \|\mathcal{L}_{a}^{s/2}f\|_{L_{w}^{p}}. \end{split}$$

Conversely, for $1 with <math>p_1 := 1 \vee \frac{d}{d-\sigma} and <math>w \in A_{p/p_1} \cap RH_{(p_2/p)'}$ we have

$$\begin{split} \|\mathcal{L}_{a}^{s/2}f\|_{L_{w}^{p}} \lesssim \left\| \left(\int_{0}^{\infty} t^{-s} |t\mathcal{L}_{a}e^{-t\mathcal{L}_{a}}f|^{2} \frac{dt}{t} \right)^{1/2} \right\|_{L_{w}^{p}} \\ \lesssim \left\| \left(\int_{0}^{\infty} t^{-s} \left| (t\mathcal{L}_{a}e^{-t\mathcal{L}_{a}} + t\Delta e^{t\Delta})f \right|^{2} \frac{dt}{t} \right)^{1/2} \right\|_{L_{w}^{p}} + \left\| \left(\int_{0}^{\infty} t^{-s} |t(-\Delta)e^{t\Delta}f|^{2} \frac{dt}{t} \right)^{1/2} \right\|_{L_{w}^{p}} \\ \lesssim \left\| \frac{f}{|x|^{s}} \right\|_{L_{w}^{p}} + \| (-\Delta)^{s/2}f \|_{L_{w}^{p}} \\ \lesssim \| (-\Delta)^{s/2}f \|_{L_{w}^{p}}, \end{split}$$

where in the last inequality we used Theorem 4.1.

This completes our proof.

Proof of Theorem 1.2. Before starting the proof, we note that the flow $e^{it\mathcal{L}_a}$ satisfies for all s the conservation laws

$$\|\mathcal{L}_a^{s/2}e^{it\mathcal{L}_a}u_0\|_{L^2} = \|\mathcal{L}_a^{s/2}u_0\|_{L^2} \qquad \forall t \in \mathbb{R}$$

by self-adjointness. By an elementary application of Theorem 1.1 in the unweighted case, this implies the almost conservation of H^s norms

(29)
$$\| (-\Delta)^{s/2} e^{it\mathcal{L}_a} u_0 \|_{L^2} \simeq \| (-\Delta)^{s/2} u_0 \|_{L^2} \quad \text{for} \quad 0 < s < 2.$$

Consider not a generic Schrödinger equation with potential

$$iu_t + \Delta u - c(x)u = 0, \quad c(x) = \frac{a}{|x|^2}$$

If u(t,x) solves this equation, then the following identity holds for any sufficiently smooth ψ : $\mathbb{R}^d \to \mathbb{R}$:

(30)
$$\Re \nabla \cdot Q^S + \Im \partial_t \{ \bar{u} \nabla \psi \cdot \nabla u \} = -\frac{1}{2} \Delta^2 \psi |u|^2 + 2 \sum_{j,k=1}^d \partial_j u \partial_j \partial_k \psi \partial_k \bar{u} - \nabla \psi \cdot \nabla c |u|^2$$

where

$$Q^{S} = \nabla u \ T\overline{u} - \frac{1}{2}\nabla\Delta\psi|u|^{2} - \nabla\psi\left[c|u|^{2} - iu_{t}\overline{u} + \nabla u \cdot \nabla\overline{u}\right]$$

Formula (30) is usually called a *virial* (or *Morawetz*) identity, and it is easy to check directly by expanding the derivative $\nabla \cdot Q^S$, and using the equation for u(t, x) (see e.g. [2], [8], [10]).

Consider the case when the weight ψ is a radial function; by abuse of notation we use the same symbol $\psi(x) = \psi(|x|)$. Then we can write

$$\nabla \psi \cdot \nabla c = \psi' \partial_r c$$

where $\partial_r = \frac{x}{|x|} \cdot \nabla$ is the radial derivative. If we denote by $\nabla^T u = \nabla u - \frac{x}{|x|} \partial_r u$ the tangential component of ∇u , we have also the identities

$$|\nabla u|^2 = |\partial_r u|^2 + |\nabla^T u|^2 \quad \text{and} \quad \sum_{j,k=1}^d \partial_j u \partial_j \partial_k \psi \partial_k \bar{u} = \psi'' |\partial_r u|^2 + \frac{\psi'}{|x|} |\nabla^T u|^2.$$

Then formula (30) reduces to

(31)
$$\Re \nabla \cdot Q^{S} + \Im \partial_{t} \{ \bar{u}\psi' \partial_{r}u \} = 2\psi'' |\partial_{r}u|^{2} + 2\frac{\psi'}{|x|} |\nabla^{T}u|^{2} - \frac{1}{2}\Delta^{2}\psi|u|^{2} - \psi'\partial_{r}c|u|^{2}.$$

We now pick an explicit radial weight

$$\psi(r) = \int_0^r \frac{s^{\epsilon}}{1+s^{\epsilon}} ds, \qquad 0 < \epsilon < 1$$

where the parameter ϵ will be chosen later. A straightforward computation gives, writing for brevity r = |x|,

$$2\psi''|\partial_r u|^2 + 2\frac{\psi'}{|x|}|\nabla^T u|^2 = \frac{2\epsilon r^{\epsilon}}{(1+r^{\epsilon})^2}\frac{1}{r}|\partial_r u|^2 + \frac{2r^{\epsilon}}{1+r^{\epsilon}}\frac{1}{r}|\nabla^T u|^2 \ge \frac{2\epsilon r^{\epsilon}}{(1+r^{\epsilon})^2}\frac{1}{r}|\nabla u|^2,$$
$$-\psi'\partial_r c|u|^2 = \frac{r^{\epsilon}}{1+r^{\epsilon}}\frac{2a}{r^3}|u|^2$$

and

$$-\frac{1}{2}\Delta^2 \psi |u|^2 = \frac{1}{r^3} \frac{r^{\epsilon}}{1+r^{\epsilon}} \left[\frac{\mu_d}{2} + \epsilon \cdot \beta(r^{\epsilon}) \right] |u|^2, \qquad \mu_d = (d-1)(d-3),$$

where

$$2\beta(r) = -\frac{d^2 - 6d + 7}{1 + r} + \epsilon(2d - 5)\frac{r^2 - 1}{(1 + r)^3} - \epsilon^2 \frac{r^2 - 4r + 1}{(1 + r)^3} \quad \Longrightarrow \quad |\beta(r)| \le 3d^2$$

since $d \geq 3$. We substitute these expressions into (31) and integrate over \mathbb{R}^d ; if u is a H^1 solution the term in divergence form ∇Q^S vanishes, and we obtain

$$\Im \int \partial_t \{ \bar{u}\psi' \partial_r u \} dx \ge \int \frac{2\epsilon r^{\epsilon}}{(1+r^{\epsilon})^2} \frac{|\nabla u|^2}{r} dx + \int \left[2a + \frac{\mu_d}{2} - 3d^2\epsilon \right] \frac{r^{\epsilon}}{1+r^{\epsilon}} \frac{|u|^2}{r^3} dx.$$

Recall that by assumption we have

$$\delta := a + \left(\frac{d-2}{2}\right)^2 - \frac{1}{4} = a + \frac{\mu_d}{4} > 0,$$

thus if we choose $\epsilon = \min\{1, d^{-2}\delta/3\}$ and note that $\epsilon < \delta$ we have proved the inequality

(32)
$$\Im \int \partial_t \{ \bar{u}\psi'\partial_r u \} dx \ge \int \frac{\epsilon r^{\epsilon}}{(1+r^{\epsilon})^2} \frac{|\nabla u|^2}{r} dx + \int \frac{\epsilon r^{\epsilon}}{1+r^{\epsilon}} \frac{|u|^2}{r^3} dx$$

We next integrate (32) with respect to time on the interval $t \in [0, T]$; we obtain

(33)
$$\Im \int \bar{u}\psi'\partial_r u dx \Big|_{t=0}^{t=T} \ge \int_0^T \int \left[\frac{\epsilon r^{\epsilon}}{(1+r^{\epsilon})^2} \frac{|\nabla u|^2}{r} d + \frac{\epsilon r^{\epsilon}}{1+r^{\epsilon}} \frac{|u|^2}{r^3}\right] dx dt.$$

We note that the bilinear form

$$B(v,w) = \int \overline{v(x)} \psi'(|x|) \partial_r w(x) dx$$

satisfies the estimate

(34)
$$|B(u,v)| \le 3||v||_{\dot{H}^{1/2}} ||w||_{\dot{H}^{1/2}}.$$

where we used the notation

$$\|v\|_{\dot{H}^s} := \|(-\Delta)^{s/2}v\|_{L^2}.$$

Indeed, by Cauchy-Schwartz and the inequality $|\psi'| \leq 1$ we have

$$|B(v,w)| \le \|v\|_{L^2} \|w\|_{\dot{H}^1}$$

On the other hand, integrating by parts we have

$$B(v,w) = \int \overline{v(x)} \nabla \psi(x) \nabla w(x) dx = \int [w \overline{\nabla v} \nabla \psi + w \overline{v} \Delta \psi] dx.$$

Note that $|\nabla \psi| \leq 1$ while

$$|\Delta \psi| = \frac{\epsilon r^{\epsilon - 1}}{(1 + r^{\epsilon})^2} + \frac{(d - 1)r^{\epsilon - 1}}{1 + r^{\epsilon}} \le \frac{d}{r}$$

so that, by Hardy's inequality,

$$\left| \int w \overline{v} \Delta \psi dx \right| \le d \|w\|_{L^2} \left\| \frac{v}{|x|} \right\|_{L^2} \le \frac{2d}{d-2} \|w\|_{L^2} \|v\|_{\dot{H}^1}$$

Thus we obtain, for $d \ge 3$,

$$|B(v,w)| \le \left(1 + \frac{2d}{d-2}\right) \|w\|_{L^2} \|v\|_{\dot{H}^1} \le 7 \|w\|_{L^2} \|v\|_{\dot{H}^1}$$

Summing up we have proved that $B : \dot{H}^1 \times L^2 \to \mathbb{C}$ with norm 1 and $B : L^2 \times \dot{H}^1 \to \mathbb{C}$ with norm ≤ 7 . By complex bilinear interpolation this implies

$$B: \dot{H}^{1/2} \times \dot{H}^{1/2} \to \mathbb{C}$$

with norm $\sqrt{7} \leq 3$ as claimed. Using (34) we can write

$$\Im \int \bar{u}\psi' \partial_r u dx \Big|_{t=0}^{t=T} \le 4 \|u(T)\|_{\dot{H}^{1/2}}^2 + 4 \|u(0)\|_{\dot{H}^{1/2}}^2 \le C_0 \|u(0)\|_{\dot{H}^{1/2}}^2$$

where in the last inequality we used the almost conservation law (29), and the constant C_0 is independent to of T. Using the last inequality in (33) we obtain

$$\int_0^T \int \left[\frac{\epsilon r^{\epsilon}}{(1+r^{\epsilon})^2} \frac{|\nabla u|^2}{r} d + \frac{\epsilon r^{\epsilon}}{1+r^{\epsilon}} \frac{|u|^2}{r^3} \right] dx dt \le C_0 \|u(0)\|_{\dot{H}^{1/2}}^2$$

and letting $T \to +\infty$ we arrive at (4).

In order to prove (5), we isolate the first term in (4) and we use again Theorem 1.1 in the unweighted case:

$$\int \int \frac{r^{\epsilon-1}}{(1+r^{\epsilon})^2} |\nabla e^{it\mathcal{L}_a} f|^2 dx dt \le C\epsilon^{-1} \|(-\Delta)^{1/4} f\|_{L^2}^2 \simeq \epsilon^{-1} \|\mathcal{L}_a^{1/4} f\|_{L^2}^2.$$

We note that the weight $w(x) = \frac{r^{\epsilon-1}}{(1+r^{\epsilon})^2} = \frac{|x|^{\epsilon-1}}{(1+|x|^{\epsilon})^2}$ satisfies the conditions of Theorem 1.1, so that we can write

$$\int \frac{r^{\epsilon-1}}{(1+r^{\epsilon})^2} |\mathcal{L}_a^{1/2} e^{it\mathcal{L}_a} f|^2 dx \lesssim \int \frac{r^{\epsilon-1}}{(1+r^{\epsilon})^2} |\nabla e^{it\mathcal{L}_a} f|^2 dx \lesssim \epsilon^{-1} \|\mathcal{L}_a^{1/4} f\|_{L^2}^2 dx$$

Since $\mathcal{L}_a^{1/4}$ commutes with the flow, this implies

$$\int \frac{r^{\epsilon-1}}{(1+r^{\epsilon})^2} |\mathcal{L}_a^{1/4} e^{it\mathcal{L}_a} f|^2 dx \lesssim \epsilon^{-1} ||f||_{L^2}^2.$$

Finally, again by Theorem 1.1, we have

$$\int \frac{r^{\epsilon-1}}{(1+r^{\epsilon})^2} |\mathcal{L}_a^{1/4} e^{it\mathcal{L}_a} f|^2 dx \gtrsim \int \frac{r^{\epsilon-1}}{(1+r^{\epsilon})^2} |(-\Delta)^{1/4} e^{it\mathcal{L}_a} f|^2 dx$$

and this gives (5).

Denote now by R(z) the resolvent operator of \mathcal{L}_a and by $\Im R(z)$ its imaginary part:

$$R(z) = (\mathcal{L}_a - z)^{-1}, \qquad \Im R(z) = (2i)^{-1} (R(z) - R(\bar{z})).$$

Moreover, let A be the operator

$$A = w(x)^{1/2} (-\Delta)^{1/4}, \qquad w(x) = \frac{|x|^{\epsilon - 1}}{(1 + |x|^{\epsilon})^2}.$$

Estimate (5) can be written

$$||Ae^{it\mathcal{L}_a}f||_{L^2(\mathbb{R}^{d+1})} \lesssim \epsilon^{-1/2}||f||_{L^2}.$$

By Kato smoothing theory, applying e.g. Theorem 2.2 in [7], we obtain that this estimate is equivalent to the resolvent estimate

(35)
$$\|A\Im R(z)A^*f\|_{L^2(\mathbb{R}^d)} \lesssim \epsilon^{-1} \|f\|_{L^2}, \qquad z \notin \mathbb{R},$$

uniformly in $z \notin \mathbb{R}$. (In the terminology of Kato's theory, the closed operator A is \mathcal{L}_a -smoothing). Then we are in position to apply Theorem 2.4 from [7] (with $\nu = 0$) and we obtain that the operator $A\mathcal{L}_a^{-1/4}$ is $\mathcal{L}_a^{1/2}$ -smoothing, i.e., the following estimate holds:

$$\|Ae^{it\mathcal{L}_{a}^{1/2}}f\|_{L^{2}(\mathbb{R}^{d+1})} \lesssim \epsilon^{-1/2}\|\mathcal{L}_{a}^{1/4}f\|_{L^{2}}$$

that is to say, we have proved that

$$\|w(x)^{1/2}(-\Delta)^{1/4}e^{it\mathcal{L}_a^{1/2}}f\|_{L^2(\mathbb{R}^{d+1})} \lesssim \epsilon^{-1/2}\|\mathcal{L}_a^{1/4}f\|_{L^2}$$

Finally, using Theorem 1.1 exactly as in the proof of (5), we can cancel the operators $(-\Delta)^{1/4}$ and $\mathcal{L}_a^{1/4}$, and we obtain estimate (6).

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