

Nanostructured Semiconductor Materials for Dye-Sensitized Solar Cells

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Since their first report in 1991 by O'Regan and Grätzel, Dye-Sensitized Solar Cells (DSSCs) appeared immediately as a promising low-cost photovoltaic technology. In fact, though being far less efficient than conventional silicon-based photovoltaics (being the maximum, lab scale prototype reported efficiency around 13%), the simple design of the device and the absence of the strict and expensive manufacturing processes needed for conventional photovoltaics, make them attractive in small-power applications especially in low-light conditions, where they outperform their silicon counterparts. Nanomaterials are at the very heart of a DSSC, being the success of its design due to the use of nanostructures at both the anode and the cathode.

In this review we present the state of the art for both n-type and p-type semiconductors used in the photoelectrodes of DSSCs, showing the evolution of the materials during the 25 years of history of this kind of devices. In the case of p-type semiconductors, also some other energy conversion applications are touched upon.

Part I: n-type semiconductors

Introduction

A dye-sensitized solar cell (DSSC) is a semiconductor-based photovoltaic device that directly converts both artificial and natural (solar) radiation into electric current. In contrast to the conventional systems where the semiconductor assumes both the tasks of light absorption and charge carrier separation and transport, the two functions are separated in a DSSC. In a conventional DSSC, light is absorbed by a sensitizer, which is anchored to the surface of a wide band gap n-type semiconductor^[1]. Often referred to as a Grätzel cell, this hybrid device was first reported to the scientific world in 1991 by the seminal publication of Brian O'Regan and Michael Grätzel presenting a device made of sensitized nanocrystalline TiO₂ with a power conversion efficiency of 7.1%^[2]. DSSCs constitute probably the cheapest photovoltaic technology today available and their efficiency has been constantly improved in the last 25 years. The confirmed efficiency record is 11.9%^[3], achieved by the Sharp company. However, it was Mathew *et al.* that reported the best performing DSSC in literature up to now. DSSCs fabricated utilizing the [Co(bpy)₃]^{2+/3+} as redox couple and the SM315 dye, demonstrated panchromatic light harvesting without the use of co-sensitization, leading to a record efficiency value of 13% under one sun illumination^[4]. DSSC is an eco-friendly technology that can be used to produce electricity in a wide range of light conditions, indoor and outdoor. If one considers that dye sensitized solar cells work properly also under low irradiation conditions, the efficiencies they reach make the technology extremely interesting with respect to conventional solar cells in such conditions. The DSSCs and their inventor received prestigious awards, including the Balzan Prize in 2009 and the 2010 Millennium Technology Prize, the most important technology prize in the world^[5].

Device structure and operation principles: brief discussion of components and their role in DSSC.

A simple scheme of the components and of the basic operating principles of a DSSC is reported in figure 1. The device is comprised of four components only:

- Nanostructured n-type semiconductor (wide band gap metal oxide) coated over transparent conducting substrate (TCO (ITO^[6] or FTO^[7])).

- Visible-light absorber dye (several organic dye can be used, such as N3^[8], N719^[9], N749^[10] (the so-called black dye), K8^[11], K19^[12], CYC-B11^[13] and C101^[2]).
- Electrolyte^[14, 15].
- Counter electrode (TCO coated with a platinum layer or other suitable catalyst).

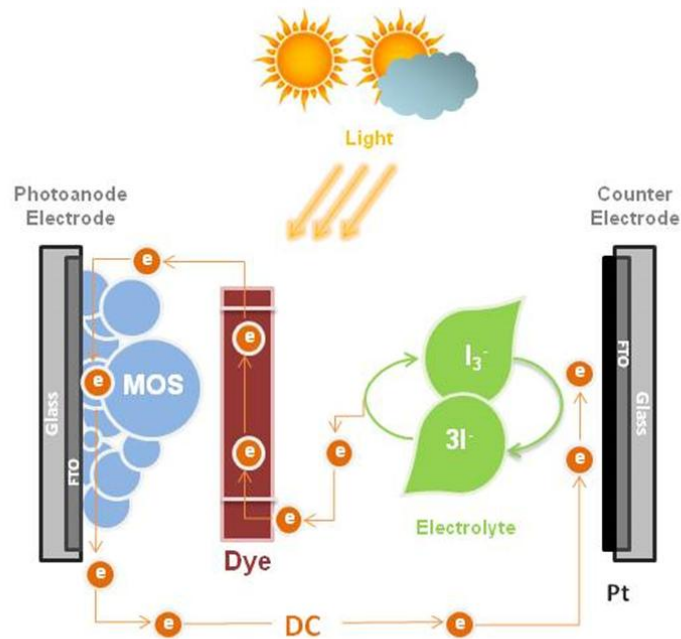


Figura 1. Device structure of a DSSC.

In the case of n-type materials, current is generated when a photon absorbed by a dye molecule causes an electron injection into the conduction band of the semiconductor. The dye molecules are anchored to the surface of the semiconductor particles by a chemical bond (sensitization). In a nanostructured oxide the semiconductor provides a large surface area for dye molecule anchoring. The photons are absorbed at the surface of the nanostructure by the dye, promoting its excitation and then the electron injection. The photogenerated electrons are injected into metal oxide and the dye is regenerated by a redox species present in the electrolyte. The electrons move through the semiconductor to a current collector and then to an external circuit. The redox mediator is regenerated at the cathode, and the process is cyclic. Each part of the device heavily determines the cost and the efficiency of DSSCs. The optimization of light absorption properties can be achieved by modifying the dye alone, while charge transport properties can be improved by the optimization

of the semiconductor and of the electrolyte composition^[16]. Thus, in the last years almost all research efforts have been focused on the modification of each component for practical applications. Ye *et al.*^[17] reported a scheme of the number of publications for sensitized solar cells. The major area of interest was represented by the fabrication of n-type nanostructured semiconductor photoanodes with improved architectures for high dye loading and fast electron transport. Therefore, this section of the present paper aims at bringing together the various contributions brought about by researchers to improve the efficiency of the n-type sensitized semiconductors for DSSCs. Moreover, the detailed discussion of the role of band gap, morphology, composition and doping of n-type semiconductors for the development of efficient photoanodes will be reported. Several contributions are summarized here with the description of nanostructures (0D, 1D, 2D, 3D and mesoporous nanostructures) and of the effect of doping of photoanodes materials (mostly TiO₂ and ZnO) on the performances of DSSCs.

Photoanode materials for n-type DSSC.

Semiconducting nanostructured films are the main core of DSSCs photoanodes. The main issues in DSSCs are the charge recombination processes and their relatively low light harvesting capability. The photoanode performs a dual function as both the support for the sensitizer and carrier of photogenerated electrons from the sensitizer to the external circuit^[17]. The photoanode material should not absorb visible light and should have sufficiently high surface area for optimum dye adsorption.

In semiconductor terminology, the top of the valence band and the bottom of the conduction band are called the *valence band edge* and the *conduction band edge*, respectively. The energy of the conduction band edge is denoted as E_{CB} ; similarly, the energy of the valence band edge is denoted as E_{VB} . The energy difference between these levels is called the *band gap* (E_g). The size of this band gap is perhaps the most important property of a semiconductor, as it influences all of the most important electronic properties of the material^[18]. The conduction band edge E_{CB} of the photoanode material should match with that of the excited dye molecules. To efficiently collect the photogenerated electrons, the photoanode should have high charge carrier mobility.

Desirable properties of the material are also the ease to preparation, stability, low cost and environmental friendliness. These properties are the defining characteristics of an ideal photoanode^[19].

N-type semiconductors

Probably it was Michael Faraday^[20] that made the first significant observation in the semiconductor science in 1833 when he discovered the negative temperature resistivity coefficient of silver sulfide^[21]. Nowadays, semiconductors are defined by the unique behaviour of their electrical properties. As their name implies, semiconductors possess an electrical conductivity intermediate between conductors such as metals and insulators such as ceramics^[22]. For high-density electron ensembles such as valence electrons in metals, Fermi statistics are applicable. The Fermi level, E_F (defined at 0 K as the energy at which the probability of finding an electron is 1/2) can be regarded as the electrochemical potential of the electron in a particular phase (in this case, a solid). Thus, all electronic energy levels below E_F are occupied and those above E_F are likely to be empty. Electrons in semiconductors may be regarded as low-density particle ensembles such that their occupancy in the valence (VB) and conduction bands (CB) may be approximated by the Boltzmann function^[23, 24]:

$$n_e \approx N_0 \exp\left(-\frac{E_0 - E_F}{kT}\right) \quad (1)$$

Now we come to another important distinction between metals and semiconductors in that two types of electronic carriers (n_i, p_i) are possible in the latter. Consider the thermal excitation of an electron from VB to CB . This gives rise to a free electron in the CB and a vacancy or hole in the VB . Thus, eq (1) becomes:

$$n_i \approx N_c \exp\left(-\frac{E_F - E_{CB}}{kT}\right) \quad (2)$$

$$p_i \approx N_v \exp\left(-\frac{E_{VB} - E_F}{kT}\right) \quad (3)$$

Where N_c and N_v are the effective density of states (in cm^{-3}) at the lower edge and top edge of CB and VB , respectively. These expressions can be combined with the recognition that $n_i = p_i$ to yield

$$n_i^2 \approx N_0 \exp\left(-\frac{E_{VB} - E_{CB}}{kT}\right) \approx N_0 \exp\left(-\frac{E_g}{kT}\right) \quad (4)$$

To provide a numerical sense of the situation, N_c and N_v are typically both approximately 10^{19} cm^{-3} so that the constant $N_0(N_c N_v)$ in eq (1) is about 10^{38} cm^{-3} .

The latter case refers to the semiconductor in its intrinsic state with very low carrier concentrations under ambient conditions. The Fermi level, E_F , in this case lies approximately in the middle of the energy band gap^[25]. This simply reflects the fact that the probability of the electron occupancy is very high in VB and very low in CB and does not imply an energy level at E_F itself that can be occupied.

In extrinsic semiconductors the carrier concentrations are perturbed such that $n \neq p$. The analogy with the addition of an acid or base to water is quite instructive here. Consider the case when donor impurities are added to an intrinsic semiconductor. Since the intrinsic carrier concentrations are low (sub-parts per trillion), even additions in part per billion levels can have a profound electrical effect. This process is known as doping of the semiconductor. In this particular case, the Fermi level shifts towards the one of the bands' edges. When the donor level is within a few kT in energy from the CB edge, appreciable electron concentrations are generated by donor ionization process (at ambient temperatures) such that now $n \gg p$. This is termed $n - type$ doping, and the resultant (extrinsic) semiconductor is called $n - type$. By analogy, $p - type$ semiconductor have $p \gg n$. The terms minority and majority carriers now become appropriate in these cases. For a $p - type$ semiconductor case, the Fermi level now lies close to VB edge. The Fermi level determines the tendency of the semiconductor to transfer charges to other phases. The movement of E_F with dopant concentration can also be rationalized via the Nernst formalism^[26].

Doping can be accomplished by adding aliovalent impurities to the intrinsic semiconductor. For example, P (a group 15 or VB element) will act as a donor in Si (a group 14 or IVB element). This can be rationalized on chemical terms by noting that P need only four valence electrons for tetrahedral bonding (as in Si lattice) the fifth electron is available for donation by each P atom. The donor density, N_D nominally is approximately 10^{17} cm^{-3} . Thus, assuming that $n \cong N_D$ (complete ionization at 300 K), p will be only approximately 10^3 cm^{-3} [recall the product $n_i p_i \sim 10^{20} \text{ cm}^{-6}$], bearing out the earlier qualitative assertion that $n \gg p$.

Impurity addition, however, is not the only doping mechanism. Non-stoichiometry in compound semiconductors such as CdTe also gives rise to n - or p - type behaviour, depending on whether Cd or Te is in slight excess, respectively. For metal oxides, doping with suitable cation/anion modifies band gap and may modulate their electrical properties. The defect chemistry in these solid chalcogenides controls their conductivity and doping in a complex manner. Excellent treatises are available on this topic and on the solid-state chemistry of semiconductors in general^[27]. The

distinction between metal and semiconductor electrodes is important when we consider the electrostatics across the corresponding solid-liquid interfaces (MOS/electrolyte in DSSC).

N-type semiconductor-electrolyte interface

The electrochemical potential of electrons in a redox electrolyte is given by the Nernst expression

$$E_{redox} = E_{redox}^0 + \frac{RT}{nF} \ln \left[\frac{c_{ox}}{c_{red}} \right] \quad (5)$$

where c_{ox} and c_{red} are the concentrations of the oxidized and reduced species, respectively, in the redox couple. The parameter ($E_{redox} = \mu_{e,redox}$) as defined by this equation can be identified with the Fermi level ($E_{F,redox}$) in the electrolyte. When a semiconductor is immersed in this redox electrolyte, the electrochemical potential (Fermi level) is disparate across the interface. Equilibration of this interface thus necessitates the flow of charge from one phase to the other and a “band bending” ensues within the semiconductor phase. For an n-type semiconductor (TiO_2) electrode at open circuit, the Fermi level is typically higher than the redox potential of the electrolyte, and hence electrons will be transferred from the electrode into the solution. Therefore, there is a positive charge region, and this is reflected in an upward bending of the band edges ^[28] (see figure 2). Since the majority charge carrier of the semiconductor has been removed from this region, the latter is also referred to as the *depletion layer*.

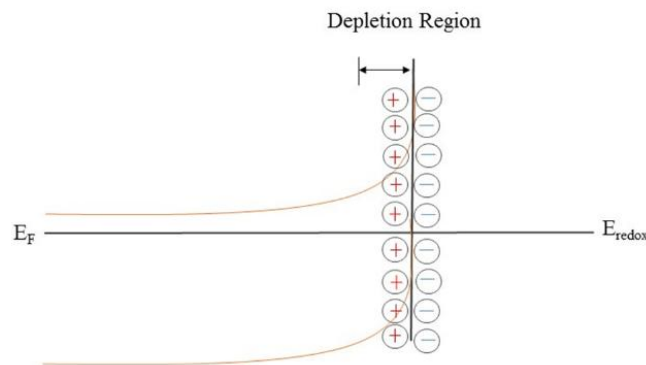


Figure 2. Band bending for n-type semiconductor in equilibrium with electrolyte.

N-type nanostructured semiconductor for DSSCs

Nanostructured materials present a specific surface area significantly larger than that of the corresponding bulk materials. The nanoscale size may also effect the behaviour of electrons transport in nanostructures in view of a limit to the electron mean free path. This is called *quantum confinement effect* ^[29-31]. The quantum confinement effect is observed when the size of the particle is too small to be comparable to the wavelength of the electron. On the optics side, by forming photonic band gap, periodic nanostructures (know as a photonic crystals) also show to be special in light manipulation and management through generating *optical confinement* or *photonic localization* ^[32-35]. These unique properties of nanomaterials received considerable attention and have been extensively investigated for applications in electronic, optoelectronic, photovoltaic, photocatalytic, and sensing devices ^[36-40]. The metal oxides in their nano-form can be synthesized under various morphologies with different shapes and sizes thus offering the possibility for modulating their properties.

The crystallinity of the metal oxides should be high to prevent the recombination of e^- and h^+ . Ohtani *et al.* clarified that the photocatalytic activity of TiO₂ powders strongly depends on its physical properties such as crystal structure, surface area, particle size, surface hydroxyls and so on ^[41]. In literature, a variety of preparation techniques, such as sol-gel ^[42], hydrothermal/solvothermal ^[43,44], electrochemical anodization ^[45], electrospinning ^[46-47], spray pyrolysis ^[48] and atomic layer deposition ^[49] have been developed and applied to obtain different morphologies in photoanode materials. Asim *et al.* summarized the different preparation and deposition methods, which have been used for photoanode materials of DSSCs, emphasizing their advantages and disadvantages, in order to allow a researcher to carefully choose and optimize a given method ^[50]. Perez-Page *et al.* presented several methods based on templates for shape-controlled nanostructuring. A variety of nanostructured materials are produced through template-based syntheses, including zero-dimensional (nanoparticles), one-dimensional (nanowires, nanotubes,...), two dimensional (nanoflakes and nanosheets) structures ^[51]. In 1995, Antonelli and Ying for the first time reported the synthesis of mesoporous TiO₂, which was accomplished through a modified sol-gel process involving TIP (titanium isopropoxide) as a precursor ^[52]. Later, various procedures have been developed for the synthesis of mesoporous materials. In 2001 Grosso *et al.* ^[53], Yun *et al.* ^[54] and Hwang *et al.* ^[55] reported syntheses of mesoporous titania films in the anatase phase. In the development of a simple and general methodology for the synthesis of highly organized mesoporous metal oxide, Sadatlu *et al.* proposed a novel and effective strategy ^[56].

Nanoparticles-based semiconductors (OD nanostructures)

A typical DSSC photoanode is composed of nanocrystalline semiconductors. Sufficient light absorption is achieved by the nanocrystalline form of the semiconductor, because a large internal surface area increases the dye concentration in the film per unit device area. Wide band gap metal oxide semiconductors (MOS) ($E_g > 3\text{eV}$)^[57] such as TiO_2 ^[58-60], ZnO ^[61-67], SnO_2 ^[60, 68, 69-71] and Nb_2O_5 ^[72-74], have been studied more or less extensively and used as photoanode materials for DSSC devices. These MOSs present good stability against photocorrosion, transparency in the major part of the solar spectrum, and good electronic properties^[75-78]. Photocorrosion, which is caused by the oxidation by holes (generated through band gap excitation) of the redox species in the electrolytes, may affect the performance of the semiconductor.

TiO_2 is a low cost, widely available, non-toxic and biocompatible material. It has been used in health care products as well as in domestic applications such as paint pigments^[16]. On the other hand, ZnO , which has a similar conduction band edge and work function compared to TiO_2 , but with a higher carrier mobility than TiO_2 , was considered as a promising photoanode materials for DSSCs. However, the instability of ZnO in acidic environment and formation of dye aggregates on its surface deteriorate its performances^[79]. This is probably due to its acidic surface and more positive (vs NHE) conduction band edge position compared to TiO_2 . Some other semiconductor materials, such as Zn_2SnO_4 ^[80], CeO_2 ^[81], WO_3 ^[82], SrTiO_3 ^[83], In_2O_3 ^[84] and $\text{Al}_{64}\text{Cu}_{25}\text{Fe}_{11}$ ^[85] were also studied. Alami *et al.* evaluated the suitability of the $\text{Al}_{64}\text{Cu}_{25}\text{Fe}_{11}$ intermetallic compound to enhance spectral solar absorption when replacing the mesoporous layer material of DSSCs. Zheng *et al.*^[86] presented for the first time a complete study on WO_3 as photoanode material for DSSC.

The efficiency of DSSCs using these materials with a suitable redox electrolyte should be around 10-15% under one sun irradiation. The main issues in DSSCs are charge recombination processes and their relatively low light harvesting efficiency. As stated before, the photoanode performs a dual function as both the support for the sensitizer and carrier of photogenerated electrons from the sensitizer to external circuit^[17]. In this regard, a major benefit of using nanostructured materials in DSSCs is that the resulting photoelectrode films are highly porous. Porous nanostructures offer a larger surface area for dye absorption compared to a bulk material. Among all semiconductor oxides, porous, anatase phase titanium dioxide (TiO_2) is the best and most common choice for high-efficiency DSSCs due to its low-cost, easy synthesis, non-toxicity and biocompatibility^[87]. A photoanode based on nanosized TiO_2 crystals with various geometries such as nanoparticles^[88], ordered meso-structured materials^[89, 90] and one-dimensional structured materials e.g. nanorods, nanowires and nanotubes^[91,92] have been extensively studied. Up to now, the most efficient DSSCs are based on TiO_2 nanoparticles owing to their integrated advantage in

porosity, dye absorption, charge transfer, and electron transport. Moreover, electron transport in nanoparticles films suffer from trapping/detrapping processes^[93]. Trapping is an important effect affecting the electron lifetime in the semiconductor, significantly reducing the electron diffusion length, Le . In order to have a quantitative collection of electrons at the anode, the condition $Le \gg d$ should be satisfied, with d being the semiconductor film thickness. Quaranta *et al.*^[94] showed that Le can be increased by a given quantity of functionalized multiwalled carbon nanotubes (MWCNTs) coated by nanostructured TiO₂-anatase. Moreover, other 0D structured photoanodes were used to prevent this drawback of nanoparticles. Chen *et al.*^[95] prepared a core-shell structure that consisted of a nanoporous TiO₂ film coated with oxides such as Nb₂O₅, ZnO, SrTiO₃, ZrO, Al₂O₃ and SnO₂. Moreover, Palomares *et al.* explained the growth of conformal metal oxide insulating overlayers (SiO₂, Al₂O₃, and ZrO₂ overlayers) on preformed nanocrystalline TiO₂ films, focusing on the ability of such insulating layers to retard interfacial recombination dynamics and thereby modulate the performance of DSSCs fabricated using these films^[96].

The best performance in DSSCs has been achieved with the anatase form of TiO₂^[97, 98]. On the other hand, relatively little attention has been paid to the rutile form of TiO₂. Park *et al.*^[99] studied a rutile TiO₂ -based DSSC^[100]. Rutile is the thermodynamic stable polymorph of TiO₂ that forms when less stable polymorphs are heated at temperatures over 700 °C. Therefore, the synthesis temperature of rutile TiO₂ is a potential issue for use in DSSCs since nanocrystalline morphology is hardly preserved at high temperatures. Moreover, the anatase form is preferred because of its higher band gap (3.2 eV) compared to rutile (~3 eV). High band gap energy makes anatase chemically more stable in DSSCs^[101].

Fan *et al.* reported band gap positions and values of several commonly used semiconductors, as is shown in figure 3:

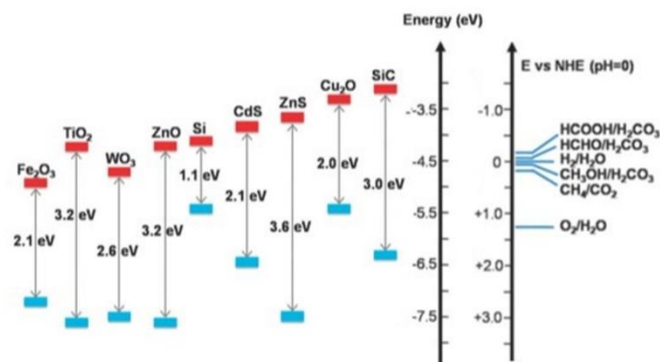


Figure 3. Band gap positions of several semiconductors relative to energy of some redox couples in contact with aqueous electrolyte at $pH=0$ ^[Adapted from 102].

The solid state physics community has adopted the electron energy in vacuum as a reference, whereas chemists have traditionally used the normal hydrogen electrode (NHE) as energy reference. NHE lies at -4.5 eV with respect to the vacuum level. We can now relate the redox potential E_{redox} (as defined with reference to NHE) with the Fermi level E_F :

$$E_{F,redox} = -4.5 \text{ eV} - eE_{redox} \quad (6)$$

Parameters related to the band gap position are the flat-band voltage (V_{FB}) and flat-band capacitance (C_{FB}). Material properties such as oxide charge, semiconductor carrier density, and the effective work function of the metal can be obtained directly from V_{FB} and C_{FB} ^[103, 104].

The standard method for determining the flat-band parameters is through Mott-Schottky plots.

Generally, a three electrode single compartment configuration was adopted for capacitance analysis used to obtain Mott-Schottky plots. MOS coated on conducting glass was used as working electrode while platinum and standard calomel electrode (SCE) were used as counter and reference electrodes respectively. To measure the capacitance, oxide coated plates were immersed in 0.5 M Na₂SO₄ solution and the capacitances were measured as a function of the potential across the space charge layer at different frequencies ^[105]. Shifting of the flat-band potentials of semiconductors plays an important role determining their photoelectrochemical properties as thin oxide films and powders. The shifting of V_{FB} can be harnessed in solar energy conversion processes in several ways. In DSSCs, the open circuit voltage (V_{OC}) is determined by the difference between the quasi-Fermi level of electrons in the oxide film and the energy of the redox couple in the electrolyte ^[106]. Consequently, the shift of V_{FB} potential in the negative direction results in increase in barrier height and hence in the attainable V_{OC} of the solar cell. Similarly, the flat band potential can be shifted negatively in order to electrolyze a redox couple that is otherwise too positive in the electrochemical series to be reactive at the illuminated semiconductor surface ^[107, 108]. Park *et al.* also reported the relative flat band potential and conduction band energies of some oxides such as TiO₂, ZnO, SnO₂, Nb₂O₅, FeTiO₃ and Ta₂O₅ and sulphides that could be used for DSSCs when considering only conduction band energy positions ^[Adapted from 99] (figure 4).

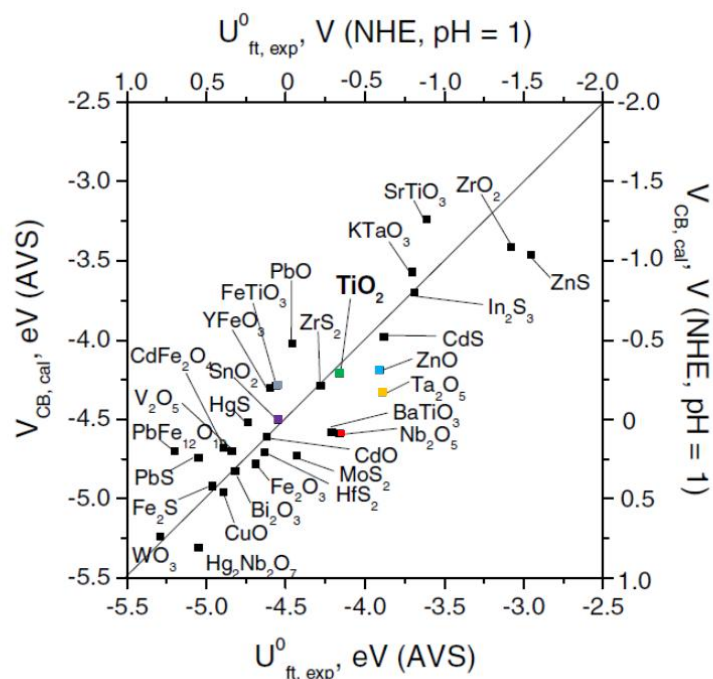


Figure 4. Conduction band energies and flat band potentials of metal oxides and sulphides at pH=1 (NHE) or vacuum level^[Adapted from 99].

High ordered nanostructures (1D nanostructures: nanowires, nanotubes, etc.)

A major feature that discriminates various types of nanostructures is their dimensionality. Low-dimensional nanostructures have an aligned structure that can act as a single crystal, thus facilitating the rapid electron transport in unidirectional manner and show potentiality for obtaining high performance devices^[109]. Zhang *et al.* classify the nanostructures, presenting nanoparticles advantages (offer large surface area to photoelectrode film for dye absorption) and drawbacks (electron transport by trapping and detrapping processes, which may result in energy losses). In this regard, one-dimensional nanostructures such as nanowires and nanotubes provide direct pathways resulting in a much faster electron transport than in the nanoparticles films^[110]. However, nanoparticles are advantageous in providing large surface area for dye absorption. Law *et al.*^[111], for the first time in 2005, presented a comparison between a DSSC made with a ZnO films made of nanoparticles and nanowires to demonstrate that one-dimensional nanostructures may provide direct pathways for electron transport in DSSCs. A high-performance nanowire photoanode must have a

large surface area for dye adsorption, comparable to that of a nanoparticle film. At one sun irradiation the device presented an efficiency value between 1.2–1.5%. In the same year, Baxter and Aydil^[112] first fabricated ZnO nanorod-based DSSC by using a hydrothermal method; they achieved an overall photoelectric conversion efficiency of 0.5%. Since then, 1D nanostructures of ZnO (nanorods and similar: Martinson *et al.*, 2007^[113]; Schlur *et al.*, 2013^[114]; Sadia Ameen *et al.*, 2012^[115]; McCune *et al.*, 2012^[116]; Guo *et al.*, 2013^[117]) have been widely investigated as photoelectrodes for enhancing the DSSCs performances^[118]. Other ZnO based 1D structures have been investigated for photoanodes. Kim *et al.* studied the performance of ZnO nanofibers^[119], Tan *et al.* constructed ZnO nanorods arrays on ITO glass^[120] and Yang *et al.* prepared ZnO nanotips on rough Zn microtip foil^[121].

ZnO based DSSCs with delicately designed nanostructures, from irregular microrods to nanosheets with simply adjusting the dosage of capping agents in the hydrothermal synthesis process, were fabricated to study the photovoltaic performances vs. morphology. In figure 5, SEM micrographs of as-synthesized ZnO structures evolving from microrods to nanosheets are reported^[122].

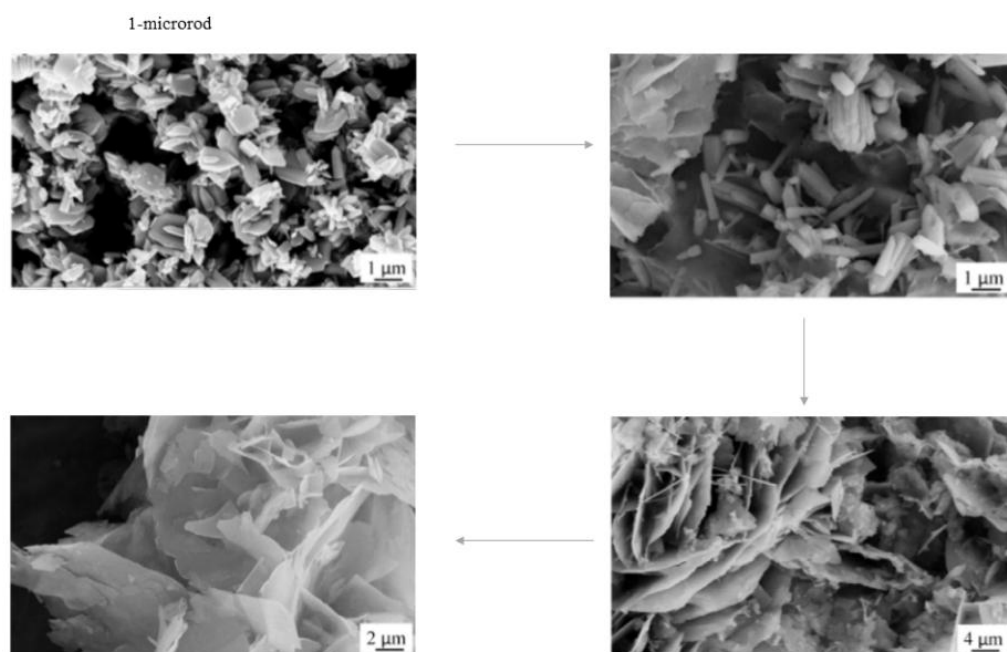


Figure 5. SEM micrographs of as-synthesized ZnO structures evolving from microrod to nanosheet with the dosage of HMTA ((CH₂)₆N₄) from 0 g to 1.2 g^[122].

TiO₂ nanowires in standard DSSC configuration gave an overall conversion efficiency of 5%, which is much higher than those reported for ZnO nanowires^[66, 123].

In addition, the other TiO₂ 1D nanostructures gave higher efficiencies than ZnO based ones. TiO₂ nanorods, for example, were synthesized by Chen *et al.* via a microemulsion electrospinning technique. The authors observed a power conversion efficiency of 8.53%, which originated from the large structure of nanorods^[124].

The most studied class of one-dimensional nanostructures is that of nanotubes. The recombination rate for TiO₂ nanotubes is ~10 times lower than that of nanoparticles^[125]. Their hollow structure usually gives a larger surface area than that nanowires or nanorods. Roy *et al.* presented a complete review of the current status of the use of TiO₂ nanotubes in Grätzel cells^[126]. Since the first results in 2005 with efficiencies of 0.04%, reproducible efficiency results in the order of 4–5% have been obtained. The best morphology reported is called bamboo-type rings. This may be achieved in various ways by modifying the tube walls with porous materials (double-walled tubes). During this last year, Momemi presented a study on Cr-doped TiO₂ nanotubes^[127]. The effect of chromium doping on the photovoltaic efficiency of dye-sensitized solar cells (DSSCs) with TiO₂ nanotubes prepared by an anodization procedure followed by an annealing process was investigated. He reported an interesting comparison between different doped nanotubes. In table 1 the obtained results are reported.

<i>Photoanodes</i>	$J_{SC}/(mA \cdot cm^{-2})$	V_{OC}/V	<i>FF</i>	η (%)	<i>References</i>
TNs	7.21	0.65	0.45	2.13	[128]
CdSe/TNs	6.19	0.438	0.495	1.56	[129]
P3HT/CdS/TNs	3.00	0.7	0.55	1.16	[130]
CdS/TNs	5.17	0.77	0.47	1.87	[131]
ZnS/CdSe/CdS/TNs	13.52	0.48	0.53	3.44	[132]
WO ₃ /TNs	6.74	0.72	0.483	2.343	[133]
Cr/TNs	10.56	0.69	0.534	3.895	[127]

Table 1. Comparison of DSSC photovoltaic parameters of Cr-doped nanotubes with similar studies as Momemi reported^[127].

ZnO nanotubes have been also studied for application in DSSCs. The utility of high surface area Al-doped ZnO (AZO) nanotubes for DSSC application was studied by Martinson *et al.*^[113]. Compared to similar ZnO-based DSSCs, the AZO nanotubes show superior photovoltage and fill factors with photoconversion efficiencies up to 1.6%. Xie *et al.* prepared a novel type of coaxial TiO₂/ZnO nanotube arrays for DSSCs. The obtained efficiency is 2.8%, due to enhanced charge separation effect in this structural arrangement^[134]. A smaller portion of literature is devoted to other TiO₂ one-dimensional nanostructures, for example nanospindles^[135, 136], electrospinning materials^[137] and nano-embossed hollow structure (NeHS-TiO₂)^[138-140]. In particular, one-dimensional hierarchical structure with hollow or porous interiors are considered promising structures due to the longitudinal pathways of the hollow or porous interior structures and the 1D interconnected nanocrystal network that allows for efficient electron transport^[141]. Wang *et al.* reported 1D titania with tuneable structural hierarchy that exhibits about 1.3-1.5 times higher power conversion efficiency than commercial P25 as photoanode material for DSSCs^[142]. Zhu *et al.* synthesized ZnO nanorods-nanosheets (NR-NS) hierarchical architecture^[143].

Combination of 1D nanostructures (2D nanostructures)

A combination of nanowires, nanotubes and nanoparticles create 2D nanostructures that are significant for their sufficiently high surface area available for dye loading. In the hybrid structures, the resulting photoelectrode film presents the advantages of the one-dimensional nanostructures and the exposed surface area could offer more space for dye absorption and electron conveyance. Most research was performed on the array films of ZnO and TiO₂ nanowires or TiO₂ nanotubes filled with ZnO or TiO₂ nanoparticles^[144]. In table 2 some 2D hybrid nanostructures used as photoanode of DSSC are reported.

Type	η (%)	Ref [n°]
TiO ₂ nanorod array	1.90	[145]
TiO ₂ nanoflower	1.53	[145]
TiO ₂ nanodisk	6.6%	[146]
TiO ₂ bilayer nanoribbons	5.6%	[147]
TiO ₂ nanoflakes	8.2%	[148]
Fill-up ZnO nanowire array with nanoparticles	2.2-3.2%	[149]
	(0.5-0.8% ZnO nanowire only)	[150]
Indium-doped ZnO (IZO) nanosheet	~7%	[151]

Table 2. Some combination of 1D nanostructures and their respective photoconversion efficiencies as reported in literature.

3D nanostructures (hierarchical nanostructures, composite, multilayer and thin film materials)

As we have seen until now, electron transport in nanoparticle films is via a random route with multiple trapping and detrapping events. This mechanism seriously affects the electron diffusion coefficient and the electron collection efficiency. Moreover, small sized particles show inefficient light scattering ability, which results in poorer light-harvesting efficiency. On the other hand, 1D and combination of 1D (2D) nanostructures attracted much attention in DSSCs owing to the improvement of the charge collection efficiency, which ascribes to a faster electron transport and slower recombination rate. However, the low surface area of 1D nanostructures hinders significant improvement of the photovoltaic performance owing to poor dye loading compared to nanoparticles. For a better balance of those three factors (surface area, light-scattering and electron transport), bi-/trifunctional hierarchical structured materials consisting of nanoparticles/nanorods, have been widely investigated as photoanodes for DSSCs. These structures demonstrated enhanced photovoltaic performance due to their superior dye adsorption properties, light scattering ability and faster electron transport compared to traditional nanoparticle only-based DSSCs^[152]. Liao *et al.* in 2011 presented the effect of TiO₂ morphology on photovoltaic performance of DSSCs with particular attention on hierarchical structures^[153]. Chen and Yang^[154] presented a series of works in which they demonstrated the viability of new photoanode architectures in DSSCs comprising building blocks of ZnO nanoplates^[155], and TiO₂ nanosplindles^[156]. Recently, they presented a systematic study about DSSCs based on ZnO nanotetrapods-based photoanodes^[157] and ZnO nanotetrapods/SnO₂ nanoparticles composite photoanode^[158, 159]. Other authors report ZnO 3D

nanostructure (e.g. caterpillar-like nanorods ($\eta \sim 5.2\%$)^[160], nanodendrites ($\eta \sim 3.74\%$)^[161]) based photoanodes that show significant photoconversion efficiencies for DSSC application.

Grätzel *et al.*^[162] introduced a novel photoanode with a nanoscale structure that resembled a forest, fabricated by pulsed laser deposition (PLD) under relatively high background gas pressures. This new nanostructure, which they named “nanoforest,” replaces the traditional random nanoparticle oxide network by vertically aligned bundles of TiO₂ oxide nanocrystals. Grätzel speculated earlier that a structure combining the high surface area of nanoparticles with the electron transport directionality of vertical rods would be optimal for DSSCs^[16].

Another approach to the improvement of DSSC is the enhancement of the conduction band level relative to the electrolyte redox/oxidation level that would result in higher photovoltage^[163]. Nb₂O₅ is an example of one such n-type transition metal oxide with a higher conduction band edge than TiO₂. Ghosh *et al.*, for the first time reported the fabrication of a Nb₂O₅-based photoanode with a nanoforest structure and the results of investigations of the photovoltaic properties of this material under a variety of key growth parameters (layer thickness, background gas pressure and composition)^[73].

Other efficient TiO₂ 3D nanostructures producing good efficiency values are branching nanorods ($\eta \sim 7.17\%$)^[164], dendritic hollow structures ($\eta \sim 7.16\%$)^[165], bridge nanotubes ($\eta \sim 6.17\%$)^[166], etc. Photoanodes made of two or more materials have attracted attention due to the obvious advantages of combining different materials. Yan *et al.*^[167] combined the high electron transport rate of ZnO and the high electron injection efficiency of TiO₂ with Ru-based dyes. Huu *et al.*^[168] presented a hierarchical photoanode comprising a SnO₂ nanoparticle underlayer and a ZnO nanorod overlayer and its photovoltaic performance was compared to photoanodes consisting of SnO₂ nanoparticles only and ZnO nanorods only.

Satapathi *et al.* presented a DSSC fabricated using photoanodes made from graphene-TiO₂ nanocomposites^[169]. The relationship between the size of graphene sheet and the cell performance was explored. It was concluded that the cells loaded with the smaller graphene sheets yielded larger enhancement. The maximum efficiency was obtained by 184 nm sheets with $\eta = 6.62\%$.

Also composite TiO₂ anatase materials with CaTiO₃^[170], and BaTiO₃^[171, 172] have been applied to electrodes for DSSCs. Among perovskite-type compound, BaTiO₃ is a n-type semiconductor with band gap energy of 3.3 eV. Such band gap value may increase the open circuit voltage of a DSSC as Suzuki and Kijima reported in 2005^[173].

Manoharam and Venkatachalam investigated the photoelectrochemical performances of DSSCs based on aluminium doped titanium dioxide (Al-TiO₂) nanoparticles/nanowires (TNPWs)

composite photoanode layers, with different electrolytes. The maximum obtained efficiency is 7.26%^[174].

An efficiency value of 3.44% has been reported by Xu *et al.*^[175] using as double scattering layer a ZnO film consisting of ZnO monodisperse light scattering layer and a submicrometer-sized plate-like ZnO film as overlayer in the photoanode of a DSSC. Ghanbari *et al.* presented a new strategy for improving the efficiency of TiO₂ DSSCs by design of a new double-layer film doped with Zn ions, with various morphologies^[176].

Mesoporous structures (3D oxide aggregate nanostructures)

Among 3D nanostructures, mesoporous MOS-based materials (in particular titania-based) with a crystalline framework, high specific surface area and tuneable pore size have received significant attention for energy conversion applications. Several morphologies such as mesoporous TiO₂ films^[177], beads^[178], monoliths^[179], networks^[180], with controlled porosity have been prepared via different synthetic strategies. The most interesting are monodisperse TiO₂ beads with a submicrometric-sized diameter. Chen *et al.*^[181] reported the synthesis of monodisperse mesoporous anatase titania beads with high surface area and tunable pore size and grain diameter via a combination of sol-gel and solvothermal processes in presence of hexadecylamine (HDA) as a structure-directing agent. In this way they enhanced the light-harvesting capability of TiO₂ electrode without sacrificing the accessible surface for dye loading. These mesoporous TiO₂ beads have a diameter of less than 1 μm and are composed of anatase TiO₂ nanocrystals. The beads have been used in the preparation of photoanodes for DSSCs and an improved efficiency was observed when compared to analogous cells prepared using standard Degussa P25 TiO₂ photoanodes of similar thickness. An overall light conversion efficiency of 7.20% (open-circuit voltage (V_{oc}) 777mV, short-circuit current density (I_{SC}) 12.79 $mAcm^{-2}$ and fill factor (FF) 0.72) was achieved using the mesoporous TiO₂ bead electrodes^[182]. The substantial improvement of J_{sc} and η for the hierarchical sphere-based DSSC when compared to other nanostructure-based DSSCs is mainly due to the larger dye loading, higher light scattering ability, faster charge transport, and longer electron lifetime^[183, 184]. Submicrometric mesoporous TiO₂ beads are also used to form a scattering layer on top of a transparent, 6-mm-thick, nanocrystalline TiO₂ film. Later, Grätzel and Chen presented DSSCs with photoanodes made of mesoporous TiO₂ beads that achieved a power conversion efficiency over 10%^[185]. Kim *et al.*^[186] reported a two-step method for the synthesis of TiO₂ aggregates which first produced TiO₂ spheres via a controlled hydrolysis and then etched the spheres under hydrothermal conditions. These aggregates, applied to DSSCs, yielded an efficiency as high as 10.5%.

In addition, ZnO spheres were used to enhance the performances of ZnO-based DSSCs. These latter were synthesized by hydrolysis of a zinc salt in polyol medium^[187]. A significant difference in the conversion efficiency values (2.4 % and 5.4% for ZnO nanoparticles and aggregates, respectively) has been verified. The explanation of this difference is based on the consideration that ZnO is not stable in an acidic dye solution and the formation of a Zn²⁺/ dye complex on the surface of ZnO may seriously hinder the electron injection process from the dye molecules to the semiconductor^[188]. In the case of film made of ZnO nanoparticles, the formation of complex on the film surface would block the pores and result in an incomplete infiltration of dye molecules^[189]. However, the situation is quite different in the case of aggregates. Due to the existence of large pores among the submicron-sized aggregates, the dye penetration can be accomplished in a very short time (for example, 30 min) prior to the formation of the complex layer^[187, 190]. Other mesoporous structures have been investigated for DSSCs applications. The use of preformed mesoporous solid as hard templates allows for the preparation of novel mesostructures of TiO₂ with high crystallinity. Moreover, hard templates can provide an excellent support and confinement framework to prevent collapse of mesoporous structures, thus allowing for high crystallinity. All these characteristics are found in mesoporous single crystals (MSC) of anatase, which were first prepared by Crossland *et al.*^[191]. The MSC anatase displays higher conductivity and electron mobility than those of conventionally used nanoparticles. In solid-state DSSCs they obtained a power conversion efficiency of 7.2%. Latini *et al.*^[192] used MSC as photoanode for conventional DSSCs in a comparison between hard and soft template based materials. Five different anatase phase mesoporous titanias were used as photoanodes in DSSCs: two materials were synthesized by using silica nanospheres (hard template) and the other three using two different organic polymeric templating agents, P123 and Brij 58 (soft template). Unexpectedly, MSCs are the least suitable for application in DSSCs because of their low specific surface area. Finally, Zhao *et al.* developed a simple surfactant-sulfuric acid carbonization method to synthesize ultrastable ordered mesoporous titania with high crystallinity^[193].

Doping and co-doping effect on nanocrystalline and mesoporous structures.

In addition to the optimization of the oxide nanostructure, the combined use of different metal oxides and/or the use of doped materials with higher conduction band energies should, in principle, allow further improvement of DSSCs performances. In metal oxides, the doping with suitable cation/anion modifies the band gap and may modulate their electrical properties. The properties of titania, including anatase to rutile phase transition as well as the photoactivity greatly depend on the presence of dopants, i.e. cations^[194] and anions^[195], impurities, as well as on the crystallinity^[196, 197],

grain size, surface area, and so on, in addition to the presence of such substances as amorphous silica at the interfaces and/or grain boundaries. All these properties have a profound influence on the mass transfer rate, diffusion, and crystallite growth of titania in composite nanoparticles^[198].

The ionic radii of niobium Nb^{5+} , aluminum Al^{3+} , gallium Ga^{3+} , yttrium Y^{3+} and scandium Sc^{3+} are quite close to that of titanium Ti^{4+} . Solid solutions of these cations in titania have been prepared and characterized^[199]. For example, Grätzel *et al.* studied the effect Ga^{3+} and Y^{3+} in mesoporous anatase and of Nb^{5+} in nanocrystalline anatase for DSSCs photoanodes^[200, 201]. Latini *et al.*^[105] reported the beneficial effect on the performances of DSSCs of Sc^{3+} doping of mesoporous anatase beads. At 0.2% of Sc atoms a maximum efficiency of 9.6% was obtained, which was found to be 6.7% greater than the efficiency of DSSCs with pure anatase. Li *et al.*^[202] reported a DSSC based on TiO_2 photoanode that was modified by the Al-doped TiO_2 layers using the chemical bath deposition method. The Al-doped TiO_2 layer improved the photocurrent density. The effect of a co-dopant M (M= Ga, Al, Sc) on the formation, crystallite growth, optical band gap, photocatalytic activity, and phase stability of anatase-type TiO_2 solid solutions ($\text{Ti}_{1-2x}\text{Nb}_x\text{M}_x\text{O}_2$) containing the same amount of dopant Nb that were directly formed as nanoparticles under mild hydrothermal conditions at 180°C for 5h, was investigated by Hirano and Ito^[203]. The effect of lanthanum ions (La^{3+}) on charge trapping in dye-sensitized solar cell (DSSC) photoanodes has been investigated with doped and surface-treated TiO_2 nanoparticles^[204]. Including those mentioned before, transition metal dopants in general could extend the photoresponse in the visible light region. This is due to low energy photon excitations of the corresponding metal oxide clusters with smaller band gap and partially from the excitations of the introduced localized states in the band gap of doped TiO_2 ^[205].

On the other hand, nonmetal dopants (N, C, B, S, P, etc.) can exist as isolated atoms rather than clusters, which have greater potential for realizing visible-light photoactivity^[206]. In anion doped TiO_2 the formation in the oxide phase of new defect levels slightly above the valence band generally cause the red shift of absorption spectrum^[207-211]. Pan *et al.*, for example, presented the enhanced efficiency of DSSCs by trace amount Ca-doping of TiO_2 photoelectrodes. The best photovoltaic performance was obtained from 50 ppm Ca-doping with a conversion efficiency of 7.45%^[212]. In table 3 some useful dopants for anatase TiO_2 have been grouped together as well as their respective photoconversion efficiencies as reported in literature.

Doped-TiO ₂	η (%)	Reference
B	6.1%	[213]
Ce	7.65%	[214]
Cr	8.4	[215]
Mg	1.2%	[216]
N	6.25%	[217]
Nb	7.41%, 7.8%, 8.7%	[218-220]
Sb	8.13%	[221]
Sn	8.31%	[222]
Tm/Yb	7.5%	[223]
Ta	7.1%	[224]
W	9.10%, 7.42%	[104, 225]

Table 3. DSSC efficiency of several doped-TiO₂ photoanodes.

The amount of papers dedicated to systematic studies of the effect of heteroatoms and their concentration in the anatase lattice on the performances of DSSCs are far less numerous. Cavallo *et al.*^[226] have undertaken a systematic work of synthesis, characterization and test of DSSC photoanodes containing mesoporous anatase bead doped with RE (rare earth) cations (Nd, Sm, Gd, Er and Yb). The maximum efficiency has been obtained for the sample containing 0.2% Er metal atoms (8.7%).

Similarly, papers dealing with doping of ZnO with different cations and anions are also present in literature, as reported in table 4.

Doped-ZnO	η (%)	Reference
Al	0.28%, 0.964%	[227, 228]
Ga	4.01%	[229]
I	4.5%	[230]
K	0.012%	[231]
Mg	4.19%, 4.11%	[232, 233]
N	5%	[234]

Table 4. DSSC efficiency of some doped-ZnO photoanodes.

Part II: p-type semiconductors

Introduction

In the recent past there has been a crescent interest towards the development and the realization of solar radiation conversion devices in which an electrochemical process of reduction is photoactivated via the absorption of light by a cathode either in the pristine^[235-239] or in a modified/sensitized state^[240-266]. These devices are photoelectrochemical cells (PECs)^[40, 87] with photoactive cathodes consisting of *p*-type semiconducting materials, and include *p*-type dye-sensitized solar cells (*p*-DSCs) (Figure 6)^[241, 242, 247, 248, 266], tandem DSCs (*t*-DSCs) (Figure 7)^[267-271] cells of photoelectrolysis for non fossil fuels production, namely molecular H₂ in the process of water splitting (Figure 8)^[272-283], and for carbon dioxide redox reduction (Figure 9)^[284-315]. The initial charge separation produced as a consequence of light absorption^[314,315] can occur either directly on the *p*-type semiconductor (Figure 9), or on the electrically connected sensitizer which accomplishes successively the transfer of charge according to the vectoriality imparted by the relative positions of the energy levels of the electronic states involved, and by the kinetics of the possible redox reactions (Figures 6-8).

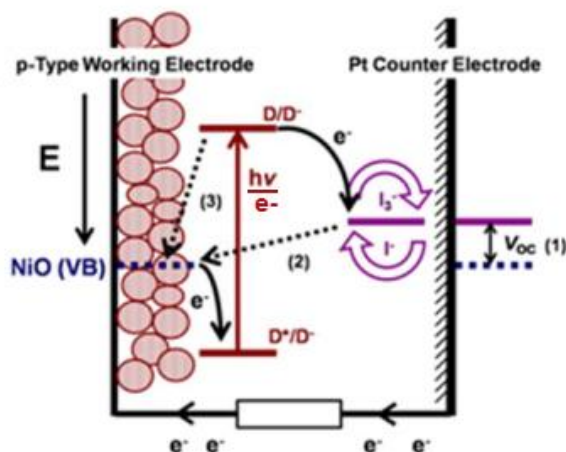


Figure 6. Diagram of the electrical potential E levels involved in the process of photoactivated electron transfer (et) at the basis of the operation of a p-DSC. The separation evidenced by the double arrow (1) represents the open circuit voltage V_{OC} of the p-DSC. The latter parameter is determined by the difference between the Nernst level of the redox couple (I_3^-/I^- , in the given example), and the Fermi level of the p-type semiconducting electrode (NiO in the reported example). Bent full arrows represent the dye-mediated et from the p-type cathode to the oxidized form of the redox couple. The et process is catalysed by the dye sensitizer D through the absorption of light with frequency ν . Dotted arrows (2) and (3) describe two processes of electronic recombination following the separation of charge induced by light absorption.

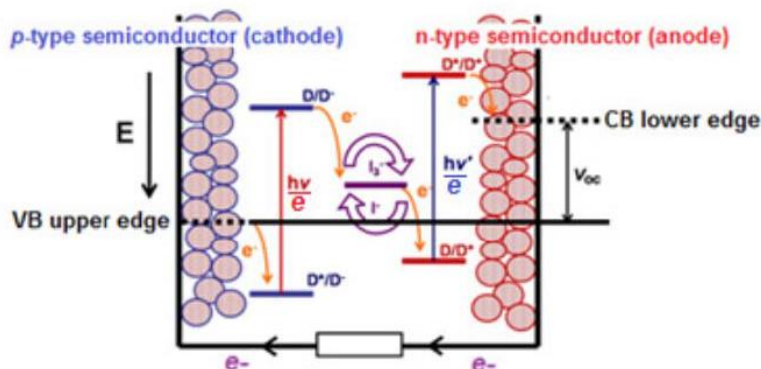


Figure 7. Diagram of the electrical potential E levels involved in the process of photoactivated electron transfer (et) at the basis of the operation of a t-DSC. The separation evidenced by the double arrow (1) represents the open circuit voltage V_{OC} of the t-DSC. The latter parameter is determined by the difference between the Fermi levels of p-type cathode and n-type anode. Bent full arrows represent the dye-mediated et processes from the photoelectrodes to the corresponding forms of the redox couple (I_3^-/I^- , in the given example). The et processes are catalysed by the dye sensitizers D and D with complementary light absorption properties.

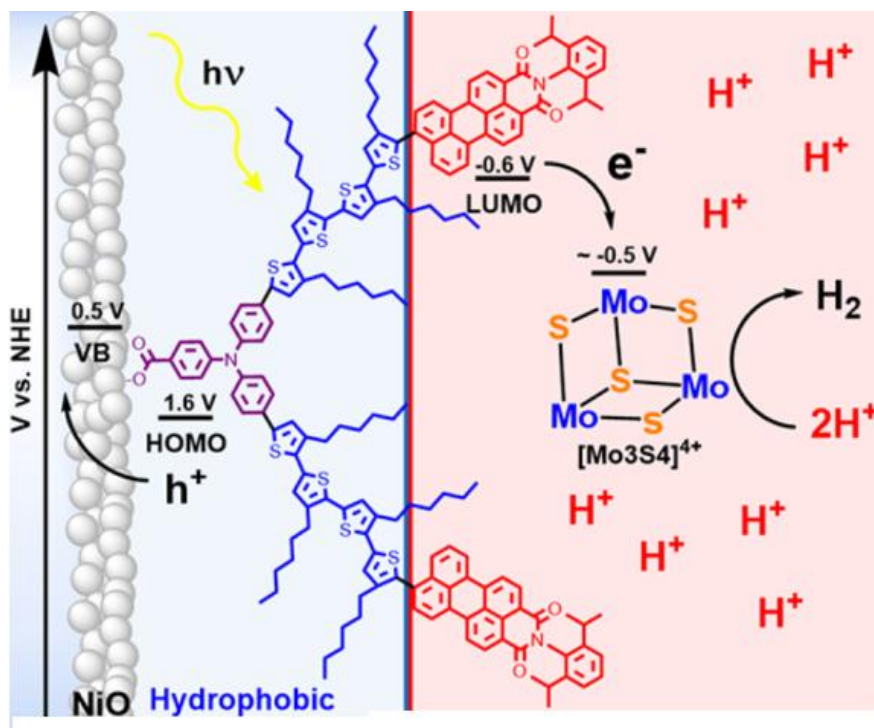


Figure 8. Depiction of the mechanism of photoactivated e^- from p -type NiO to H^+ mediated by the multifunctional dye-sensitizer PMI-4T-TPA for the production of H_2 fuel. The organic dye is immobilized onto NiO surface through the carboxylic group (in violet). The dye-sensitizer acts as electron donor (through the PMI moiety, in red) towards the molecular co-catalyst Mo_3S_4 , and behaves as electron acceptor towards NiO. The hydrophobic hexyl groups (in blue) exert a blocking effect against the direct discharge of H^+ onto bare NiO cathode. Adapted from ref. 282.

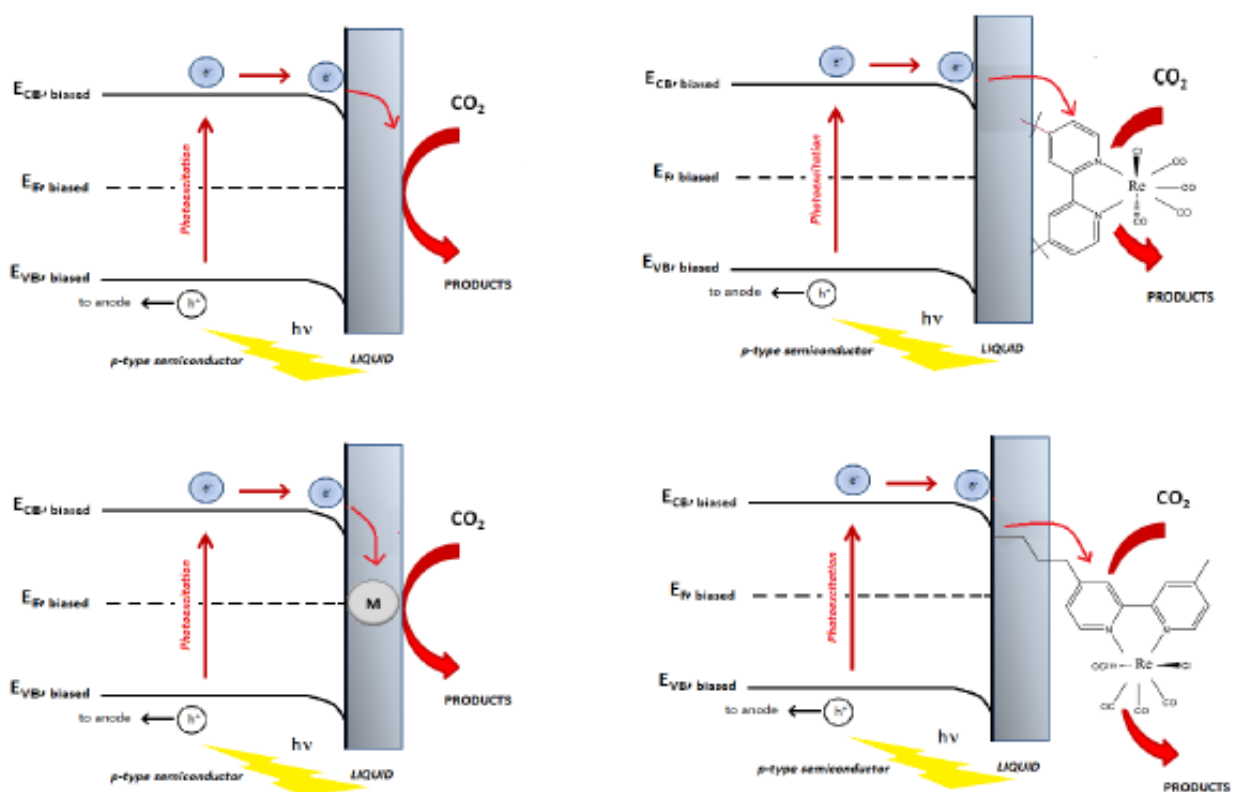


Figure 9. Mechanisms of photoactivated CO_2 reduction on a *p*-type semiconductor electrode: (top left) heterogeneous photoelectrocatalysis occurring directly at the semiconductor electrode; (top right) homogeneous catalysis mediated by a molecular catalyst in the electrolyte; (bottom left) heterogeneous photoelectrocatalysis directly on a metal-decorated semiconductor electrode; (bottom right) heterogeneous catalysis mediated by a molecular catalyst anchored on the semiconductor electrode. In all four examples photoexcitation occurs at the semiconducting cathode. E_{VB} , E_{CB} and E_{F} indicate the upper edge of the semiconductor valence band (VB), the lower edge of the semiconductor conduction band (CB), and the Fermi level of the *p*-type semiconductor, respectively. The illuminated electrode is reverse biased in these four examples. Adapted from ref. 284

Several requirements of thermodynamic as well as kinetics character must be fulfilled in order to realize an efficient process of photoconversion into the desired product. The latter can be the electrical power as in case of the *p*-DSCs and *t*-DSCs (Figures 6 and 7), and fuels/chemicals in case of the photoelectrolytic cells based on the working principles schematized in Figures 8 and 9. The most crucial aspect generally limiting the performance of radiation conversion devices is the one related to the effectiveness of charge separation. This is because the separation of charge, once effectuated, can evolve either (a) to the displacement of the charges towards the sites/reactants for the realization of the desired redox transformations (reductions in the present context), or (b) lead to the occurrence of unwanted charge recombination/trapping phenomena depending on the relative rates of processes (a) and (b), and the chemical-physical features of the various interfaces on which the events of charge separation take place^[316-318]. For this reason is fundamental to envisage photoconversion devices with rapid mechanisms of photogenerated charge displacement within the semiconductor^[319]. A relatively recent example of that has been the invention of the DSC by Grätzel^[2]. The latter represents a photoconversion device in which the electrons photoinjected at a sensitized photoanode of TiO₂ in the rutile form^[320] are swiftly drifted to the metallic cathode for the occurrence of a fast interfacial *et* from the cathode to the anode-generated oxidized form of a redox couple. Fundamentally, this successful result is based on very fast charge transport within the photoactive semiconducting electrode. This would prevent recombination between the photocharge generated at the surface of the light-absorbing semiconducting electrode and the redox species in the electrolyte initially oxidized/reduced by the photoinjected charge^[321]. The most recent advancements in the design and realization of PECs of *p*-type have been achieved when the photoelectroactive *p*-type semiconductors had nanostructured features (Figure 10)^[250, 267, 322-328]. This is motivated by the fact that a nanostructured surface increases enormously the effective area of contact for the interfaces created by the nanostructured semiconducting material with respect to the same system possessing a compact morphology. This difference improves the kinetics of those electrochemical processes occurring at an extended interface in favour to the nanostructured version^[329].

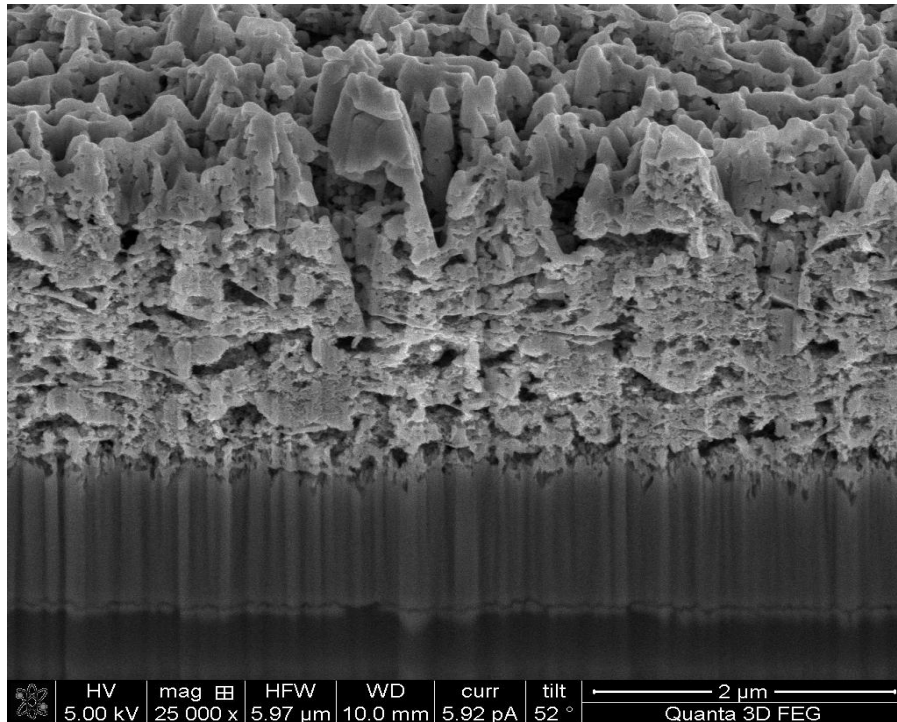


Figure 10. Image of the cross section of a nanostructured cathode made of nickel oxide. The picture has been obtained with the dual beam focused ion beam–scanning electron microscope (FIB-SEM). Adapted from ref. 329.

Another important motive of interest in adopting nanostructured semiconductor electrodes resides in the possibility of modifying the optical properties of the pristine system through sensitization with very large amounts of dye-sensitizers anchored per unit surface^[330, 331]. These colouring additives sometime represent the sole actual light harvesting species in case the semiconductor has a wide bandgap with respect to the energy of the luminous radiation^[332]. On the other hand, the main concern that might arise by the utilization of nanostructured semiconductors regards the delineation of an opportune pattern of energy bands (or a set of discrete energy levels), which still allows the electrical connection through the nanostructured semiconductor^[333, 334]. In fact, a nanostructured semiconductor represents a system characterized by not having an internal electric field because of the inexistence of a charge depletion layer^[235-237]. In this type of semiconductors the typically discontinuous pattern of frontier energy levels is also associated to electronic states delocalized disorderly at a variable extent^[238]. Consequently, the latter feature imposes a charge transport mechanism of variable range hopping between localized states^[239-342], through which the displacement of charge is originated by diffusion^[343-347].

In this review we will report the most recent developments and progresses achieved with PECs having photoactive cathodes (vide supra). Despite the fast advancement of the technology of perovskite solar cells[348-382], in the present contribution this type of device will be not reviewed since the photocurrent generated in perovskite based photoconversion devices is not associated to the primary event of a photostimulated redox reactions but is rather a consequence of a light-driven process of charge separation at a p-i-n junction. As a such, it does not introduce any redox states variations in the conducting materials during ordinary operative conditions of perovskite systems [383-385].

Devices for the conversion of the electromagnetic radiation into electrical power with photoelectroactive *p*-type semiconductors: analysis of performances recently reported in the literature

DSCs of p-type and t-type

One of the first studies on sensitized *p*-type semiconductors with nanostructured features for DSCs of *p*-type was reported in 2000^[267]. In this seminal work nanoporous nickel oxide (NiO) prepared via sol-gel was employed as cathode while erythrosine B was the dye-sensitizer. This prototypical *p*-DSC displayed a photovoltage of about 80 mV, and a photocurrent density of *ca.* 0.2 mA cm⁻² the photocurrent being originated by associated the photoinjection of mobile holes in the VB of NiO after photoexcitation of the chemisorbed colorant (Figure 6). Despite the low overall efficiency (less than 0.1 %), this results represented in principle a quite remarkable step forward especially in comparison with the modest performance of one of the first *p*-DSCs that utilized a semiconducting photocathode with compact morphology^[332]. Anyhow, when compared to the analogous *n*-type counterparts, *p*-DSCs show their relative limits especially in terms of overall efficiencies since the most performing *p*-DSCs reach at most 1-2 % of overall conversion efficiency^[386, 387], whereas there is a constant upgrade of efficiency records with *n*-DSCs that reach values in the broad range of 10-15 %^[388, 4, 389, 390]. What mainly motivates the research on *p*-type photoelectrochemical devices is the realization of *t*-DSCs (Figure 7)^[267-271], which possess both electrodes with photoactive features due to their sensitization with colorants having complementary absorption properties^[237].

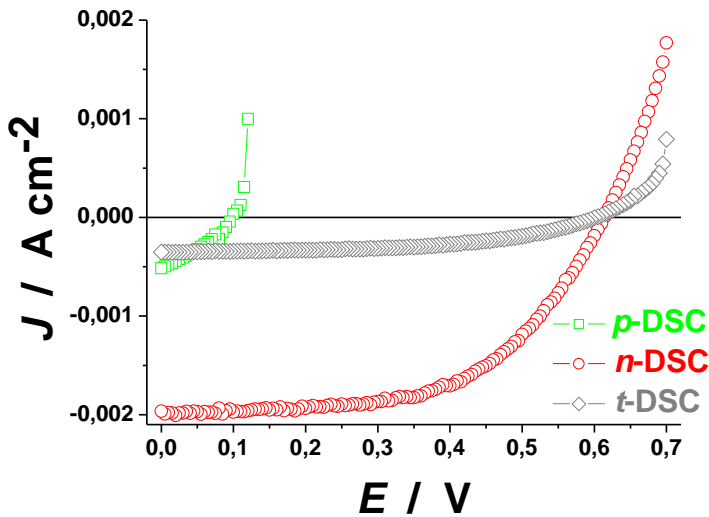


Figure 111. JV curves of a p -DSC with NiO cathode, of an n -DSC with TiO_2 anode and the corresponding t -DSC obtained upon combination of the p -DSC photocathode with the n -DSC photoanode. Anode and cathode sensitizers were N-719 and Fast Green, respectively. Redox shuttle: I_3^-/I^- . The illumination of the t -DSC with sun simulator having AM 1.5 was first directed to the side of the photoanode. Adapted from ref. 237.

Such a tandem configuration would allow the achievement of larger open circuit photovoltages with respect to the corresponding DSCs having single photoactive electrodes (Figure 11), but the gain in photopotential is obtained at expenses of current density and fill factor, that are both controlled by the less performing photoactive electrode (typically the nanoporous photocathode)^[237, 267-271]. The expected improvement of the overall photoelectrochemical performance in t -DSC^[391] in comparison to the DSC configurations with single photoactive electrode is possible only when n - and p -type devices generate separately photocurrents and reach FFs with comparable values the sensitizing agents of cathode and anode being complementary in terms of optical absorption (Figure 12). The present work updates the most recent achievements in the fields, which were not covered by the works of Daeneke *et al.* on the general limits of p -DSCs [241], and of Dini *et al.*^[239, 240] about the influence of photocathode nature and preparation method on the relative photoelectrochemical properties in the corresponding p -DSCs. Table 5 presents the most recent achievements obtained in the best performing p -DSCs when the redox shuttle was the couple I_3^-/I^- . The comparison of the p -DSCs regards PECs differing for the structure of the dye-sensitizer^[243, 250, 270, 271, 327, 329, 392-395] (Figure 12), and the nature of the nanostructured p -type photocathode^[396-400].

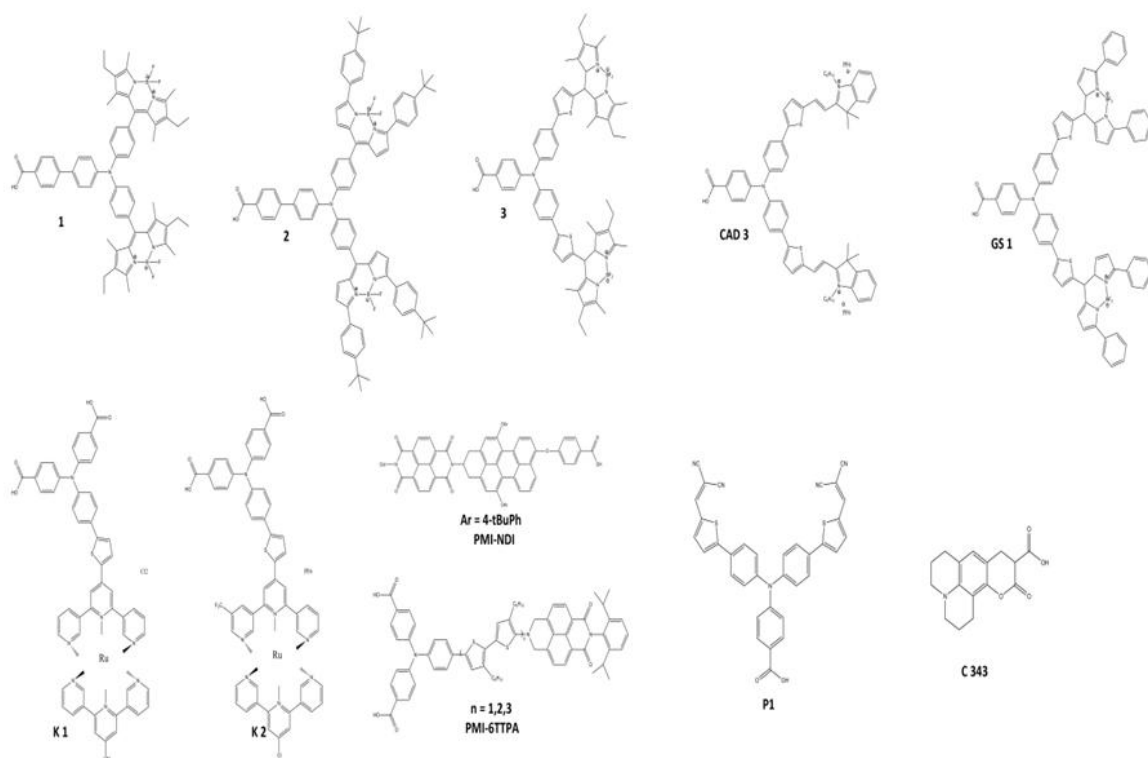


Figure 12. Structures of the most efficient dye-sensitizers insofar utilized in *p*-DSCs.

As recognized in precedence^[241], the main drawback of a *p*-DSC is the attainment of relatively low fill factors, which generally never exceed 40 % in the most efficient configurations (i.e. with $\eta > 0.14$, Table 5)^[243, 250, 270, 271, 327, 394] when the redox couple I_3^-/I^- is present. Upon adoption of a redox mediator other than the one based on iodide, the *p*-DSCs display performances with considerably ameliorated characteristics (Table 6)^[387, 401-411]. These facts would indicate that the main routes of photocathode optimization and refinement of dye-sensitizer design (Figure 13) must be necessarily accompanied by a systematic research on the selection and eventually the definition of new redox mediators in order to improve the performance of *p*-DSCs.

Cathode	Sensitizer	$\eta / \%$	V_{oc} / mV	$j_{sc} / mA\ cm^{-2}$	$FF/\%$	Ref.
NiO-RDS	P1	0.121	125	-2.84	33.7	[329]
NiO-CS	P1	0.110	128	-2.42	35.2	[329]
NiO sol-gel	P1	0.150	84	-5.48	34.0	[250]
NiO sol-gel	GS1	0.200	106	-5.87	31.0	[271]
NiO sol-gel	CAD1	0.250	101	-8.21	31.0	[271]
NiO sol-gel	K1	0.090	96	-2.91	32.0	[393]
NiO sol-gel	K2	0.070	93	-1.96	39.0	[393]
NiO sol-gel	1	0.060	97	-1.60	38.0	[243]
NiO sol-gel	2	0.140	109	-3.70	35.0	[243]
NiO sol-gel	3	0.050	95	-1.58	35.0	[243]
NiO sol-gel	3	0.080	79	-3.15	31.0	[392]
NiO ¹	PMI-6TTPA	0.460	208	-6.36	34.0	[328]
NiO ²	PMI-6TTPA	0.410	218	-5.35	35.0	[270]
NiO ³	PMI-6TTPA	0.400	292	-3.30	41.0	[394]
CuAlO ₂	PMI-6TTPA	0.040	333	-0.30	42.0	[396]
NiO sol-gel	PMI-NDI	0.073	120	-1.76	34.5	[395]
CuGaO ₂	PMI-NDI	0.023	187	-0.29	41.0	[397]
CuCrO ₂	C343	0.014	145	-0.24	39.8	[398]
K _x ZnO	C343	0.012	82	-0.41	35.8	[231]

RDS: rapid discharge sintering; CS: conventional sintering; 1: microballs; 2: nanoparticles; 3: nanorods

Table 5. *p*-DSC parameters for devices having different nanostructured photocathodes and dye-sensitizers (Figure 8). Redox shuttle: I₃⁻/I₂.

Cathode	Sensitizer	Redox couple	η / %	V_{oc} / mV	j_{sc} / mA cm ⁻²	FF/%	Ref.
NiO-SP	PMI-6TTPA	[Fe(acac) ₃] ^{0/1-}	2.51	645	-7.65	51.0	[387]
NiO-SP	PMI-6TTPA	[Co(en) ₃] ^{3+/2+}	1.20	724	-4.11	40.0	[387]
NiO-SP	PMI-6TTPA	I ₃ ⁻ /I ⁻	0.60	243	-6.26	39.0	[387]
NiO sol-gel	PMI-6TTPA	thiolate/disulfide	0.51	285	-5.30	34.0	[401]
NiO-SP ¹	PMI-6TTPA	[Co(en) ₃] ^{3+/2+}	1.30	709	-4.44	42.0	[386]
NiO-SP ²	PMI-6TTPA	[Co(en) ₃] ^{3+/2+}	1.30	660	-4.35	46.0	[386]
CuCrO ₂ ³	P1	thiolate/disulfide	0.17	309	-1.43	38.0	[402]
CuCrO ₂ ⁴	P1	thiolate/disulfide	0.17	240	-1.89	36.0	[402]
CuCrO ₂ ⁵	P1	thiolate/disulfide	0.22	293	-1.71	43.0	[402]
CuCrO ₂ ⁶	P1	thiolate/disulfide	0.23	304	-1.73	44.0	[402]
NiO-SP	P1	[Co(dtbbpy) ₃] ^{3+/2+}	0.099	280	-1.18	30.0	[403]
NiO-SP	E1	[Co(dtbbpy) ₃] ^{3+/2+}	0.130	320	-0.93	44.0	[403]
NiO-SP	E2	[Co(dtbbpy) ₃] ^{3+/2+}	0.102	320	-0.78	41.0	[403]
NiO tmp ⁷	T3	I ₃ ⁻ /I ⁻	0.180	144	-4.06	30.8	[404]
NiO tmp	T3H	I ₃ ⁻ /I ⁻	0.226	133	-5.56	30.5	[404]
NiO tmp	T4	I ₃ ⁻ /I ⁻	0.207	152	-3.94	34.5	[404]
NiO tmp	T4H	I ₃ ⁻ /I ⁻	0.317	152	-6.73	31.0	[404]
NiO np ⁸	PMI-6TTPA	I ₃ ⁻ /I ⁻	0.45	226	-5.40	36.0	[405]
NiO np	PMI-6TTPA	[Co(en) ₃] ^{3+/2+}	1.00	640	-3.72	42.0	[405]
CuCrO ₂	PMI-6TTPA	I ₃ ⁻ /I ⁻	0.23	268	-1.98	44.0	[405]
CuCrO ₂	PMI-6TTPA	[Co(en) ₃] ^{3+/2+}	0.48	734	-1.23	53.0	[405]
NiO db ⁹	CW1	I ₃ ⁻ /I ⁻	0.114	99	-2.66	35.0	[406]
NiO db	CW2	I ₃ ⁻ /I ⁻	0.160	93	-4.05	34.0	[406]
NiO db	P1	I ₃ ⁻ /I ⁻	0.143	93	-4.67	33.0	[406]
NiO	BH2	I ₃ ⁻ /I ⁻	0.13	97	-4.3	31.0	[407]
NiO	BH4	I ₃ ⁻ /I ⁻	0.28	128	-7.4	30.0	[407]
NiO	BH6	I ₃ ⁻ /I ⁻	0.13	95	-4.4	31.0	[407]
NiO db	QT1	I ₃ ⁻ /I ⁻	0.33	120	-8.2	34.0	[408]
NiO db	QT1	Co ^{3+/2+}	0.50	226	-6.5	34.0	[408]
CuCrO ₂	P1	thiolate/disulfide	0.17	309	-1.43	38.0	[409]
CuCrO ₂ /Au	P1	thiolate/disulfide	0.31	305	-2.68	38.0	[409]
NiO cl+mp ¹⁰	P1	PCBM	-	620	-0.05	-	[410]
NiO-SP	DPP-1	[Co(dtbbpy) ₃] ^{3+/2+}	0.205	330	-2.06	30.0	[411]
NiO-SP	DPP-2	[Co(dtbbpy) ₃] ^{3+/2+}	0.21	370	-1.95	29.0	[411]
NiO-SP	DPP-NDI	[Co(dtbbpy) ₃] ^{3+/2+}	0.13	292	-1.56	29.0	[411]

SP: screen printed; 1: 0.1 M LiTFSI; 2: 0.5 M LiTFSI; 3: 1.8 μ m thick; 4: 3 μ m thick; 5: PEDOT counter electrode; 6: CoS counter electrode; 7: tmp = template [93]; 8: np = nanoparticles; 9: db = doctor blade; 10: cl+mp = mesoporous layer over a compact layer

Table 6. Parameters of the *p*-DSC differing for the nature of the nanostructured photocathodes and dye-sensitizers (Figure 9). The redox shuttles here considered are alternatives to the reference couple I₃⁻/I⁻ here considered for comparative purposes.

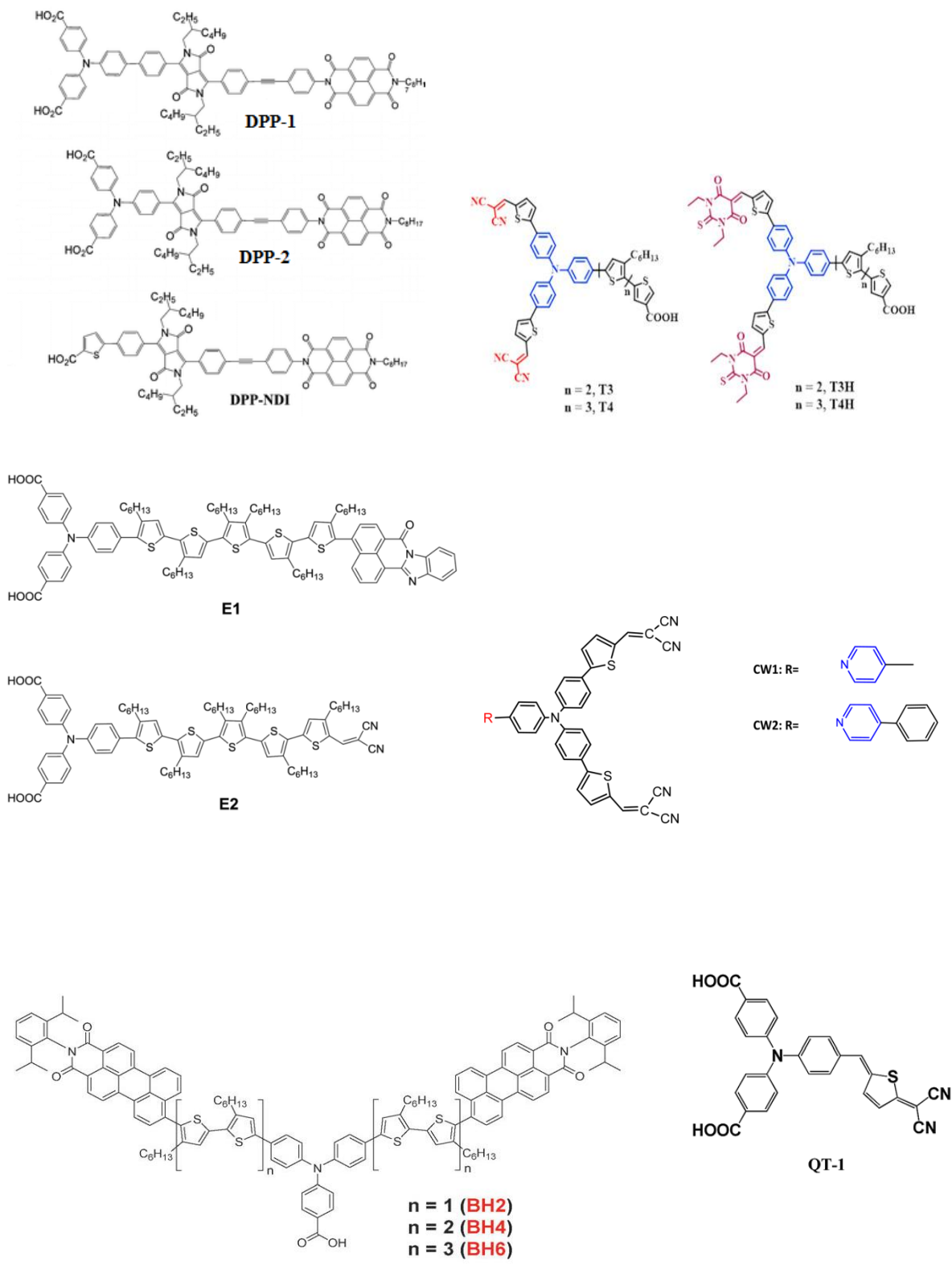


Figure 13. Structures of newly synthesized dye-sensitizers for NiO based p -DSCs with $\text{Co}^{2+/3+}$ and I_3^-/I^- redox shuttles.

Closely related to the progress of *p*-DSC is the one of *t*-DSC with both electrodes nanostructured and opportunely sensitized^[237, 267-271, 401, 412]. In terms of performance, the *t*-DSC represents a device still in its infancy, which strongly needs of a further stage of evolution for reaching and eventually surpassing the performances of the corresponding *n*-DSCs in accordance to the predictions of thermodynamics^[391]. Table 7 presents the list of the parameters characterizing the performance of the *t*-DSCs reported so far.

Cathode	<i>p</i> -sensitizer	<i>n</i> -sensitizer	$\eta / \%$	V_{OC} / mV	$j_{sc} / mA\ cm^{-2}$	$FF / \%$	Ref.
NiO	PMI-6T-TPA	SQ2	1.19	924	1.9	67	[401]
NiO ^a	PMI-6T-TPA	SQ2	1.33	814	2.5	65	[401]
NiO	CAD 3	D35	1.7	613	5.15	54	[271]
NiO	GS 1	D35	1.3	638	4.54	43	[271]
NiO	P1	D35	1.1	732	3.71	38	[271]
NiO	3	N719	1.91	1079	2.40	74	[270]
NiO ^b	PMI-NDI	N719	0.55	910	0.97	62	[270]
Se		N719	0.98	940	2.72	39	[412]
NiO	MC ^c	N3 ^d	0.66	918	3.62	19	[268]
NiO	ERY B ^e	N719	0.39	732	2.15		[267]

a: cell utilizing thiolate/disulfide redox mediator; b: cell employing Co^{3+/2+} as redox shuttle; c: MC = 3-carboxymethyl-5-[2-(3-octadecyl-2-benzothiazolinyldene) ethylidene]-2-thioxo-4-thiazolidine; d: N3 = cis-di(thiocyanato)bis(4,4'-dicarboxy-2,2'-bipyridine) ruthenium(II); e: ERY B = erythrosine b

Table 7. Parameters of the *t*-DSC differing for the combination of the nanostructured electrodes and corresponding dye-sensitizers. The redox shuttle is represented by the couple I₃⁻/I⁻ if not otherwise indicated. The photoanodic material is mesoporous TiO₂. The *t*-DSC is illuminated from the side of the *n*-electrode (TiO₂ photoanode).

Conclusions

The use of *n*-type semiconductors as photoanodes and *p*-type semiconductors as photoactive cathodes of *p*- and *t*-DSCs has been reviewed. Concerning the *n*-type semiconductor, a substantial performance improvement has been obtained by optimizing the dimensions of the particles that can be achieved by tailoring the synthesis conditions. The role of band gap, morphology, composition and doping of *n*-type semiconductors for the development of efficient photoanodes has been reported. Several contributions were summarized here with the description of nanostructures (0D, 1D, 2D, 3D and mesoporous nanostructures) and of the effect of doping of photoanodes materials (mostly TiO₂ and ZnO) on the performances of DSSCs. The maximum photoconversion efficiency (13-15%) has been reported using titania based photoanodes. About photocathodes, we have posed our attention towards the most recent developments on *p*-type semiconductors having

nanostructured features and mesoporous morphology. Such features impart large surface areas and allow the anchoring of a large number of dye-sensitizers per unit area with favourable consequences on the efficaciousness with which the desired photoactivated electrochemical process occurs. An important issue related to the employment of nanostructured semiconducting cathodes is the determination of the position of the energy levels for the frontier states directly involved in the exchange of photocarriers with the dye-sensitizer in the excited state. Among nanostructured semiconducting cathodes, nickel oxide having non-stoichiometric features is the most important representative of *p*-type photocathodes due to its intrinsic photoelectroactivity. Other examples of nanostructured photocathodes have been also mentioned in the present review thus indicating new possible directions of research in the materials science of semiconducting electrodes of *p*-type. Also the most important examples of dye-sensitizers for *p*-type nanostructured electrodes have been reviewed. The aspects of electronic conjugation, presence of spacers and substituents with peculiar electronic effects, state of surface immobilization have been briefly considered in the analysis of the most important structural factors of the dye-sensitizers, which control the process of charge photogeneration and injection in the photocathode. An aspect that has not been considered in depth is the feasibility of the synthesis and purification methods for purposely designed photo-/electrocatalytic systems for *p*-type semiconductors. In terms of short-term goals, the main purpose of the development of increasingly efficient *p*-DSCs is the achievement of efficiencies comparable to those of the *n*-type counterparts in order to produce tandem devices with efficiencies approaching the theoretical limit.

Author Contributions

Part I has been written and compiled by A.L., C.C. and F.D.P. who contributed equally to the completion of this first part of the review. Part II has been written and compiled by D.D. and M.B. who contributed equally to the completion of this second part of the review.

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