



Isoperimetric Inequalities for the Magnetic Neumann and Steklov Problems with Aharonov–Bohm Magnetic Potential

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Abstract

We discuss isoperimetric inequalities for the magnetic Laplacian on bounded domains of \mathbb{R}^2 endowed with an Aharonov–Bohm potential. When the flux of the potential around the pole is not an integer, the lowest eigenvalue for the Neumann and the Steklov problems is positive. We establish isoperimetric inequalities for the lowest eigenvalue in the spirit of the classical inequalities of Szegő–Weinberger, Brock and Weinstock, the model domain being a disk with the pole at its center. We consider more generally domains in the plane endowed with a rotationally invariant metric, which include the spherical and the hyperbolic case.

Keywords Magnetic Laplacian · Aharonov–Bohm magnetic potential · Ground state · Neumann problem · Steklov problem · Reverse Faber–Krahn inequality

Mathematics Subject Classification 35J10 · 35P15 · 49Rxx · 58J50 · 81Q10

1 Introduction

The question of the *isoperimetric inequalities* for the eigenvalues of the Laplacian (in particular for the first nonzero eigenvalue) is a long standing problem. Let us give a short and partial summary in the case of bounded domains of the Euclidean space

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which will be the main topic of the present paper. It began with the celebrated Faber–Krahn inequality [17]: for Dirichlet boundary conditions, among all bounded domains of given volume, the *first* eigenvalue is minimized by the ball. For spaces of constant curvature the result can be found in [7]. For Neumann boundary conditions, among all bounded domains of given volume with Lipschitz boundary, the *second* eigenvalue (i.e., the first nonzero) is maximized by the ball. This is the Szegő–Weinberger inequality [23, 24]. This inequality has been extended to bounded domains in spaces of constant curvature by Ashbaugh and Benguria [3] (see also [4]). For Robin boundary conditions with positive parameter, the ball also realizes the minimum [8]. For other operators, similar results exist. For the Steklov problem, the *second* eigenvalue (the first nonzero) is maximized by the ball among the bounded open domains of given volume with Lipschitz boundary: this is the inequality of Brock [5]. However, if we consider the domains of \mathbb{R}^2 with boundary of given length, the second eigenvalue is maximized by the disk only among all simply connected domains. This is the inequality of Weinstock [25]. There exist annuli with larger second eigenvalue. Again, we refer to [4, 7] for more discussion and generalizations. Note that, even if we will not go in this direction, the maximization or minimization of higher eigenvalues is intensively studied, see for example [6] for the second nonzero eigenvalue of the Neumann problem and [14] for the third eigenvalue of the Robin problem.

In this paper we will be mainly concerned with the Neumann problem for the Aharonov–Bohm magnetic Laplacian on domains of \mathbb{R}^2 (see (1)–(3)) and with the corresponding Steklov problem (see (1)–(4)).

The problems under consideration are instances of how the Aharonov–Bohm effect has an influence on the bound states of a quantum particle. Roughly speaking, the Aharonov–Bohm effect can be described as follows: consider an impenetrable region (typically, an ideal solenoid) where a magnetic field is confined, while a charged quantum particle is placed outside the impenetrable region (i.e., the wave function describing the particle vanishes near the boundary of the solenoid). It turns out that the Hamiltonian of the particle is influenced by a vector potential A which generates the magnetic field, even if the magnetic field vanishes outside the solenoid. This effect has been predicted by Ehrenberg and Siday [10], and discussed analytically ten years later by Aharonov and Bohm [1], and has been validated then by a series of experiments, especially in the framework of the scattering of the electrons from a magnetic solenoid. The Aharonov–Bohm effects illustrates the physicality of magnetic potentials (energies) in quantum mechanics, which were classically considered just mathematical artifices for computational purposes, contrarily to fields (forces), the core of Newtonian physics. Other theoretical consequences concerning local/global nature of electromagnetic effects can be deduced. Nevertheless, there have been many disputes on the effective existence of the Aharonov–Bohm effect, and it has been argued that experiments (mainly involving scattering situations) do not really confirm the theory. On the other hand, if we consider the Aharonov–Bohm effect on bound states, which is the main topic of our paper from a mathematical point of view, there are no more ambiguities, as discussed in [15, 22]. The mathematical problems which we study, namely problem (3)–(4) with potential (1), describe the motion of a quantum particle in a region Ω of a plane perpendicular to a magnetic solenoid of arbitrarily small radius R (in experiments, smaller than de Broglie wavelength of the electron).

It turns out that the eigenvalues of the corresponding Hamiltonian feel in some sense a shift which is related to the flux of the vector potential along closed paths (modulo integers), as described in [22]. In particular, for Neumann boundary conditions on the boundary of Ω , the ground state needs not to be zero, and this poses fundamental mathematical questions concerning isoperimetric inequalities and eigenvalue bounds.

In the case of the Neumann problem (3), we will also consider the Aharonov–Bohm magnetic Laplacian on domains of surfaces of revolution (in particular, the standard sphere \mathbb{S}^2 and the standard hyperbolic space \mathbb{H}^2). Most of the time, the first eigenvalue of this kind of problem is strictly positive and its study difficult. For example, for the magnetic Laplacian with constant non zero magnetic field and Dirichlet boundary condition in \mathbb{R}^2 , it is known that the first eigenvalue is minimized by the disk among all domains of given area: this was shown in [11] and the proof is quite involved. To our knowledge, a similar result is not known in \mathbb{S}^2 or \mathbb{H}^2 . However, for the magnetic Laplacian with constant magnetic field and magnetic Neumann boundary condition, it is no longer true that the disk maximizes the first eigenvalue: even for simply connected domains, the question is open, see [12, Question 1, Remark 2.4 and Proposition 3.3]. More information can be found also in [13, §4 and §5]. Still for the case of constant magnetic field, we mention [19, 20] for bounds on Dirichlet and Neumann eigenvalues of certain families of domains, and [16] for an isoperimetric inequality for the Robin problem.

In the case of \mathbb{R}^2 , we show that, among all domains of given area, the disk with the singularity of the magnetic field at the center is the unique maximizer of the *first* eigenvalue, which is positive provided that the flux is not an integer. This is a reverse Faber–Krahn inequality, that we obtain in the spirit of Szegő–Weinberger [23, 24]. For the Steklov problem we prove two isoperimetric inequalities for the *first* eigenvalue, which, again, is positive if the flux is not an integer. These correspond to the inequalities of Weinstock and of Brock [5, 25].

For the Neumann problem, we obtain similar results for domain of \mathbb{S}^2 and \mathbb{H}^2 . The result will be a consequence of a general isoperimetric inequality for the Schrödinger operator on a manifold of revolution with radial, non-negative, and radially decreasing potential, that we will prove here (see Theorem 6).

We finally remark that the Faber–Krahn inequality for the magnetic Dirichlet problem with Aharonov–Bohm potential is trivial: the first eigenvalue is minimized by that of the usual Laplacian on the disk, among all bounded domains of given area. A simple argument to see this comes from the diamagnetic inequality [21, Theorem 7.21], namely $|\nabla|u|| \leq |\nabla^A u|$, where u is any smooth complex-valued function in the magnetic Sobolev space associated with the magnetic Dirichlet problem with Aharonov–Bohm potential A , and $\nabla^A u$ is the magnetic gradient of u (see “Appendix A” for precise definitions). This pointwise identity implies that the first Dirichlet magnetic eigenvalue on any domain is lower bounded by the first eigenvalue of the Dirichlet Laplacian on that domain. Therefore, by the standard Faber–Krahn inequality, it is lower bounded by the first eigenvalue of the Dirichlet Laplacian on a disk with the same area.

2 Notation and Statement of Results

Let Ω be a smooth bounded domain of \mathbb{R}^2 with a distinguished point $x_0 = (a, b)$ and consider the one-form

$$A_0 = -\frac{x_2 - b}{(x_1 - a)^2 + (x_2 - b)^2} dx_1 + \frac{x_1 - a}{(x_1 - a)^2 + (x_2 - b)^2} dx_2. \tag{1}$$

The one-form $A_{x_0, \nu} = \nu A_0$ will be called *Aharonov–Bohm potential with pole x_0 and flux ν* . Note that A_0 is smooth, closed, co-closed (hence harmonic) on $\mathbb{R}^2 \setminus \{x_0\}$, and is singular at x_0 ; it gives rise to a zero magnetic field ($B = dA_{x_0, \nu} = 0$). Recall that the flux of A around a loop $c : [0, L] \rightarrow \mathbb{R}^2$ is

$$\frac{1}{2\pi} \oint_c A = \frac{1}{2\pi} \int_0^L A(c'(t)) dt \tag{2}$$

where we assume that c is travelled once in the counterclockwise direction (even though the direction or the number of turns have no effect on our results). We also say that a closed form A has flux ν around x_0 if the flux of A around any loop enclosing x_0 is ν : since A closed, this definition does not depend on the loop c . In particular, let Γ be the outer boundary of Ω ; if x_0 is inside Γ , then ν is the flux of $A_{x_0, \nu}$ around Γ . Let $\Delta_{A_{x_0, \nu}}$ be the magnetic Laplacian with potential $A_{x_0, \nu}$: it is the operator

$$\Delta_{A_{x_0, \nu}} u = \Delta u + |A_{x_0, \nu}|^2 u + 2i \langle \nabla u, A_{x_0, \nu} \rangle$$

acting on complex valued functions u (the sign convention is that $\Delta u = -\sum_j \partial_{x_j}^2 u$). Of course, we can always assume that x_0 is the origin.

We will also consider the case when the ambient space is a two-dimensional manifold of revolution (M, g) with pole x_0 and polar coordinates (r, t) , where r is the distance to x_0 . In this case, we consider the form $A_0 = dt$, which is harmonic (closed and co-closed), and with flux 1 around x_0 .

In this paper, we consider the eigenvalue problem for $\Delta_{A_{x_0, \nu}}$ with magnetic Neumann conditions:

$$\begin{cases} \Delta_{A_{x_0, \nu}} u = \lambda u, & \text{in } \Omega, \\ \langle \nabla u - iuA_{x_0, \nu}, N \rangle = 0, & \text{on } \partial\Omega, \end{cases} \tag{3}$$

and also the magnetic Steklov eigenvalue problem:

$$\begin{cases} \Delta_{A_{x_0, \nu}} u = 0, & \text{in } \Omega, \\ \langle \nabla u - iuA_{x_0, \nu}, N \rangle = \sigma u, & \text{on } \partial\Omega. \end{cases} \tag{4}$$

Here N is the outer unit normal to $\partial\Omega$. With abuse of notation we still denote by $A_{x_0, \nu}$ the potential dual to the 1-form $A_{x_0, \nu}$. We also denote by $\nabla^{A_{x_0, \nu}} u$ the vector

field

$$\nabla^{A_{x_0,v}} u = \nabla u - iuA_{x_0,v}$$

which is called *magnetic gradient*. Therefore, the magnetic Neumann condition reads $\langle \nabla^{A_{x_0,v}} u, N \rangle = 0$, while the magnetic Steklov condition is $\langle \nabla^{A_{x_0,v}} u, N \rangle = \sigma u$.

We will prove in “Appendix A” that each of these two problems admits an infinite discrete sequence of eigenvalues of finite multiplicity. We will denote by $\lambda_1(\Omega, A_{x_0,v})$ the first eigenvalue of Problem (3) and by $\sigma_1(\Omega, A_{x_0,v})$ the first eigenvalue of Problem (4). These two eigenvalues are non-negative for all $v \in \mathbb{R}$ and are strictly positive if and only if $v \notin \mathbb{Z}$ (in particular, when $x_0 \in \Omega^c$ and the flux of $A_{x_0,v}$ is 0 in Ω , we set $v = 0$); in particular, when $v \in \mathbb{Z}$, the two spectra of problems (3) and (4) reduce to the corresponding spectra of the Laplacian Δ (i.e., when $A_{x_0,v} = 0$, see “Appendix A”).

In the sequel, we will often suppose that $v \notin \mathbb{Z}$, so $\lambda_1(\Omega, A_{x_0,v})$ and $\sigma_1(\Omega, A_{x_0,v})$ are both positive.

The first result is a reverse Faber–Krahn inequality for the first eigenvalue of the Neumann problem whose proof is based on the well-known Szegő–Weinberger approach [23, 24].

Through all the paper, by $|\Omega|$ we denote the Lebesgue measure of a smooth bounded domain Ω , and by $|\partial\Omega|$ the length of its boundary.

Theorem 1 *Let Ω be a smooth bounded domain in \mathbb{R}^2 or \mathbb{H}^2 , and let $A_{x_0,v}$ be the Aharonov–Bohm potential with pole at x_0 and flux v . Let $B(x_0, R)$ be the disk with center x_0 and radius R such that $|B(x_0, R)| = |\Omega|$. Then*

$$\lambda_1(\Omega, A_{x_0,v}) \leq \lambda_1(B(x_0, R), A_{x_0,v}); \tag{5}$$

if $v \notin \mathbb{Z}$, equality holds if and only if $\Omega = B(x_0, R)$.

We will observe that Theorem 1 extends to domains in the manifold (\mathbb{R}^2, g) where g is a complete, non-positively curved, rotationally invariant metric around x_0 , the pole of the magnetic potential (see Sect. 4.1).

This theorem will be a consequence of a more general result about an isoperimetric inequality for Schrödinger operators on revolution manifolds with pole x_0 and radial potential V that we will present in Sect. 3. In fact, Theorem 1 holds also in this setting, under suitable hypothesis on the function describing the density of the Riemannian metric in standard polar coordinates.

The case of the sphere \mathbb{S}^2 is more involved. We are able to show a similar result to Theorem 1 only if the domain is contained in a hemisphere centered at the pole x_0 .

Theorem 2 *Let Ω be a smooth domain contained in a hemisphere centered at x_0 , and let $A_{x_0,v}$ be the Aharonov–Bohm potential with pole at x_0 and flux v . Let $B(x_0, R)$ be the disk in \mathbb{S}^2 with center x_0 and radius R such that $|B(x_0, R)| = |\Omega|$. Then*

$$\lambda_1(\Omega, A_{x_0,v}) \leq \lambda_1(B(x_0, R), A_{x_0,v}); \tag{6}$$

if $\nu \notin \mathbb{Z}$, equality holds if and only if $\Omega = B(x_0, R)$.

Note that the analogous result for the second eigenvalue of the Neumann Laplacian is proved in [3]. However, for simply connected domains we can do better.

Theorem 3 *Let Ω be a smooth simply connected domain in \mathbb{S}^2 with $|\Omega| \leq 2\pi$ and $-x_0 \notin \Omega$, and let $A_{x_0,\nu}$ be the Aharonov–Bohm potential with pole at x_0 and flux ν . Let $B(x_0, R)$ be the disk in \mathbb{S}^2 with center x_0 and radius R such that $|B(x_0, R)| = |\Omega|$. Then*

$$\lambda_1(\Omega, A_{x_0,\nu}) \leq \lambda_1(B(x_0, R), A_{x_0,\nu}); \tag{7}$$

if $\nu \notin \mathbb{Z}$, equality holds if and only if $\Omega = B(x_0, R)$.

The next result is the analogous of Brock’s inequality [5] for the first Steklov eigenvalue on planar domains:

Theorem 4 *Let Ω be a smooth bounded domain in \mathbb{R}^2 and let $A_{x_0,\nu}$ be the Aharonov–Bohm potential with pole at x_0 and flux ν . Let $B(x_0, R)$ be the disk with center x_0 and radius R such that $|B(x_0, R)| = |\Omega|$. Then*

$$\sigma_1(\Omega, A_{x_0,\nu}) \leq \sigma_1(B(x_0, R), A_{x_0,\nu}) = \frac{\sqrt{\pi}}{\sqrt{|\Omega|}} \inf_{k \in \mathbb{Z}} |\nu - k|. \tag{8}$$

If $\nu \notin \mathbb{Z}$, equality holds if and only if $\Omega = B(x_0, R)$.

Finally, we prove the analogue of Weinstock’s inequality [25]:

Theorem 5 *Let Ω be a smooth bounded and simply connected domain in \mathbb{R}^2 and let $A_{x_0,\nu}$ be the Aharonov–Bohm potential with pole at x_0 and flux ν . Let $B(x_0, R)$ be the disk with center x_0 and radius R such that $|\partial B(x_0, R)| = |\partial\Omega|$. Then*

$$\sigma_1(\Omega, A_{x_0,\nu}) \leq \sigma_1(B(x_0, R), A_{x_0,\nu}) = \frac{2\pi}{|\partial\Omega|} \inf_{k \in \mathbb{Z}} |\nu - k|. \tag{9}$$

If $\nu \notin \mathbb{Z}$, equality holds if and only if $\Omega = B(x_0, R)$.

Note that the upper bounds of Theorems 4 and 5 correctly reduce to zero whenever the flux is an integer.

We stated Theorem 5 for planar, simply connected domains, however it extends to any Riemannian surface with boundary.

We conclude this section with a few remarks. It is natural to ask what happens for the second eigenvalue of (3) and (4), at least on planar domains. One immediately observes that Theorems 1 and 4 no longer hold, in the sense that the ball punctured at the origin is not a maximiser. In fact, $\lambda_2(B(x_0, R), A_{x_0,\nu}) = \frac{(z'_{1-\inf_{k \in \mathbb{Z}} |\nu - k|, 1})^2}{R^2} < \frac{(z'_{1,1})^2}{R^2}$ when $\nu \notin \mathbb{Z}$. Here $z'_{\mu,1}$ denotes the first positive zero of the derivative of the Bessel function J_μ (see ‘‘Appendix B.3’’). We recall that $\frac{(z'_{1,1})^2}{R^2}$ is exactly the second

Neumann eigenvalue of the Laplacian on a ball of radius R . Analogously, we have $\sigma_2(B(x_0, R), A_{x_0, v}) = \frac{1 - \inf_{k \in \mathbb{Z}} |v - k|}{R} < \frac{1}{R}$ (see “Appendix B.4”), and $\frac{1}{R}$ is the second Steklov eigenvalue of the Laplacian on $B(x_0, R)$. However, for problem (3) (4), it can be shown that the disjoint union of two balls with suitable radii, and one of them centered at the pole, and total area π , has second eigenvalue strictly greater than that of the standard Neumann (Steklov) eigenvalue on $B(0, 1)$. Therefore we are left with the following

Open problem 1. Find (if it exists) a maximiser for the second Neumann (Steklov) Aharonov–Bohm eigenvalue among all smooth bounded domains in \mathbb{R}^2 .

As for inequality (9), preliminary calculations show that it holds for all circular annuli in \mathbb{R}^2 , but it fails in the case of long cylinders. In fact, when $\Omega = \mathbb{S}^1 \times (-L, L)$, the first Steklov eigenvalue is given by $\inf_{k \in \mathbb{Z}} |v - k| \tanh(\inf_{k \in \mathbb{Z}} |v - k|L)$, hence (9) does not hold for $L > L_0$, with L_0 sufficiently large. We are left with the following

Open problem 2. Does inequality (9) hold for all doubly connected domains of the plane?

The present paper is organized as follows. In Sect. 3 we prove an isoperimetric inequality for the first (positive) Neumann eigenvalue of the Schrödinger operator $\Delta + V$ on domains in manifolds of revolution, under suitable hypothesis on the potential V and on the density of the Riemannian metric (Theorem 6). In Sect. 4, Theorem 6 is applied to the magnetic Neumann spectrum. In particular, in Sect. 4.1 the reverse Faber–Krahn inequality is proved for manifolds of revolution (Theorem 9). As a consequence, we prove that it holds for domains in \mathbb{R}^2 and \mathbb{H}^2 , (Corollary 10). In Sect. 4.2 it is proved for spherical domains contained in a hemisphere centered at the pole (Theorem 12). In Sect. 4.3 we prove the isoperimetric inequality for spherical simply connected domains with area less than 2π (Theorem 14). In Sect. 5 we prove Brock’s inequality for planar domains (Theorem 15), and Weinstock’s inequality for planar domains (Theorem 16).

We have included in this article a quite complete set of appendices, where we discuss the functional and geometrical setting for the magnetic problems that we consider. In particular, we will compute explicitly the Neumann and Steklov spectrum for the unit disk.

In “Appendix A” we provide the basic spectral theory for problems (3) and (4). “Appendix B” contains a more explicit description of the eigenvalues and the eigenfunctions of the magnetic Neumann and Steklov problems on disks in manifolds of revolution (see “Appendices B.1” and B.2). These facts, which have an interest on their own, are crucial for the proofs of the main Theorems. In “Appendices B.3” and B.4 we describe the eigenfunctions and eigenvalues on disks in \mathbb{R}^2 . Finally, in “Appendix C” we prove the conformal invariance of the Aharonov–Bohm energy which is crucial in the proof of Weinstock’s inequality.

3 Isoperimetric inequality for Schrödinger operators

In this section, Ω will be a bounded smooth domain in a n -dimensional manifold of revolution (M, g) with pole x_0 . With D we denote the diameter of M (which can be infinite) and with D_Ω we denote the diameter of Ω .

We recall that a smooth n -dimensional Riemannian manifold (M, g) with a distinguished point x_0 is called a *revolution manifold with pole x_0* if $M \setminus \{x_0\}$ is isometric to $(0, D] \times \mathbb{S}^{n-1}$ whose metric is, in normal coordinates based at the pole, $g = dr^2 + \Theta(r)^2 g_{\mathbb{S}^{n-1}}$, for $r \in (0, D)$. Here $\Theta(0) = \Theta'(0) = 0$, $\Theta''(0) = 1$, and $g_{\mathbb{S}^{n-1}}$ is the standard metric on the $n - 1$ -dimensional sphere. The density of the Riemannian metric on M in normal coordinates is given by $\sqrt{\det g} = \Theta^{n-1}(r) = \theta(r)$.

It is known that, for space forms of constant curvature $K = 0, -1, 1$ we have:

$$\theta(r) = \begin{cases} r^{n-1} & \text{if } K = 0, \\ \sinh^{n-1}(r) & \text{if } K = -1, \\ \sin^{n-1}(r) & \text{if } K = 1. \end{cases}$$

In general, we have $\theta > 0$ on $(0, D)$. We refer to ‘‘Appendix B’’ for more information on manifolds of revolution.

We discuss here an isoperimetric inequality for the first eigenvalue of the Schrödinger operator:

$$\begin{cases} \Delta u + Vu = \lambda u, & \text{in } \Omega, \\ \langle \nabla u, N \rangle = 0, & \text{on } \partial\Omega. \end{cases} \tag{10}$$

Note that the results of this section can be applied to manifolds of revolution of any dimension $n \geq 2$.

Assumptions on the potential V .

1. the potential V is smooth on $M \setminus \{x_0\}$, non-negative and radial with respect to x_0 , that is, $V = V(r)$;
2. V is non-increasing on $(0, D_\Omega) : V'(r) \leq 0$ on $(0, D_\Omega)$;
3. $\theta'V' + 2V^2\theta \leq 0$ on $(0, R)$, where $R > 0$ is such that $|B(x_0, R)| = |\Omega|$;
4. there exists a first eigenfunction u of (10) on $B(x_0, R)$ which is non-negative, radial and non-decreasing in the radial direction: $u' \geq 0$.

We consider the following number:

$$\lambda_1(\Omega, \Delta + V) = \inf_{0 \neq u \in H_V^1(\Omega)} \frac{\int_\Omega |\nabla u|^2 + Vu^2}{\int_\Omega u^2}, \tag{11}$$

where $H_V^1(\Omega) = \{u \in H^1(\Omega) : V^{1/2}u \in L^2(\Omega)\}$, and $H^1(\Omega)$ is the standard Sobolev space of square integrable functions with square integrable weak first derivatives. Since V is non-negative, the infimum in (11) exists and is non-negative. We are ready to state the main result of this section.

Theorem 6 *Let Ω be a smooth bounded domain in a manifold of revolution M with pole at x_0 . Let $B = B(x_0, R)$ be the ball centered at x_0 with the same volume of Ω . Let Assumptions 1–4 hold. Then*

$$\lambda_1(\Omega, \Delta + V) \leq \lambda_1(B(x_0, R), \Delta + V).$$

Equality holds if and only if $\Omega = B(x_0, R)$.

If the spectrum of (10) is discrete in its lower portion, the number $\lambda_1(\Omega, \Delta + V)$ is the first eigenvalue. This is the case of regular potentials (e.g., $V \in L^{n/2}$ for $n \geq 3$ or $V \in L^{1+\delta}$, $\delta > 0$ for $n = 2$), but also of singular potentials of the form $\frac{\nu^2}{r^2}$ (inverse-square potentials). In both these cases, the whole spectrum is purely discrete and made of non-negative eigenvalues of finite multiplicity diverging to $+\infty$.

Let now $B(x_0, R)$ be the ball of radius R centered at the pole x_0 and assume that there exists a first eigenfunction of (10) on $B(x_0, R)$ which is non-negative (and therefore radial, as V is radial) and non-decreasing with respect to r . Let us denote this function by $u = u(r)$. It satisfies

$$\begin{cases} u'' + \frac{\theta'}{\theta}u' + (\lambda - V)u = 0, & \text{in } (0, R), \\ u'(R) = 0, \end{cases} \tag{12}$$

where $\lambda = \lambda_1(B(x_0, R), \Delta + V)$ is the first eigenvalue.

In order to prove Theorem 6 we need the following lemma:

Lemma 7 *Let $u = u(r)$ be a solution of (12) such that $u \geq 0$ and $u' \geq 0$ on $(0, R)$. Let*

$$F(r) = u'(r)^2 + V(r)u(r)^2.$$

If $V' \leq 0$ on $(0, R)$ and $\theta'V' + 2V^2\theta \leq 0$ on $(0, R)$, then one has:

$$F'(r) \leq 0$$

on $(0, R)$.

Proof One has:

$$\begin{aligned} F' &= 2u'u'' + V'u^2 + 2Vuu' = 2u' \left(-\frac{\theta'}{\theta}u' - (\lambda - V)u \right) + V'u^2 + 2Vuu' \\ &= -2\frac{\theta'}{\theta}u'^2 - 2\lambda uu' + V'u^2 + 4Vuu' \leq -2\frac{\theta'}{\theta}u'^2 + V'u^2 + 4Vuu' \end{aligned}$$

because $u \geq 0$ and $u' \geq 0$. Now:

$$V'u^2 + 4Vuu' = V' \left(u + 2\frac{V}{V'}u' \right)^2 - 4\frac{V^2}{V'}u'^2,$$

and we have:

$$F' \leq -2\left(\frac{\theta'}{\theta} + 2\frac{V^2}{V'}\right)u'^2 + V'\left(u + 2\frac{V}{V'}u'\right)^2.$$

As $V' \leq 0$ we conclude:

$$F' \leq -2\left(\frac{\theta'}{\theta} + 2\frac{V^2}{V'}\right)u'^2.$$

If $\theta'V' + 2V^2\theta \leq 0$ then, dividing by $\theta V'$ (which is non-positive) we indeed have

$$\frac{\theta'}{\theta} + 2\frac{V^2}{V'} \geq 0$$

which guarantees that $F' \leq 0$. □

Proof (Proof of Theorem 6) Define the radial function $f : M \rightarrow \mathbb{R}$ as follows:

$$f(r) = \begin{cases} u(r) & \text{for } r \leq R, \\ u(R) & \text{for } r \geq R. \end{cases}$$

We note that, by construction, $f|_{\Omega} \in H^1_V(\Omega)$, therefore it is possible to use it as test function in (11).

We start by observing that, by assumption $|\Omega \cap B^c| = |\Omega^c \cap B|$, so that, since u is increasing, we have $u(r) \leq u(R)$ and

$$\int_{\Omega} f^2 \geq \int_B u^2.$$

In fact:

$$\begin{aligned} \int_{\Omega} f^2 &= \int_{\Omega \cap B} f^2 + \int_{\Omega \cap B^c} f^2 = \int_{\Omega \cap B} u^2 + u(R)^2|\Omega \cap B^c| \\ &= \int_{\Omega \cap B} u^2 + u(R)^2|\Omega^c \cap B| \geq \int_{\Omega \cap B} u^2 + \int_{\Omega^c \cap B} u^2 = \int_{\Omega} u^2. \end{aligned} \tag{13}$$

We have to control the energy. Since $F(r) = u^2(r) + V(r)u^2(r)$ is decreasing, f is constant, equal to $u(R)$ on $\Omega \cap B^c$, $V' \leq 0$ on $(0, D_{\Omega})$, and $u(R) \leq u(r)$ on B^c :

$$\begin{aligned} \int_{\Omega \cap B^c} |\nabla f|^2 + Vf^2 &= u(R)^2 \int_{\Omega \cap B^c} V \leq u(R)^2 V(R)|\Omega \cap B^c| \\ &= u(R)^2 V(R)|\Omega^c \cap B| \leq F(R)|\Omega^c \cap B| \leq \int_{\Omega^c \cap B} F, \end{aligned} \tag{14}$$

where we have used the monotonicity of V in the first inequality and the monotonicity of F in the last inequality. Therefore, from (13), (14) and from the fact that $F = |\nabla f|^2 + Vf^2 = |\nabla u|^2 + Vu^2$ on B , we deduce:

$$\begin{aligned} \lambda_1(\Omega, \Delta + V) \int_B u^2 &\leq \lambda_1(\Omega, \Delta + V) \int_\Omega f^2 \leq \int_\Omega |\nabla f|^2 + Vf^2 \\ &= \int_{\Omega \cap B} |\nabla f|^2 + Vf^2 + \int_{\Omega \cap B^c} |\nabla f|^2 + Vf^2 \leq \int_{\Omega \cap B} F + \int_{\Omega^c \cap B} F = \int_B F \\ &= \int_B |\nabla u|^2 + Vu^2 = \lambda_1(B, \Delta + V) \int_B u^2 \end{aligned}$$

and the assertion follows. □

Remark 8 Note that Assumption 3 may look quite involved. However, when we will apply Theorem 6 to the particular case of the Aharonov–Bohm operator we will choose $V = \frac{\nu^2}{\theta^2}$ and condition 3 will take a simpler and more natural form (see Theorem 9).

4 Application to the Aharonov–Bohm Spectrum of the Neumann Problem

We apply now the results of Sect. 3 to the lowest eigenvalue of problem (3). We consider first the general case of manifolds of revolution, then we concentrate on \mathbb{R}^2 , \mathbb{H}^2 and \mathbb{S}^2 .

4.1 Aharonov–Bohm Spectrum on Domains of Manifolds of Revolution

We take a 2-dimensional manifold of revolution M^2 with pole x_0 , and $\theta(r)$ the density of the Riemannian metric in polar coordinates (r, t) around the pole. The 1-form $A_{x_0, \nu} = \nu dt$ is closed, harmonic, and has flux ν around x_0 . It will be called Aharonov–Bohm potential with flux ν .

On a smooth bounded domain Ω of M^2 we have that the spectrum of $\Delta_{A_{x_0, \nu}}$ with Neumann condition, namely problem (3), is made of an increasing sequence of non-negative eigenvalues of finite multiplicity diverging to $+\infty$ (see “Appendix A”).

For all radial functions $u = u(r)$ we have

$$|A_{x_0, \nu}|^2 = \frac{\nu^2}{\theta^2}, \quad \langle \nabla u, A_{x_0, \nu} \rangle = 0,$$

therefore $\Delta_{A_{x_0, \nu}}$ applied to a real, radial function $u = u(r)$, can be written as:

$$\Delta_{A_{x_0, \nu}} u = \Delta u + Vu$$

where $V = \frac{\nu^2}{\theta^2}$.

In “Appendix B.1” we will prove that for a disk centered at the pole $B(x_0, R)$ the spectrum is the union of spectra of a countable family of Sturm–Liouville problems indexed by an integer k (see Lemma 24). In particular, the first eigenvalue, denoted by $\lambda_1(\Omega, A_{x_0, \nu})$, is non-negative, and is positive if and only if $\nu \notin \mathbb{Z}$.

As explained in “Appendix B.1”, thanks to gauge invariance we can take $\nu \in (0, \frac{1}{2}]$ and in that case the first eigenfunction is real and radial.

We denote it by $u = u(r)$. Moreover, we prove in “Appendix B.1” that $u > 0$ and $u' > 0$ for all $R \in (0, \bar{R})$, where \bar{R} is the first zero of θ' (see Theorem 28).

We can apply Theorem 6, taking $V = \frac{\nu^2}{\theta^2}$. The conditions

$$V' \leq 0, \quad \theta' V' + 2V^2 \theta \leq 0$$

reduce to the conditions

$$\theta' \geq 0, \quad \theta'^2 \geq \nu^2.$$

Assumption. Through all this subsection, we shall always assume $\nu \in (0, \frac{1}{2}]$.

We therefore have the following:

Theorem 9 *Let Ω be a smooth bounded domain with diameter D_Ω in a revolution manifold M^2 with pole x_0 , and with density θ , and let $B(x_0, R)$ be the disk centered in x_0 with the same volume as Ω . Assume $\nu \in (0, \frac{1}{2}]$. If*

- (i) $\theta' \geq 0$ on $(0, D_\Omega)$,
- (ii) $\theta'^2 \geq \nu^2$ on $(0, R)$,

then

$$\lambda_1(\Omega, A_{x_0, \nu}) \leq \lambda_1(B(x_0, R), A_{x_0, \nu}). \tag{15}$$

Equality holds if and only if $\Omega = B(x_0, R)$.

When the flux is 0, inequality (15) reduces to the identity $0 = 0$ for all domains Ω .

Note that Theorem 9 provides sufficient conditions to have a reverse Faber–Krahn inequality on a manifold of revolution, which are quite simple to understand. Theorem 9 works well, as we shall see, in the case of \mathbb{R}^2 and \mathbb{H}^2 . However, in some cases (e.g., spherical domains), condition (ii) is somehow restrictive. In “Appendix B.1” we prove Theorem 29, where we show that (ii) can be replaced by some other (more involved) condition, which, in the case of the sphere, turns out to be less restrictive than (ii). Nevertheless, we decide to keep Theorem 9 here for two reasons: it is simpler and has immediate application in many contexts; it is a consequence of a more general result valid for Schrödinger operators, namely Theorem 6, which we believe has an interest per se.

We now assume that M^2 has infinite diameter, so that we can identify it with the manifold (\mathbb{R}^2, g) , where g is a complete, rotationally invariant metric around x_0 , and consider the 1-form $A_{x_0, \nu}$. If $\theta(r)$ is the density of the Riemannian measure, then it is

well-known that the Gaussian curvature of (\mathbb{R}^2, g) is given by $K = -\frac{\theta''}{\theta}$. Assuming $K \leq 0$, we will get $\theta'' \geq 0$, and since $\theta'(0) = 1$, we immediately obtain $\theta'(r) \geq 1$ for all $r > 0$. The assumptions (i) and (ii) of Theorem 9 are met, thus we have the following:

Corollary 10 *Let Ω be a smooth bounded domain in the manifold of revolution (\mathbb{R}^2, g) with pole at x_0 and non-positive Gaussian curvature. Let $B(x_0, R)$ be the disk of radius R centered at x_0 such that $|B(x_0, R)| = |\Omega|$. Then*

$$\lambda_1(\Omega, A_{x_0, \nu}) \leq \lambda_1(B(x_0, R), A_{x_0, \nu}),$$

with equality if and only if $\Omega = B(x_0, R)$.

In particular, Corollary 10 applies to \mathbb{R}^2 with its Euclidean Riemannian metric, and to \mathbb{H}^2 . Since all points in \mathbb{R}^2 and \mathbb{H}^2 can be chosen as poles of the manifold, this gives a proof of Theorem 1.

Another consequence of Theorem 9 is the following

Corollary 11 *Let $B(p, R)$ be a disk in \mathbb{R}^2 or \mathbb{H}^2 , punctured at $x_0 \in B(p, R)$. Then*

$$\lambda_1(B(p, R), A_{x_0, \nu}) \leq \lambda_1(B(x_0, R), A_{x_0, \nu}),$$

that is, among all disks of the same measure, the first eigenvalue is maximized by the disk punctured at its center.

4.2 Aharonov–Bohm Spectrum on Domains of \mathbb{S}^2

For the standard sphere \mathbb{S}^2 of curvature 1 we have

$$\theta(r) = \sin(r),$$

hence a direct application of Theorem 9 is possible only for domains Ω such that $|\Omega| \leq |B(x_0, \pi/3)|$ and contained in a hemisphere centered at x_0 .

Assumption. Through all this subsection, we shall always assume $\nu \in (0, \frac{1}{2}]$.

The condition $|\Omega| \leq |B(x_0, \pi/3)|$ ensures $\theta'^2 \geq \nu^2$ on $B(x_0, R)$, in fact $R \leq \frac{\pi}{3}$ and then $\cos(r)^2 \geq \frac{1}{4}$.

The condition that Ω is contained in a hemisphere centered at x_0 ensures that $\theta' \geq 0$ on $\Omega \cup B(x_0, R)$.

Note that, restrictions to the class of spherical domains for which one usually proves isoperimetric inequalities are natural and common. For example, the Szegő–Weinberger inequality for the second eigenvalue of the Neumann Laplacian on spherical domains is proved under the assumption that the domain is contained in a hemisphere [3] or that it is simply connected with total area less than 2π [4]. In the first case, the approach is that of Weinberger for planar domains [24] (which is the one that we have used up to now), while in the second case the approach is that of Szegő by conformal transplantation [23].

The hypothesis of Theorem 9 are not sufficient to cover the case of spherical domains contained in a hemisphere centered at the pole of the magnetic field, which is the natural counterpart of the results in [3]. However, we prove in Theorem 29 that the function $F(r) = u'(r)^2 + \frac{v^2}{\theta(r)^2}u(r)^2$ is decreasing in $(0, R)$ for all $R \in (0, \bar{R})$ under suitable assumptions, even if $\theta' < v$. Recall that \bar{R} is the first zero of θ' and R is such that $|B(x_0, R)| = |\Omega|$. The fact that $F' \leq 0$, together with $\theta' > 0$, implies the hypotheses of Theorem 6 for $V = \frac{v^2}{\theta^2}$ and henceforth that of Theorem 10. In particular, the assumptions of Theorem 29 for \mathbb{S}^2 reduce to

- (i) $\theta' > 0$,
- (ii) $v(v + 1)\theta^2 - v^2 + v^2(\theta')^2 + v\theta\theta'' \geq 0$,

which are clearly satisfied as long as $R \leq \frac{\pi}{2}$. In view of this, we only need to assume that Ω is contained in a hemisphere centered at the pole x_0 . Then we have the expected result

Theorem 12 *Let Ω be a smooth bounded domain contained in a hemisphere centered at x_0 . Let $B(x_0, R)$ be the disk centered in x_0 with the same volume as Ω . Then*

$$\lambda_1(\Omega, A_{x_0, v}) \leq \lambda_1(B(x_0, R), A_{x_0, v}),$$

with equality if and only if $\Omega = B(x_0, R)$.

Note that this result coincides with that of Ashbaugh–Benguria [3].

We also deduce the following

Corollary 13 *Let $B(p, R) \subset \mathbb{S}^2$ be a spherical disk punctured at $x_0 \in B(p, R)$. If $R \leq \frac{\pi}{4}$, then*

$$\lambda_1(B(p, R), A_{x_0, v}) \leq \lambda_1(B(x_0, R), A_{x_0, v}),$$

that is, among all disks of the same measure and radius smaller than $\frac{\pi}{4}$, the first eigenvalue is maximized by the disk punctured at its center.

4.3 Szegő's Isoperimetric Inequality on Spheres

Theorem 12 is valid for all bounded domains which are contained in a hemisphere centered at the pole x_0 of $A_{x_0, v}$. We have used the Weinberger's argument for its proof. If we take the Szegő's point of view, we are able to extend the result to the class of simply connected domains on the sphere with area less than 2π . Namely, we have:

Theorem 14 *Let Ω be a bounded and simply connected domain in \mathbb{S}^2 with $|\Omega| \leq 2\pi$ and $-x_0 \notin \Omega$, and let $B(x_0, R)$ be the disk in \mathbb{S}^2 centered at x_0 with $|B(x_0, R)| = |\Omega|$. Then*

$$\lambda_1(\Omega, A_{x_0, v}) \leq \lambda_1(B(x_0, R), A_{x_0, v}).$$

If $v \notin \mathbb{Z}$, equality holds if and only if $\Omega = B(x_0, R)$.

Proof Through the stereographic projection f we identify a point $re^{it} \in \mathbb{S}^2$ (r is the distance to x_0 , and t is the angular coordinate) with $z = f(re^{it}) = \tan(r/2)e^{it} \in \mathbb{C}$. Then, for any function v defined on Ω we have

$$\int_{\Omega} v = \int_{f(\Omega)} (v \circ f^{-1}) \frac{4}{(1 + |z|^2)^2} dz.$$

Moreover,

$$|\Omega| = \int_{f(\Omega)} \frac{4}{(1 + |z|^2)^2} dz, \quad |\partial\Omega| = \int_{f(\partial\Omega)} \frac{2}{1 + |z|^2} dz.$$

Let u be an eigenfunction associated to $\lambda_1(B(x_0, R), A_{x_0, v})$. As proved in Theorem 28, we know that $u = u(r)$ is real, radial, and can be chosen such that $u, u' > 0$ on $(0, R)$. In fact, from the our assumptions we have $R \leq \frac{\pi}{2}$, thus Theorem 28 applies. Then, $u \circ f^{-1}$ is positive and increasing as well.

Let now $g : f(B(x_0, R)) \rightarrow f(\Omega)$ be a conformal map with $f(0) = 0$. To simplify our notation, we will set $\tilde{\Omega} = f(\Omega)$ and $\tilde{B} = f(B(x_0, R))$. Note that \tilde{B} is a ball centered at 0 of radius $T := \tan(R/2)$.

We set

$$\hat{u} := u \circ f^{-1} \circ g^{-1} \circ f.$$

Then \hat{u} is a function defined in Ω , the conformal transplantation of u through the map $f^{-1} \circ g^{-1} \circ f$. The conformal invariance of the Aharonov–Bohm energy, proved in ‘‘Appendix C’’ tells that

$$\int_{\Omega} |\nabla^{A_{x_0, v}} \hat{u}|^2 = \int_{B(x_0, R)} |\nabla^{A_{x_0, v}} u|^2,$$

then $\hat{u} \in H^1_{A_{x_0, v}}(\Omega)$ and it is a suitable test function for the min–max principle (30) for $\lambda_1(\Omega, A_{x_0, v})$. In particular

$$\lambda_1(\Omega, A_{x_0, v}) \leq \frac{\int_{\Omega} |\nabla^{A_{x_0, v}} \hat{u}|^2}{\int_{\Omega} \hat{u}^2} = \frac{\int_{B(x_0, R)} |\nabla^{A_{x_0, v}} u|^2}{\int_{\Omega} \hat{u}^2}. \tag{16}$$

If

$$\int_{\Omega} \hat{u}^2 \geq \int_{B(x_0, R)} u^2 \tag{17}$$

then we conclude from (16)

$$\lambda_1(\Omega, A_{x_0, v}) \leq \frac{\int_{B(x_0, R)} |\nabla^{A_{x_0, v}} u|^2}{\int_{B(x_0, r)} u^2} = \lambda_1(B(x_0, R), A_{x_0, v})$$

which is what we want. Note that (17) is equivalent to

$$\int_{\tilde{B}} U |g'|^2 (\sigma \circ g) \geq \int_{\tilde{B}} U \sigma, \tag{18}$$

where

$$U = (u \circ f^{-1})^2$$

and $\sigma(z) = \frac{4}{(1+|z|^2)^2}$. We note that $U' > 0$ on $(0, T) = (0, \tan(R/2))$. Let us set $B_r = B(0, r)$ the disk in \mathbb{C} centered at 0 of radius r . In particular $\tilde{B} = B_T = B_{\tan(R/2)}$. We define

$$a(r) := \int_{B_r} |g'|^2 (\sigma \circ g), \quad v(r) = \int_{B_r} \sigma = 2\pi \int_0^r \frac{4s}{(1+s^2)^2} ds = \frac{4\pi r^2}{1+r^2}.$$

Then, since $|\Omega| = |B(x_0, R)|$, we have $a(T) = v(T)$.

We write a differential inequality for $a(r)$:

$$a'(r) = \int_{\partial B_r} |g'|^2 (\sigma \circ g) \geq \frac{\left(\int_{\partial B_r} |g'| (\sigma \circ g)^{1/2}\right)^2}{2\pi r}. \tag{19}$$

If $\Omega_r = (g \circ \sigma)^{-1}(B_r) \in \mathbb{S}^2$, then

$$|\partial \Omega_r| = \int_{\partial B_r} |g'| (\sigma \circ g)^{1/2}, \quad |\Omega_r| = \int_{B_r} |g'|^2 (\sigma \circ g) = a(r).$$

The isoperimetric inequality $|\partial \Omega_r|^2 \geq |\Omega_r|(4\pi - |\Omega_r|)$ holds for spherical domains, hence from (19)

$$a'(r) \geq \frac{a(r)(4\pi - a(r))}{2\pi r}. \tag{20}$$

Then, the function

$$r \mapsto \frac{a(r)}{r^2(4\pi - a(r))}$$

is not decreasing as long as $a(r) \leq 4\pi$, just take the derivative and use (20). We see now that

$$\frac{v(r)}{r^2(4\pi - v(r))} = 1$$

for all r , and in particular, since $v(T) = a(T)$,

$$\frac{a(r)}{r^2(4\pi - a(r))} \leq \frac{a(T)}{T^2(4\pi - a(T))} = 1$$

and then

$$a(r) \leq \frac{4\pi r^2}{1 + r^2} = \int_{B_r} \sigma = v(r).$$

This gives the desired result. In fact, since U is radial with $U' > 0$,

$$\begin{aligned} \int_{\bar{B}} U|g'|^2(\sigma \circ g) &= \int_0^T U(r)a'(r)dr \\ &= U(T)a(T) - \int_0^T U'(r)a(r)dr = U(T)v(T) - \int_0^T U'(r)a(r)dr \\ &\geq U(T)v(T) - \int_0^T U'(r)v(r)dr = \int_{\bar{B}} U\sigma. \end{aligned} \tag{21}$$

This proves (18) and then the isoperimetric inequality for $\lambda_1(\Omega, A_{x_0, v})$. Finally, if equality holds in the isoperimetric inequality for $\lambda_1(\Omega, A_{x_0, v})$, then equality holds in the isoperimetric inequality $|\partial\Omega_r|^2 \geq |\Omega_r|(4\pi - |\Omega_r|)$ used in (19) for all r , hence all Ω_r are spherical disks, and $\Omega = \Omega_T$ as well.

□

5 The Magnetic Steklov Problem: Brock’s and Weinstock’s Inequalities

We now focus on the magnetic Steklov problem on a bounded smooth domain $\Omega \subset \mathbb{R}^2$, namely problem (4). The min-max principle for the first eigenvalue reads

$$\sigma_1(\Omega, A_{x_0, v}) = \inf_{0 \neq u \in H^1_{A_{x_0, v}}(\Omega, \mathbb{C})} \frac{\int_{\Omega} |\nabla^{A_{x_0, v}} u|^2}{\int_{\partial\Omega} |u|^2}, \tag{22}$$

where $H^1_{A_{x_0, v}}(\Omega, \mathbb{C})$ denotes the standard magnetic Sobolev space (see “Appendix A” for the precise definition).

If we consider the maximisation problem for the lowest eigenvalue under volume constraint, we have Brock’s Theorem for $\sigma_1(\Omega, A_{x_0, v})$:

Theorem 15 *Let Ω be a smooth bounded domain in \mathbb{R}^2 , $x_0 \in \mathbb{R}^2$ a fixed pole, and let $B(x_0, r)$ be the disk with the same measure as Ω . Let $v \in (0, \frac{1}{2}]$. Then:*

$$\sigma_1(\Omega, A_{x_0, v}) \leq \sigma_1(B(x_0, R), A_{x_0, v}) = \frac{\sqrt{\pi} v}{|\Omega|^{\frac{1}{2}}}.$$

Equality holds if and only if $\Omega = B(x_0, r)$.

Proof From (22) we have:

$$\sigma_1(\Omega, A_{x_0, v}) \leq \frac{\int_{\Omega} |\nabla^A u|^2}{\int_{\partial\Omega} |u|^2}$$

for all $u \in H_A^1(\Omega, \mathbb{C})$. Then we choose $u = r^v$ which is the first eigenfunction for any disk centered at x_0 (see ‘‘Appendix B.4’’). One has

$$|\nabla^A u|^2 = 2v^2 r^{2v-2}.$$

In particular, since $v \in (0, \frac{1}{2}]$,

$$\int_{\Omega} |\nabla^A u|^2 = 2v^2 \int_{\Omega} r^{2v-2} \leq 2v^2 \int_{B(0, R)} r^{2v-2} = 2\pi v R^{2v}.$$

In fact,

$$\int_{\Omega \cap B(0, R)^c} r^{2v-2} \leq R^{2v-2} |\Omega \cap B(0, R)^c| = R^{2v-2} |\Omega^c \cap B(0, R)| \leq \int_{\Omega^c \cap B(0, R)} r^{2v-2}.$$

Here $R = \frac{|\Omega|^{1/2}}{\pi^{1/2}}$ because $B(x_0, R)$ has the same volume of Ω .

We recall a well-known fact: for all $p \geq 0$,

$$\int_{\partial\Omega} r^p \geq 2\pi^{\frac{1-p}{2}} |\Omega|^{\frac{p+1}{2}}.$$

When $p = 0$ this is just the classical isoperimetric inequality. For $p > 0$ this inequality says that the infimum of $\int_{\partial\Omega} r^p$ among all domains with fixed measure is attained by the ball centered at x_0 , which is the unique minimizer. This result is proved in [2]. Using u as test function for $\sigma_1(\Omega, A_{x_0, v})$ and the isoperimetric inequality above with $p = 2v$ we obtain

$$\sigma_1(\Omega, A_{x_0, v}) \leq \frac{2v\pi R^{2v}}{2\pi^{\frac{1-2v}{2}} |\Omega|^{\frac{2v+1}{2}}}, \tag{23}$$

that is

$$|\Omega|^{\frac{1}{2}} \sigma_1(\Omega, A_{x_0, v}) \leq \pi^{\frac{1}{2}} v. \tag{24}$$

□

By gauge invariance (see ‘‘Appendix A.4’’), if $v \notin (0, \frac{1}{2}]$, we can replace v in (23) and (24) by $\inf_{k \in \mathbb{Z}} |v - k|$.

If we consider instead the problem of maximising the lowest eigenvalue under perimeter constraint, we have Weinstock’s Theorem for $\sigma_1(\Omega, A_{x_0, v})$:

Theorem 16 *Let Ω be bounded simply connected domain in \mathbb{R}^2 , $x_0 \in \mathbb{R}^2$ be a fixed pole, and let $B(x_0, r)$ the disk with the same perimeter of Ω . Let $v \in (0, \frac{1}{2}]$. Then:*

$$\sigma_1(\Omega, A_{x_0, v}) \leq \sigma_1(B(x_0, R), A_{x_0, v}) = \frac{2\pi}{|\partial\Omega|} v.$$

Equality holds if and only if $\Omega = B(x_0, R)$.

Proof Assume for simplicity that $x_0 = 0$. Take the unique conformal map $\Phi : \Omega \rightarrow B$, where B is the unit disk centered at the origin, and with $\Phi(0) = 0$, and fix the eigenfunction $u = r^v$ of the unit disk, associated to $\sigma_1(B, A_{x_0, v}) = v$. We refer to ‘‘Appendix B.4’’ for more details. We take as test-function

$$\hat{u} = u \circ \Phi.$$

Then:

$$\begin{aligned} \sigma_1(\Omega, A) \int_{\partial\Omega} |\hat{u}|^2 &= \sigma_1(\Omega, \hat{A}) \int_{\partial\Omega} |\hat{u}|^2 \leq \int_{\Omega} |d^{\hat{A}} \hat{u}|^2 \\ &= \int_D |d^A u|^2 = \sigma_1(D, A) \int_{\partial D} |u|^2 = 2\pi v \end{aligned}$$

where on the first line, we used gauge invariance (Lemma 36) and in the third we used the conformal invariance of the magnetic energy (Lemma 35). Here $\hat{A} = \Phi^* A$. On the other hand, $\hat{u} = 1$ on $\partial\Omega$ so that

$$\int_{\partial\Omega} |\hat{u}|^2 = |\partial\Omega|.$$

The conclusion is

$$\sigma_1(\Omega, A_{x_0, v}) \leq \frac{2\pi v}{|\partial\Omega|}$$

as asserted. □

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A Setting

In this section we prove that problems (3) and (4) admit purely discrete spectrum made of non-negative eigenvalues of finite multiplicity diverging to $+\infty$. Through all this section Ω will denote a bounded domain in \mathbb{R}^2 . We will also prove the gauge invariance property of the two problems.

A.1 Functional Setting

Let $H^1(\Omega, \mathbb{C})$ be the standard Sobolev space of complex-valued functions. As for problem (3), we shall assume that Ω is such that the embedding $H^1(\Omega, \mathbb{C}) \subset L^2(\Omega, \mathbb{C})$ is compact. As for problem (4) we shall assume that the trace operator $\gamma_0 : H^1(\Omega, \mathbb{C}) \rightarrow L^2(\partial\Omega, \mathbb{C})$ is compact. In both cases the compactness is guaranteed if Ω is smooth (Lipschitz is enough).

Let $x_0 \in \mathbb{R}^2$ and let $C_{x_0}^\infty(\Omega, \mathbb{C})$ be the space of smooth functions on Ω vanishing in a neighborhood of x_0 . We introduce the magnetic Sobolev space $H_A^1(\Omega, \mathbb{C})$ defined as the closure of $\{u \in C_{x_0}^\infty(\Omega, \mathbb{C}) : \nabla^A u, u \in L^2(\Omega, \mathbb{C})\}$ with respect to the norm

$$\|u\|_A^2 := \int_{\Omega} |\nabla^A u|^2 + |u|^2, \quad \forall u \in C_0^\infty(\Omega, \mathbb{C}) : \nabla^A u, \quad u \in L^2(\Omega, \mathbb{C}).$$

To simplify the notation, through all this section we shall denote by A the Aharonov–Bohm potential vA_0 with pole at x_0 . We will make use of equivalent norms, suitable to our problems. One fundamental equivalent norm is given in the following proposition.

Proposition 17 *If $v \notin \mathbb{Z}$, the norm $\|u\|_A$ is equivalent to the following norm:*

$$\|u\|_{A'}^2 := \int_{\Omega} |\nabla u|^2 + |u|^2 + \frac{|u|^2}{|x - x_0|^2}.$$

If Ω does not contain the pole x_0 , then the equivalence is immediate to check. In any case, the proof of Proposition 17 is standard, and follows from the well-known Hardy-type inequality proved in [18]:

Lemma 18 *For any $R > 0$ and any $u \in H^1(B(x_0, R), \mathbb{C})$ we have*

$$\int_{B(x_0, R)} |\nabla^A u|^2 dx \geq C^2 \int_{B(x_0, R)} \frac{|u|^2}{|x - x_0|^2}, \tag{25}$$

where $C = \inf\{|v - k| : k \in \mathbb{Z}\}$. In particular, if $v \in (0, \frac{1}{2}]$, then $C = v$.

Proof Inequality (25) is proved in [18]. We recall briefly here the proof. We use polar coordinates (r, θ) in \mathbb{R}^2 centered at x_0 , and express $\int_{B(x_0, R)} |\nabla^A u|^2 dx$ in this new coordinate system. We have

$$\int_{B(x_0, R)} |\nabla^A u|^2 dx = \int_0^R \int_0^{2\pi} \left(|\partial_r u|^2 + \frac{1}{r^2} |\partial_\theta u - i v u|^2 \right) r d\theta dr.$$

We focus on $\int_0^{2\pi} |\partial_\theta u - i v u|^2 d\theta$. We estimate

$$\inf_{0 \neq u \in H^1(\mathbb{S}^1, \mathbb{C})} \frac{\int_0^{2\pi} |\partial_\theta u - i v u|^2 d\theta}{\int_0^{2\pi} r^2 |u|^2 d\theta}.$$

This infimum corresponds to the first eigenvalue, which we denote by μ_1 , of

$$-\partial_{\theta\theta}^2 u + 2i v \partial_\theta u + v^2 u = \mu u$$

on \mathbb{S}^1 . A set of $L^2(\mathbb{S}^1, \mathbb{C})$ -normalized eigenfunctions is given by $(2\pi)^{-1/2} e^{i\theta k}$, $k \in \mathbb{Z}$, with corresponding eigenvalue $\mu = (k - v)^2$. If $v \notin \mathbb{Z}$, then $\mu_1 = \inf\{(k - v)^2 : k \in \mathbb{Z}\}$.

This concludes the proof. □

We recall the following well-known result

Lemma 19 *Let Ω be a bounded domain in \mathbb{R}^n with Lipschitz boundary. Then the norm $\|u\|_{H^1(\Omega, \mathbb{C})}^2$ is equivalent to the following norm:*

$$\|u\|_{\partial}^2 := \int_{\Omega} |\nabla u|^2 + \int_{\partial\Omega} |u|^2, \tag{26}$$

for all $u \in H^1(\Omega, \mathbb{C})$. With abuse of notation, we write $\int_{\partial\Omega} |u|^2$ in place of $\int_{\partial\Omega} |\gamma_0(u)|^2$.

A consequence of Lemmas 17 and 19 is the following

Lemma 20 *For $v \notin \mathbb{Z}$, the norm $\|u\|_A$ is equivalent to the following norms:*

$$\|u\|_{A, \partial}^2 := \int_{\Omega} |\nabla^A u|^2 + \int_{\partial\Omega} |u|^2, \tag{27}$$

and

$$\|u\|_{A', \partial}^2 := \int_{\Omega} |\nabla u|^2 + \frac{|u|^2}{|x - x_0|^2} + \int_{\partial\Omega} |u|^2. \tag{28}$$

Remark 21 Lemma 17 essentially says that, if $v \notin \mathbb{Z}$

$$H_A^1(\Omega, \mathbb{C}) = \left\{ u \in H^1(\Omega, \mathbb{C}) : \frac{u}{|x - x_0|} \in L^2(\Omega, \mathbb{C}) \right\}.$$

In particular, $H_A^1(\Omega, \mathbb{C}) \subset H^1(\Omega, \mathbb{C})$.

A.2 The Neumann Eigenvalue Problem

Problem (3) is understood in the weak sense, namely, find a function $u \in H_A^1(\Omega, \mathbb{C})$ and number $\lambda \in \mathbb{C}$ such that

$$\int_{\Omega} \langle \nabla^A u, \overline{\nabla^A \phi} \rangle = \lambda \int_{\Omega} u \bar{\phi}, \quad \forall \phi \in H_A^1(\Omega, \mathbb{C}).$$

We rewrite it as

$$\int_{\Omega} \langle \nabla^A u, \overline{\nabla^A \phi} \rangle + \int_{\Omega} u \bar{\phi} = (\lambda + 1) \int_{\Omega} u \bar{\phi}, \quad \forall \phi \in H_A^1(\Omega, \mathbb{C}). \tag{29}$$

Since the quadratic form on the left-hand side of (29) is bounded and coercive on $H_A^1(\Omega, \mathbb{C})$ and the embedding $H_A^1(\Omega, \mathbb{C}) \subset L^2(\Omega, \mathbb{C})$ is compact, problem (29) is recast to an eigenvalue problem for a compact self-adjoint operator on the Hilbert space $H_A^1(\Omega, \mathbb{C})$. Standard Spectral Theory implies that problem (29) admits a sequence of eigenvalues

$$-\infty < \lambda_1 \leq \lambda_2 \leq \dots \leq \lambda_k \leq \dots \nearrow +\infty.$$

Moreover, there exists a Hilbert basis of $H_A^1(\Omega, \mathbb{C})$ of eigenfunctions of (29). The eigenfunctions can be chosen to form an orthonormal basis of $L^2(\Omega, \mathbb{C})$ as well.

We finally recall the variational characterization of the eigenvalues:

$$\lambda_k = \min_{\substack{U \subset H_A^1(\Omega, \mathbb{C}) \\ \dim U = k}} \max_{0 \neq u \in U} \frac{\int_{\Omega} |\nabla^A u|^2}{\int_{\Omega} |u|^2}. \tag{30}$$

In particular, $\lambda_1 \geq 0$. If $\nu \notin \mathbb{Z}$ and $x_0 \in \Omega$, we deduce from Lemma 18 that $\lambda_1 > 0$. On the other hand, if $\nu \in \mathbb{Z}$, then $e^{i\nu\theta} \in H_A^1(\Omega, \mathbb{C})$ is an eigenfunction corresponding to the eigenvalue $\lambda = 0$. We deduce by gauge invariance (see ‘‘Appendix A.4’’) that in this case the eigenvalues of (3) coincide with those of the Neumann Laplacian (i.e., $\nu = 0$). If x_0 belongs to the unbounded component of Ω^c , then the flux of $A_{x_0, \nu}$ is zero, hence $A_{x_0, \nu} = \nabla \phi$ with $\Delta \phi = 0$ on Ω , and the eigenvalues of (3) coincide with those of the Neumann Laplacian. If x_0 belongs to some bounded component of Ω^c , the flux is ν , and if it is not an integer, inequality (25) with $C > 0$ holds with $B(x_0, R)$ replaced by an annular region $A \subset \Omega$ such that x_0 belongs to the bounded component of A^c (one immediately realizes that (25) holds if we replace $B(x_0, R)$ by $B(x_0, R) \setminus B(x_0, r)$ for any $0 < r < R$). Then $\lambda_1 > 0$.

A.3 The Steklov Eigenvalue Problem

Problem (4) is understood in the weak sense, namely, find a function $u \in H_A^1(\Omega, \mathbb{C})$ and number $\sigma \in \mathbb{C}$ such that

$$\int_{\Omega} \langle \nabla^A u, \overline{\nabla^A \phi} \rangle = \sigma \int_{\partial\Omega} u \bar{\phi}, \quad \forall \phi \in H_A^1(\Omega, \mathbb{C}).$$

We rewrite it as

$$\int_{\Omega} \langle \nabla^A u, \overline{\nabla^A \phi} \rangle + \int_{\partial\Omega} u \bar{\phi} = (\sigma + 1) \int_{\partial\Omega} u \bar{\phi}, \quad \forall \phi \in H_A^1(\Omega, \mathbb{C}). \tag{31}$$

For the study of problem (31) we consider the space $H_A^1(\Omega, \mathbb{C})$ endowed with the equivalent norm $\|u\|_{A,\partial}$.

Since the quadratic form on the left-hand side of (31) is bounded and coercive on $H_A^1(\Omega, \mathbb{C})$ (with the equivalent norm $\|u\|_{A,\partial}$, see Lemma 20) and the trace operator $\gamma_0 : H_A^1(\Omega, \mathbb{C}) \rightarrow L^2(\partial\Omega, \mathbb{C})$ is compact, problem (31) is recast to an eigenvalue problem for a compact self-adjoint operator on the Hilbert space $H_A^1(\Omega, \mathbb{C})$. Standard Spectral Theory implies that problem (31) admits a sequence of eigenvalues

$$-\infty < \sigma_1 \leq \sigma_2 \leq \dots \leq \sigma_k \leq \dots \nearrow +\infty.$$

Moreover, there exists a Hilbert basis of $H_{A,h}^1(\Omega, \mathbb{C})$ of eigenfunctions of (31), where

$$H_{A,h}^1(\Omega, \mathbb{C}) = \left\{ u \in H_A^1(\Omega, \mathbb{C}) : \int_{\Omega} \langle \nabla^A u, \overline{\nabla^A \phi} \rangle = 0, \forall \phi \in H_A^1(\Omega, \mathbb{C}) : \gamma_0(u) = 0 \right\}.$$

The eigenfunctions can be chosen to form an orthonormal basis of $L^2(\partial\Omega, \mathbb{C})$ as well.

We finally recall the variational characterization of the eigenvalues:

$$\sigma_k = \min_{\substack{U \subset H_A^1(\Omega, \mathbb{C}) \\ \dim U = k}} \max_{0 \neq u \in U} \frac{\int_{\Omega} |\nabla^A u|^2}{\int_{\partial\Omega} |u|^2}. \tag{32}$$

In particular, $\sigma_1 \geq 0$. As in the Neumann case, when $x_0 \in \Omega$, $\sigma_1 > 0$ if and only if $v \notin \mathbb{Z}$. If $v \in \mathbb{Z}$ or x_0 belongs to the unbounded connected component of Ω^c , the eigenvalues coincide with those of the usual Steklov problem (i.e., $v = 0$). If $v \notin \mathbb{Z}$ and x_0 belongs to some bounded connected component of Ω^c , then $\sigma_1 > 0$.

We finally remark that the discussion contained in this section applies with no essential modifications to the case of bounded domains in two-dimensional Riemannian manifolds M^2 , and also to the case in which multiple singularities (of Aharonov–Bohm type) occur, i.e., the domain is punctured in many points.

A.4 Gauge Invariance

It is well-known that, if $A_{x_0, \nu}$ and $A_{\nu'}$ have fluxes ν, ν' which differ by an integer, i.e., $\nu - \nu' \in \mathbb{Z}$, then $\Delta_{A_{x_0, \nu}}$ and $\Delta_{A_{\nu'}}$ are unitarily equivalent. In fact, one can define a multivalued function

$$\psi(x) = \int_{y_0}^x (A_{x_0, \nu} - A_{\nu'}),$$

where y_0 is any reference point (but not the pole). By the given condition, ψ is multivalued but $e^{i\psi}$ is well-defined. Gauge invariance identity is:

$$\Delta_{A_{\nu'}} e^{-i\psi} = \Delta_{A_{x_0, \nu} - \nabla\psi} e^{i\psi} = e^{-i\psi} \Delta_{A_{x_0, \nu}}.$$

One also observes that, if

$$\nabla^{A_{x_0, \nu}} u = \nabla u - iuA_{x_0, \nu},$$

then one has the identity:

$$\nabla^{A_{x_0, \nu} - \nabla\psi} (e^{-i\psi} u) = e^{-i\psi} \nabla^{A_{x_0, \nu}} u.$$

From this one gets easily that, if u is an eigenfunction of the magnetic Neumann problem (3) associated to λ , then $e^{-i\psi} u$ is an eigenfunction of the corresponding problem with potential $A_{x_0, \nu} - \nabla\psi$, associated to the same eigenvalue. The same assertion holds for the magnetic Steklov problem (4).

Since we are interested in the case $\nu \notin \mathbb{Z}$, we can always assume, without loss of generality, that

$$\nu \in \left(0, \frac{1}{2}\right]. \tag{33}$$

B Aharonov–Bohm Eigenvalues in Two-Dimensional Manifolds of Revolution

Let (M^2, g) be a 2-dimensional compact manifold of revolution with pole x_0 and diameter D . We assume that the metric, in polar coordinates (r, t) around the pole, is given by:

$$g = \begin{pmatrix} 1 & 0 \\ 0 & \theta^2(r), \end{pmatrix}$$

with $\theta(r)$ smooth and positive. Here r is the distance to x_0 (i.e., the radial coordinate). We shall always assume that

$$\theta > 0 \text{ on } (0, D), \quad \theta'(0) = 1.$$

We also define

$$\bar{R} = \min\{R : \theta'(R) = 0\} \in (0, D).$$

If such minimum does not exist, then we set $\bar{R} = D$.

If $f = f(r, t)$ the Laplacian writes:

$$\Delta f = -f'' - \frac{\theta'}{\theta} f' - \frac{1}{\theta^2} \frac{\partial^2 f}{\partial t^2},$$

where by f' we denote the derivative of f with respect to the radial variable r .

We consider the harmonic 1-form $A_{x_0, \nu}$ which is written in polar coordinates as

$$A_{x_0, \nu} = \nu dt,$$

where ν is the flux around x_0 .

Lemma 22 For any $f = f(r, t)$:

$$\Delta_{A_{x_0, \nu}} f = -f'' - \frac{\theta'}{\theta} f' - \frac{1}{\theta^2} \frac{\partial^2 f}{\partial t^2} + \frac{\nu^2}{\theta^2} f + 2i \frac{\nu}{\theta^2} \frac{\partial f}{\partial t}.$$

If $f(r, t) = u(r)e^{ikt}$ then

$$\frac{\partial f}{\partial t} = iku e^{ikt}, \quad \frac{\partial^2 f}{\partial t^2} = -k^2 u e^{ikt}$$

and therefore:

$$\Delta_{A_{x_0, \nu}} f = \left(-u'' - \frac{\theta'}{\theta} u' + \frac{(k - \nu)^2}{\theta^2} u \right) e^{ikt}.$$

Proof The magnetic Laplacian writes, for co-closed potentials $A_{x_0, \nu}$:

$$\Delta_{A_{x_0, \nu}} f = \Delta f + |A_{x_0, \nu}|^2 f + 2i \langle df, A_{x_0, \nu} \rangle.$$

Normalizing the basis $\left(\frac{\partial}{\partial r}, \frac{\partial}{\partial t} \right)$ we obtain the orthonormal basis:

$$(e_1, e_2) = \left(\frac{\partial}{\partial r}, \frac{1}{\theta} \frac{\partial}{\partial t} \right)$$

hence we get:

$$|A_{x_0, \nu}|^2 = \frac{\nu^2}{\theta^2}$$

and

$$2i \langle df, A_{x_0, \nu} \rangle = 2i \frac{\nu}{\theta^2} \frac{\partial f}{\partial t}.$$

The claim immediately follows. □

Corollary 23 *If $f(r, t) = u(r)e^{ikt}$, then $\Delta_{A_{x_0, \nu}} f = \lambda f$ reads:*

$$u'' + \frac{\theta'}{\theta} u' + \left(\lambda - \frac{(k - \nu)^2}{\theta^2} \right) u = 0. \tag{34}$$

B.1 Eigenfunctions of the Neumann Problem

We consider now problem (3) on disks $B(x_0, R)$ of radius R centered at the pole of a manifold of revolution M^2 .

Looking for eigenfunctions of the form $f(r, t) = u(r)e^{ikt}$, $k \in \mathbb{Z}$, we find, thanks to (34), that u is a bounded solution of the following singular Sturm–Liouville problem.

$$\begin{cases} u'' + \frac{\theta'}{\theta} u' + \left(\lambda - \frac{(k-\nu)^2}{\theta^2} \right) u = 0 & \text{in } (0, R), \\ u'(R) = 0. \end{cases} \tag{35}$$

If $\nu \notin \mathbb{Z}$, we see from (35) that necessarily $u(0) = 0$ (and in fact, any function in $H^1_A(\Omega, \mathbb{C})$ vanishes at x_0). If $\nu \in \mathbb{Z}$, we see that any bounded solution of (35) with $\nu = k$ satisfies $u'(0) = 0$, while for $\nu \neq k$, $u(0) = 0$. This is exactly the case of the standard Laplacian. The condition $u'(R) = 0$ is the Neumann condition; in fact, $\langle \nabla^{A_{x_0, \nu}} u, N \rangle = \langle \nabla u, N \rangle = u'(R)e^{ikt}$ in the case of a disk centered at x_0 .

The associated quadratic form is

$$\int_0^R \left(u'^2 + V_k u^2 \right) \theta \, dr,$$

where $V_k(r) = \frac{(k-\nu)^2}{\theta^2(r)}$. Note that, if $\nu \in (0, \frac{1}{2}]$, then V_k is increasing in $|k|$, hence:

$$V_k \geq V_0,$$

for all k .

Lemma 24 *Let $\nu \in (0, \frac{1}{2}]$. Then:*

(i) *For each $k \in \mathbb{Z}$, problem (35) admits an infinite sequence of eigenvalues:*

$$0 < \lambda_{k1} < \lambda_{k2} < \dots \leq \lambda_{kj} < \dots \nearrow +\infty$$

with associated bounded eigenfunctions u_{kj} . All eigenvalues are simple and u_{kj} has j zeros in $(0, R)$.

(ii) If $|k| \leq |h|$ then:

$$\lambda_{k1} \leq \lambda_{h1}.$$

In particular, $\min_{k,j} \{\lambda_{kj}\} = \lambda_{01}$.

Proof Point (i) follows from standard theory for singular Sturm–Liouville problems, since $\theta(r)$ is smooth and positive in $(0, R]$, and $\theta(r) \sim \theta'(0)r$ as $r \rightarrow 0^+$.

To prove (ii) it is sufficient to consider $k = 0$. The claim is a consequence of the fact that V_h is increasing in h . Let us take $u = u_{h1}$, an eigenfunction associated to λ_{h1} , as test function for λ_{01} in its variational characterization (which is the analogue of (11) weighted with θ). We get

$$\lambda_{01} \int_0^R u^2 \theta dr \leq \int_0^R (u'^2 + V_0 u^2) \theta dr \leq \int_0^R (u'^2 + V_h u^2) \theta dr = \lambda_{h1} \int_0^R u^2 \theta dr.$$

□

Corollary 25 *We have*

$$\lambda_1(B(x_0, R), A_{x_0, \nu}) = \lambda_{01}.$$

A first eigenfunction is given by $u_{01}(r)$, and it is real and radial.

Note that $\lambda_1(B(x_0, R), A_{x_0, \nu})$ is not necessarily simple (this is the case of $\nu = \frac{1}{2}$ and the unit disk, see “Appendix B.3”). At any rate, Corollary 25 states that we can always find a first eigenfunction which is real and radial. When $\nu \in \mathbb{Z}$, Lemma 24 gives exactly the eigenvalues of the Neumann Laplacian, and in particular $\lambda_{01} = 0$.

We have the following result on completeness of the family of eigenfunctions $\{u_{kj}(r)e^{ikt}\}$.

Theorem 26 *Let*

$$\psi_{kj}(r, t) = u_{kj}(r)e^{ikt}$$

for all $k \in \mathbb{Z}$, $j = 1, 2, \dots$, where u_{kj} are the eigenfunctions of (35). Then $\{\psi_{kj}(r, t)\}_{k,j}$ is a complete system in $L^2(B(x_0, R))$.

Proof If $\nu \in \mathbb{Z}$ the result is well-known: in fact the eigenfunctions $\{\psi_{kj}\}$ are the eigenfunctions of the Neumann Laplacian. Assume then $\nu \notin \mathbb{Z}$. Let us assume that the system of eigenfunctions $\{\psi_{kj}(r, t)\}_{k,j}$ is not complete. Then, there exists an eigenfunction f of (3) which is orthogonal to the span of all ψ_{kj} , that is:

$$\int_{B(x_0, R)} f \overline{\psi_{hi}} = 0 \tag{36}$$

for all $h \in \mathbb{Z}, i = 1, 2, \dots$. Expand $f(r, t)$ in a Fourier series, for every fixed r , and get:

$$f(r, t) = \sum_{k \in \mathbb{Z}} a_k(r) e^{ikt}.$$

As by assumption $f \neq 0$, there exists $h \in \mathbb{Z}$ such that $a_h \neq 0$.

We prove in Lemma 27 here below that, for all $k \in \mathbb{Z}$,

$$a_k(r) = c_k u_{k, j_k}(r) \tag{37}$$

for some $j_k = 1, 2, \dots$ and $c_k \in \mathbb{C}, c_k \neq 0$. The conclusion follows from (37). In fact, we can write

$$f(r, t) = \sum_{k \in \mathbb{Z}} c_k u_{k, j_k}(r) e^{ikt} = \sum_{k \in \mathbb{Z}} c_k \psi_{k, j_k}(r, t)$$

and then, by the orthogonality of ψ_{kj} 's:

$$\int_{B(x_0, R)} f \bar{\psi}_{h, j_h} = c_h \neq 0 \tag{38}$$

which contradicts (36). □

Lemma 27 *Let $\{u_{kj}\}$ denote the eigenfunctions of (35) and assume that $f(r, t) = \sum_{k \in \mathbb{Z}} a_k(r) e^{ikt}$ is a nonzero eigenfunction of (3). Then, for all $k \in \mathbb{Z}$*

$$a_k(r) = c_k u_{k, j_k}(r)$$

for some $j_k = 1, 2, \dots$

Proof We show that for each $k \in \mathbb{Z}$ the function $a_k : [0, R] \rightarrow \mathbb{R}$ is a bounded solution of

$$\begin{cases} a_k'' + \frac{\theta'}{\theta} a_k' + \left(\lambda - \frac{(k-\nu)^2}{\theta^2} \right) a_k = 0 & \text{in } (0, R), \\ a_k'(R) = 0. \end{cases}$$

In particular, since $\nu \notin \mathbb{Z}, a_k(0) = 0$. Hence, since the eigenvalues are simple, for every k , there exists $j = j_k$ and $c_k \in \mathbb{C}$ such that

$$a_k(r) = c_k u_{k, j_k}(r), \quad \lambda = \lambda_{k, j_k}.$$

If we assume $a_k \neq 0$, then $c_k \neq 0$.

First observe that $f(x_0) = 0$ because otherwise f is not in $H_A^1(\Omega, \mathbb{C})$ when $\nu \notin \mathbb{Z}$. Then:

$$0 = f(x_0) = \sum_{k \in \mathbb{Z}} a_k(0) e^{ikt}$$

for all t . Hence, by multiplying by e^{-iht} and integrating on $[0, 2\pi]$ we see that $a_h(0) = 0$. Then $a_k(0) = 0$ for all k . Similarly, the Neumann condition $\frac{\partial f}{\partial r}(R, t) = 0$ implies $a'_k(R) = 0$ for all k .

It remains to show that a_k is an eigenfunction of the problem (35). First observe that

$$a_k(r) = \int_0^{2\pi} f(r, t)e^{-ikt} dt.$$

Hence:

$$-k^2 a_k = \int_0^{2\pi} f(r, t) \frac{\partial^2}{\partial t^2} e^{-ikt} dt = \int_0^{2\pi} \frac{\partial^2 f}{\partial t^2}(r, t) e^{-ikt} dt$$

therefore

$$-\frac{k^2}{\theta^2} a_k = \int_0^{2\pi} \frac{1}{\theta^2} \frac{\partial^2 f}{\partial t^2}(r, t) e^{-ikt} dt. \tag{39}$$

Now we know that $\Delta_{A_{x_0, v}} f = \lambda f$; hence

$$\frac{1}{\theta^2} \frac{\partial^2 f}{\partial t^2} = -\lambda f - f'' - \frac{\theta'}{\theta} f' + \frac{v^2}{\theta^2} f + 2i \frac{v}{\theta^2} \frac{\partial f}{\partial t}.$$

so that

$$\begin{aligned} \int_0^{2\pi} \frac{1}{\theta^2} \frac{\partial^2 f}{\partial t^2} e^{-ikt} dt &= -\lambda \int_0^{2\pi} f e^{-ikt} dt - \int_0^{2\pi} f'' e^{-ikt} dt \\ &\quad - \frac{\theta'}{\theta} \int_0^{2\pi} f' e^{-ikt} dt + \frac{v^2}{\theta^2} \int_0^{2\pi} f e^{-ikt} dt + 2i \frac{v}{\theta^2} \int_0^{2\pi} \frac{\partial f}{\partial t} e^{-ikt} dt. \end{aligned}$$

Integrating by parts:

$$\int_0^{2\pi} \frac{\partial f}{\partial t} e^{-ikt} dt = ik \int_0^{2\pi} f e^{-ikt} dt = ika_k.$$

Then:

$$\int_0^{2\pi} \frac{1}{\theta^2} \frac{\partial^2 f}{\partial t^2} e^{-ikt} dt = -\lambda a_k - a''_k - \frac{\theta'}{\theta} a'_k + \frac{v^2}{\theta^2} a_k - 2k \frac{v}{\theta^2}$$

which substituted in (39), gives:

$$a''_k + \frac{\theta'}{\theta} a'_k + \left(\lambda - \frac{(k - v)^2}{\theta^2} \right) a_k = 0$$

which is the final assertion. □

In the following theorems we collect a few useful properties of the eigenfunction $u = u_{01}$ associated with the first eigenvalue $\lambda = \lambda_{01}$.

Theorem 28 *Let $v \in (0, \frac{1}{2}]$. Let $\lambda > 0$ denote the first eigenvalue of problem (35) with $k = 0$ and $R \in (0, \bar{R}]$, and let u denote a corresponding first eigenfunction. Then the following statements hold:*

- (i) $u(0) = 0$ and $u \neq 0$ on $(0, R)$. In particular, we can choose $u > 0$ on $(0, R)$.
- (ii) $u' > 0$ on $(0, R)$ and $\lambda > \frac{v^2}{\theta(R)^2}$.
- (iii) The first eigenvalue λ is strictly decreasing in R for $R \in (0, \bar{R})$. In particular $\lambda > \bar{\lambda}$, where $\bar{\lambda}$ is the first eigenvalue of (35) with $R = \bar{R}$.

Proof Point (i) has been already discussed in Lemma 24 and follows from standard Sturm–Liouville theory and the fact that $\theta, \theta' > 0$ on $(0, R), \theta(r) \sim \theta'(0)r$ as $r \rightarrow 0^+$.

Consider now (ii). Set $N(r) := \theta(r)u'(r)$. Then $N(0) = N(R) = 0$ and

$$N'(r) = \left(\frac{v^2}{\theta^2(r)} - \lambda \right) \theta(r)u(r).$$

Since $N'(r) > 0$ on $(0, \delta)$ for some $\delta > 0$, we deduce that the function N increases from zero and decreases to zero. It has only one maximum point in $(0, R)$. In fact, since θ and u are strictly positive in $(0, R)$, $N'(r) = 0$ if and only if $\lambda = \frac{v^2}{\theta^2(r)}$. Since $\frac{1}{\theta^2}$ is strictly decreasing in $(0, R)$, it vanishes exactly once in $(0, R)$. This implies that $N > 0$ on $(0, R)$ and therefore $u' > 0$ on $(0, R)$. Moreover $N'(R) < 0$.

As for (iii), let us write $\lambda = \lambda(R)$ to highlight the dependence on R . Then, as $\lambda(R)$ is simple, we derive the identity

$$\int_0^R u^2 \theta \lambda(R) dr = \int_0^R \left(u'^2 + \frac{v^2}{\theta^2} u^2 \right) \theta dr$$

and obtain

$$\lambda'(R) = \left(\frac{v^2}{\theta^2(R)} - \lambda(R) \right) \frac{u^2(R)\theta(R)}{\int_0^R u^2 \theta dr},$$

which is strictly negative by point (ii) as long as $\theta' < 0$, that is, $R \in (0, \bar{R})$. More details can be found e.g., in [9]. □

Theorem 29 *Let $v \in (0, \frac{1}{2}]$. Let $\lambda > 0$ denote the first eigenvalue of problem (35) with $k = 0$ and $R \in (0, \bar{R})$, and let u denote a corresponding eigenfunction. Let $F(r) = u'(r)^2 + \frac{v^2 u(r)^2}{\theta(r)^2}$. If either*

- (a) $\theta'(r) \geq v$ on $(0, R)$, or

(b) $\bar{\lambda}\theta^2 - v^2 + v^2(\theta')^2 + v\theta\theta'' \geq 0$ on $(0, R)$,

then $F'(r) \leq 0$ on $(0, R)$. If moreover θ^v is a solution to (35) for $R = \bar{R}$, then $F'(r) \leq 0$ on $(0, \bar{R})$

Proof Define

$$q(r) = \theta(r) \frac{u'(r)}{u(r)}.$$

Then $F = u'^2 + \frac{v^2u}{\theta^2} = \frac{u^2}{\theta^2}(q^2 + v^2)$, and with standard computations we see that

$$F' = 2qq' \frac{u^2}{\theta^2} + \frac{2u^2}{\theta^3}(q^2 + v^2)(q - \theta').$$

Since $u, \theta, \theta' > 0$ on $(0, R)$, the claim follows provided $q' \leq 0$ and $q \leq \theta'$ on $(0, R)$.

We start by proving $q' \leq 0$ on $(0, R)$. We see that

$$q' = -\theta\lambda + \frac{v^2}{\theta} - \frac{q^2}{\theta}.$$

The local behavior of u near $r = 0^+$ depends only on the differential equation (35), and in particular, using the Frobenius–Taylor expansion, we deduce that $u \sim cr^v$ as $r \rightarrow 0^+$ for some $c > 0$. Therefore $q(0) = v, q'(0) = 0$. Clearly there exists $a > 0$ such that $q \neq v$ on $(0, a)$, otherwise $q' = -\theta\lambda \neq 0$ on $(0, a)$. Choose then (a possibly smaller) $a > 0$ such that q, q' have constant sign on $(0, a)$ (in particular, $q > 0$ on $(0, R)$). We have $q' < 0$ in $(0, a)$. Otherwise, $q > q(0) = v$, but then the differential equation tells then $q' < 0$. Then q starts decreasing. The same argument shows that $0 < q < v$ on $(0, R)$. If q' changes sign in $(0, R)$, let then $b, c, d \in (0, R), b < c < d$, be such that $q'(b) < 0, q'(c) > 0, q'(d) < 0$ and $q(b) = q(c) = q(d) = \bar{q} < v$. Such points exist if q' changes sign since $q(R) = 0$. We note now that, since $\theta' > 0$ in $(0, R), \theta(c) = t\theta(b) + (1 - t)\theta(d)$ for some $t \in (0, 1)$. Then, from the differential equation for q' , we see that

$$\begin{aligned} 0 < q'(c) &= -\lambda\theta(c) + \frac{v^2 - \bar{q}^2}{\theta(c)} \\ &\leq t \left(-\lambda\theta(b) + \frac{v^2 - \bar{q}^2}{\theta(b)} \right) + (1 - t) \left(-\lambda\theta(d) + \frac{v^2 - \bar{q}^2}{\theta(d)} \right) \\ &= tq'(b) + (1 - t)q'(d) < 0, \end{aligned}$$

a contradiction. Then $q' < 0$ on $(0, R)$. Note that $q' < 0$ also in the case $R = \bar{R}$.

It remains to prove that $q \leq \theta'$. Note that, if $\theta' \geq v$ on $(0, R)$, since q is decreasing and $q(0) = v$, we have $q \leq \theta'$. We have used condition (a). Note that this condition has already been found in Theorem 10.

Assume now (b) holds. Let $R \in (0, \bar{R})$. In order to conclude, it is sufficient to prove

$$q \leq (v + \delta)\theta' \tag{40}$$

on $(0, R)$, for some $\delta \in (0, 1 - \nu)$. For any $\delta > 0$ we have $q(0) < (\nu + \delta)\theta'(0)$ and

$$q'(r) = G(r, q(r)),$$

where $G(r, q)$ is defined as

$$G(r, q) := -\lambda\theta(r) + \frac{\nu^2}{\theta(r)} - \frac{q^2}{\theta(r)}.$$

Note also that, since $R \in (0, \bar{R})$, we have from (iii) of Theorem 28 that $\lambda = \bar{\lambda} + \epsilon$ for some $\epsilon > 0$.

Let us compute

$$\begin{aligned} G(r, (\nu + \delta)\theta'(r)) - ((\nu + \delta)\theta'(r))' &= -\lambda\theta(r) + \frac{\nu^2}{\theta(r)} \\ &\quad - (\nu + \delta)^2 \frac{\theta'(r)^2}{\theta(r)} - (\nu + \delta)\theta''(r) \\ &= -\epsilon\theta(r) - \frac{1}{\theta} \left(\bar{\lambda}\theta^2(r) - \nu^2 + \nu^2\theta'(r)^2 + \nu\theta(r)\theta''(r) \right) \\ &\quad - \frac{\delta}{\theta} \left(\theta'(r)^2(\delta + 2\nu) + \theta''(r)\theta(r) \right). \end{aligned}$$

The second summand is less or equal than zero by hypothesis, therefore we can choose $\delta \in (0, 1 - \nu)$ arbitrarily small so that

$$G(r, (\nu + \delta)\theta'(r)) - ((\nu + \delta)\theta'(r))' < 0$$

on $(0, R)$. From Lemma 30 here below with $v = q$ and $w = (\nu + \delta)\theta'$, $a = 0$, $b = R$, we deduce that

$$q < (\nu + \delta)\theta'$$

as claimed.

In the case $R = \bar{R}$, if $u = \theta^\nu$, then $q = \nu\theta' \leq \theta'$ on $(0, \bar{R})$. □

Lemma 30 *Suppose v, w continuous on an interval $[a, b]$ and differentiable on (a, b) , and let $G : (a, b] \times \mathbb{R} \rightarrow \mathbb{R}$ continuous. Suppose that*

$$v(a) < w(a)$$

and

$$v'(r) - G(r, v(r)) < w'(r) - G(r, w(r)) \tag{41}$$

on $(a, b]$. Then $v < w$ on $[a, b]$.

Proof Assume that there exists $c \in (a, b]$ such that $v(c) = w(c)$ and $v - w < 0$ on $[a, c)$. Then clearly $v'(c) \geq w'(c)$. On the other hand, inserting $v(c) = w(c)$ in (41), we see that $v'(c) < w'(c)$. A contradiction. \square

B. 2 Eigenfunctions of the Steklov Problem

The same results of the previous subsection hold also for the Steklov problem (4). For the reader’s convenience, we briefly resume them, even if in this note we almost only consider the Steklov problem for domains in \mathbb{R}^2 .

When looking for solutions on $B(x_0, R)$ of the form $f(r, t) = u(r)e^{ikt}$ we get the following singular Sturm–Liouville problem, for $k \in \mathbb{Z}$:

$$\begin{cases} u'' + \frac{\theta'}{\theta} u' - \frac{(k-v)^2}{\theta^2} u = 0 & \text{in } (0, R), \\ u'(R) = \sigma u(R), \end{cases} \tag{42}$$

with the requirement that a solution is bounded near 0 (and therefore, that it vanishes at 0 when $v \notin \mathbb{Z}$).

We have the following

Lemma 31 *Let $v \in (0, \frac{1}{2}]$. Then problem (42) admits a unique solution $u_k(r)$ with $u_k(0) = 0$ for all $k \in \mathbb{Z}$. Let us set*

$$\eta_k = \frac{u'_k(R)}{u_k(R)}.$$

If $|k| \leq |h|$ then:

$$\eta_k \leq \eta_h.$$

Then the set $\{\sigma_k(B(x_0, R), A_{x_0,v})\}_{k=1}^\infty$ of the Steklov eigenvalues on $B(x_0, R)$ coincides with $\{\eta_k\}_{k \in \mathbb{Z}}$. In particular, $\min_k \{\eta_k\} = \eta_0$, therefore

$$\sigma_1(B(x_0, R), A_{x_0,v}) = \eta_0.$$

The eigenfunctions corresponding to an eigenvalues $\sigma = \eta_k$ of (4) are given by

$$\psi_k(r, t) = u_k(r)e^{ikt}.$$

In particular, an eigenfunction associated to $\sigma_1(B(x_0, R), A_{x_0,v})$ is given by $\psi_0(r, t)$ which is real and radial. The restrictions of $\{\psi_k(r, t)\}_k$ to $\partial B(x_0, R)$ form a orthonormal system in $L^2(\partial B(x_0, R))$.

Proof The proof is analogous of that of Lemma 24. It is sufficient to note that a bounded solution of (42) is of the form.

$$u_k(r) = e^{\int_R^r \frac{|k-v|}{\theta(s)} ds}.$$

The last statement is straightforward. □

B.3 The Magnetic Neumann Spectrum of the Unit Disk

We consider the particular case $M = B(x_0, 1)$, where $B(x_0, 1)$ is the unit disk in \mathbb{R}^2 centered at x_0 . We recall that $A = \nu dt$ and we assume $\nu \in (0, \frac{1}{2}]$. Here (r, t) are the standard polar coordinates in \mathbb{R}^2 centered at x_0 . In this case $\theta(r) = r$.

We can describe the spectrum of (3) on $B(x_0, 1)$ more explicitly:

Theorem 32 *Let $\nu \in (0, \frac{1}{2}]$. Then:*

- (i) *The spectrum of (3) on $B(x_0, 1)$ consists of the numbers λ_{kj} , $k \in \mathbb{Z}$, $j = 1, 2, \dots$, where $\sqrt{\lambda_{kj}} = z'_{|k-\nu|,j}$ and $z'_{\mu,j}$ denotes the j th positive zero of $J'_\mu(\sqrt{\lambda})$. Here $J_\mu(z)$ denotes the Bessel function of the first kind and order μ .*
- (ii) *The eigenspace associated to λ_{kj} is spanned by $\psi_{kj}(r, t) = J_{|k-\nu|}(\sqrt{\lambda_{kj}}r)e^{ikt}$.*
- (iii) *The lowest eigenvalue is $\lambda_1(B(x_0, 1), A_{x_0,\nu}) = \lambda_{01}$ and a first eigenfunction is $\psi_{01}(r, t) = J_\nu(\sqrt{\lambda_{01}}r)$. It is real and radial.*
- (iv) *The set of eigenfunctions $\{\psi_{kj}\}_{k,j}$ is complete in $L^2(B(x_0, 1))$.*

Proof We consider (35) with $\theta(r) = r$ and $R = 1$. As in Sect. 1, when looking for solutions of the form $f(r, t) = u(r)e^{ikt}$ in polar coordinates (r, t) , we obtain the following differential equation for u :

$$u'' + \frac{1}{r}u' + \left(\lambda - \frac{(k - \nu)^2}{r^2}\right)u = 0. \tag{43}$$

With the substitution $\sqrt{\lambda}r = z$, this equation can be recast to a standard Bessel equation. Therefore, a couple of linearly independent solutions of (43) is given by $J_{|k-\nu|}(\sqrt{\lambda}r)$, $J_{-|k-\nu|}(\sqrt{\lambda}r)$, since by hypothesis $\nu \notin \mathbb{Z}$. Therefore

$$u(r) = b_{k1}J_{|k-\nu|}(\sqrt{\lambda}r) + b_{k2}J_{-|k-\nu|}(\sqrt{\lambda}r).$$

Note that, while $J_{|k-\nu|}$ vanishes at the origin whenever $\nu \notin \mathbb{Z}$, the function $J_{-|k-\nu|}$ is unbounded near $r = 0$. Any bounded solution of (43) is given by setting $b_{k2} = 0$. Therefore, any solution of (35) with $\theta(r) = r$ has the form

$$u_k(r) = b_{k1}J_{|k-\nu|}(\sqrt{\lambda}r).$$

We impose now the Neumann boundary condition $u'_k(1) = 0$. We recall that for any $\mu > 0$, the function J_μ has an infinite set of positive zeroes, denoted by $z_{\mu,j}$, $j = 1, 2, \dots$. Moreover, the function J'_μ has an infinite number of positive zeroes as well, namely $z'_{\mu,j}$, $j = 1, 2, \dots$.

Then, points (i)–(iv) follow from Lemma 24 and Theorem 26 with $\sqrt{\lambda_{kj}} = z'_{|k-\nu|,j}$ and $\psi_{kj}(r, t) = J_{|k-\nu|}(\sqrt{\lambda_{kj}}r)e^{ikt}$. □

If $\nu \in (0, \frac{1}{2})$ then the set $\{|k - \nu|\}_{k \in \mathbb{Z}}$ is given by $\nu, 1 - \nu, 1 + \nu, 2 - \nu, 2 + \nu, \dots$. For example, when $\nu = \frac{1}{4}$ we have $\frac{1}{4}, \frac{3}{4}, \frac{5}{4}, \frac{7}{4}, \dots$. If $\nu \in (0, \frac{1}{2})$, we immediately see that $\lambda_1(B(x_0, 1), A_{x_0, \nu})$ is simple, and a first eigenfunction is real and radial:

$$\psi_{01}(r, t) = J_\nu(\sqrt{\lambda_{01}}r).$$

Such function is positive and increasing on $(0, 1)$.

When $\nu = \frac{1}{2}$, we have that the set of $\{|k - \nu|\}_{k \in \mathbb{Z}}$ is given by $\frac{1}{2}, \frac{1}{2}, \frac{3}{2}, \frac{3}{2}, \dots$, and the first eigenvalue has multiplicity 2 with eigenspace spanned by the eigenfunctions with $k = 0$ and $k = 1$:

$$\begin{cases} \psi_{01}(r, t) = J_{\frac{1}{2}}(\sqrt{\lambda_{01}}r), \\ \psi_{11}(r, t) = J_{\frac{1}{2}}(\sqrt{\lambda_{01}}r)e^{it}, \end{cases}$$

and $\lambda_{01} = \lambda_{11} = z'_{\frac{1}{2}, 1}$ is the first zero of $J'_{\frac{1}{2}}$.

B.4 The Magnetic Steklov Spectrum of the Unit Disk

In this section we investigate the spectrum of problem (4) on $B(x_0, 1) \subset \mathbb{R}^2$. Also in this section we assume that $\nu \in (0, \frac{1}{2}]$.

We apply the results of Sect. 1 with $\theta(r) = r$ and $R = 1$.

Theorem 33 *Let $\nu \in (0, \frac{1}{2}]$. Then:*

- (i) *The spectrum of (4) on $B(x_0, 1)$ consists of the numbers $\eta_k = |k - \nu|$, $k \in \mathbb{Z}$.*
- (ii) *The eigenspace associated to η_k is spanned by $\psi_k(r, t) = r^{|k-\nu|}e^{ikt}$.*
- (iii) *The lowest eigenvalue is $\sigma_1(B(x_0, 1), A_{x_0, \nu}) = \eta_0 = \nu$. A first eigenfunction is $\psi_1(r, t) = r^\nu$. It is real and radial.*
- (iv) *The restrictions of $\{\psi_{kj}\}_{k, j}$ to $\partial B(x_0, 1)$ is complete in $L^2(\partial B(x_0, 1))$.*

Proof It is sufficient to note that any solution to the differential equation (42) is of the form

$$u_k(r) = b_{k1}r^{|k-\nu|} + b_{k2}r^{-|k-\nu|},$$

and in order to have a solution of (42) we need to impose $b_{k2} = 0$. The theorem now easily follows from Lemma 31. □

When $\nu \in (0, \frac{1}{2})$, the first eigenvalue is $\nu > 0$. It is positive and simple and a corresponding eigenfunction is given by r^ν and it is radial, positive and increasing on $(0, 1)$.

We can list the Steklov eigenvalues $\{\sigma_k(B(x_0, 1), A_{x_0, \nu})\}_{k=1}^\infty$ as follows:

$$\nu, 1 - \nu, 1 + \nu, 2 - \nu, 2 + \nu, \dots$$

If $\nu = \frac{1}{2}$ all eigenvalues are double and are given by $\frac{1}{2}, \frac{1}{2}, \frac{3}{2}, \frac{3}{2}, \dots$

In particular, $\sigma_1(B(x_0, 1), A_{x_0,1/2}) = \sigma_2(B(x_0, 1), A_{x_0,1/2}) = \frac{1}{2}$ and two linearly independent eigenfunctions are $u_1(x) = \sqrt{|x|}$, $u_2(x) = \sqrt{|x|} \left(\frac{x_1}{|x|} + i \frac{x_2}{|x|} \right)$. However, also in this case there exist a radial, real eigenfunction associated with $\sigma_1(B(x_0, 1), A_{x_0,v})$, positive and increasing on $(0, 1)$.

C Conformal Invariance of Magnetic Energy

Let (Ω_1, g) , (Ω_2, g) be two 2-dimensional manifolds, and let

$$\Phi : (\Omega_1, g_1) \rightarrow (\Omega_2, g_2)$$

be a conformal map between them. We can assume then that $\Phi^* g_2 = e^{2f} g_1$, where f is the conformal factor. Let ω be a 1-form on Ω_2 . We start from:

Lemma 34 *We have*

(i) *In the above notation*

$$\delta_{g_1} \Phi^* \omega = e^{2f} \Phi^* (\delta_{g_2} \omega).$$

In particular, ω is co-closed if and only if $\Phi^ \omega$ is co-closed.*

(ii) *For any complex 1-form ω on Ω_2 , one has:*

$$\int_{\Omega_1} |\Phi^* \omega|_{g_1}^2 d\mu_{g_1} = \int_{\Omega_2} |\omega|_{g_2}^2 d\mu_{g_2}.$$

Proof The first fact is standard; for the second, fix a g_1 -orthonormal basis (e_1, e_2) and let

$$E_1 = d\Phi(e^{-f} e_1), \quad E_2 = d\Phi(e^{-f} e_2).$$

Using these orthonormal frames to compute the norms we end-up with the identity:

$$|\Phi^* \omega|_{g_1}^2 = e^{2f} (|\omega|_{g_2}^2 \circ \Phi).$$

Integrating the identity on Ω_1 and using the change of variables formula, we obtain the assertion. □

Now fix a potential one-form A on Ω_2 . We get a potential one-form $\hat{A} = \Phi^* A$ on Ω_1 by pull-back. Recall the magnetic gradient on Ω_2 :

$$d^A u = du - iuA$$

for a complex valued function u on Ω_2 . Consider the function

$$\hat{u} = u \circ \Phi = \Phi^* u$$

on Ω_1 . Then, since $d\Phi^* = \Phi^*d$:

$$d^{\hat{A}}\hat{u} = d\Phi^*A\Phi^*u = d\Phi^*u - i\Phi^*u\Phi^*A = \Phi^*(du - iuA) = \Phi^*d^A u.$$

Applying the L^2 invariance property to $\omega = d^A u$ we conclude with the following fact, expressing the conformal invariance of the magnetic energy.

Proposition 35 *Let $\Phi : \Omega_1 \rightarrow \Omega_2$ be a conformal map between 2-manifolds, let A be a potential one-form on Ω_2 and let $\hat{A} = \Phi^*A$. For any complex function u on Ω_2 , let $\hat{u} = u \circ \Phi$. Then we have:*

$$\int_{\Omega_1} |d^{\hat{A}}\hat{u}|^2 d\mu_1 = \int_{\Omega_2} |d^A u|^2 d\mu_2.$$

C.1 Aharonov–Bohm Potentials

Let now (Ω, x_0) be a simply connected plane domain punctured at x_0 , with Aharonov–Bohm potential having pole at x_0 and flux ν . We can assume without loss of generality that x_0 is the origin. Then:

$$A = \frac{\nu}{r^2}(-ydx + xdy). \tag{44}$$

Take the unit disk D centered at the origin. By the standard theory, there is a conformal map $\Phi : \Omega \rightarrow D$ fixing the origin. Note that the form Φ^*A on Ω is closed (clear) and co-closed (from previous section, because Φ is conformal). Then, Φ^*A is harmonic. It is also clear that Φ^*A has flux ν around the pole, the origin of \mathbb{R}^2 (by elementary change of variable in dimension one). It follows that Φ^*A differs from A by an exact 1-form, and gauge invariance applies. Precisely:

Lemma 36 *Let Ω be a plane domain and let D be the unit disk, both punctured at $O \in \Omega$. Let $\Phi : \Omega \rightarrow D$ be the unique conformal map fixing O . If A is the canonical Aharonov–Bohm potential with pole O and flux ν , as in (44), then*

$$\sigma_k(\Omega, A) = \sigma_k(\Omega, \Phi^*A)$$

for all k .

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